

# Predictions for the LHC: an Overview

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CERN Theory Institute, May-June 2007

***Heavy Ion Collisions at the LHC - Last Call for Predictions***

organized by N.A., N. Borghini, S. Jeon and U.A. Wiedemann,

arXiv:0711.0974 (93/170/82 contributions/authors/institutes).

# Contents:

A. The bulk:

- ↑ A.1. Multiplicities.
- A.2. Azimuthal asymmetries.
- A.3. Hadronic flavor observables.
- ↓ A.4. Correlations at low  $p_T$ .

B. Hard and  
EM probes:

- ↑ B.1. High- $p_T$  observables and jets.
- B.2. Heavy quarks and quarkonium.
- ↓ B.3. Leptonic probes and photons.

Only predictions for PbPb.

See also Borghini and Wiedemann, arXiv:0707.0654.

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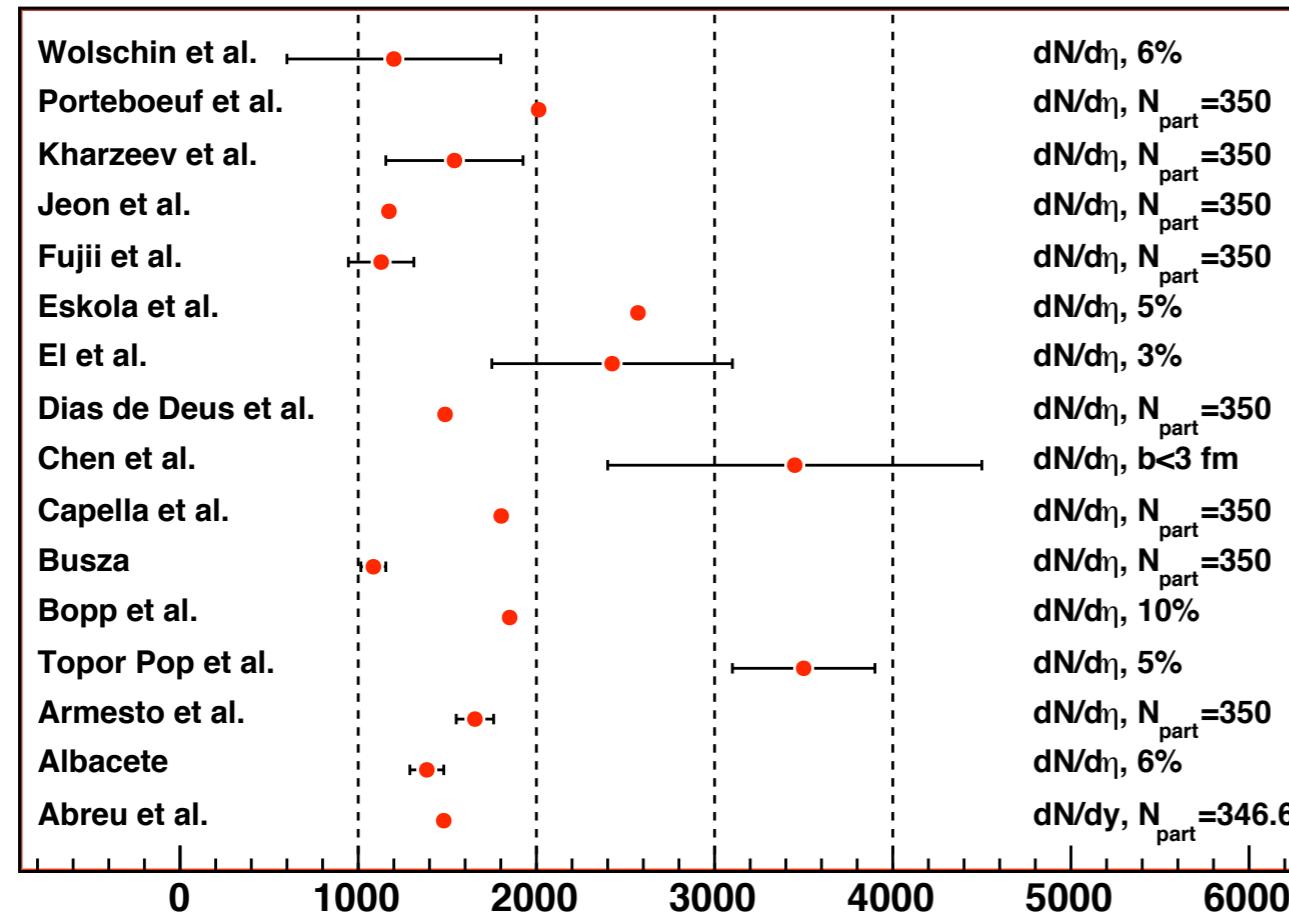
Only predictions for PbPb.

See also Borghini and Wiedemann, arXiv:0707.0654.

⇒ Remark: I assume that a RHIC- (and possibly SPS-) tested model can be reliably extrapolated to the LHC. The huge lever arm in energy will eventually falsify some of them.

# A. I. Multiplicities (I):

Charged multiplicity for  $\eta=0$  in central Pb+Pb at  $\sqrt{s_{NN}}=5.5$  TeV



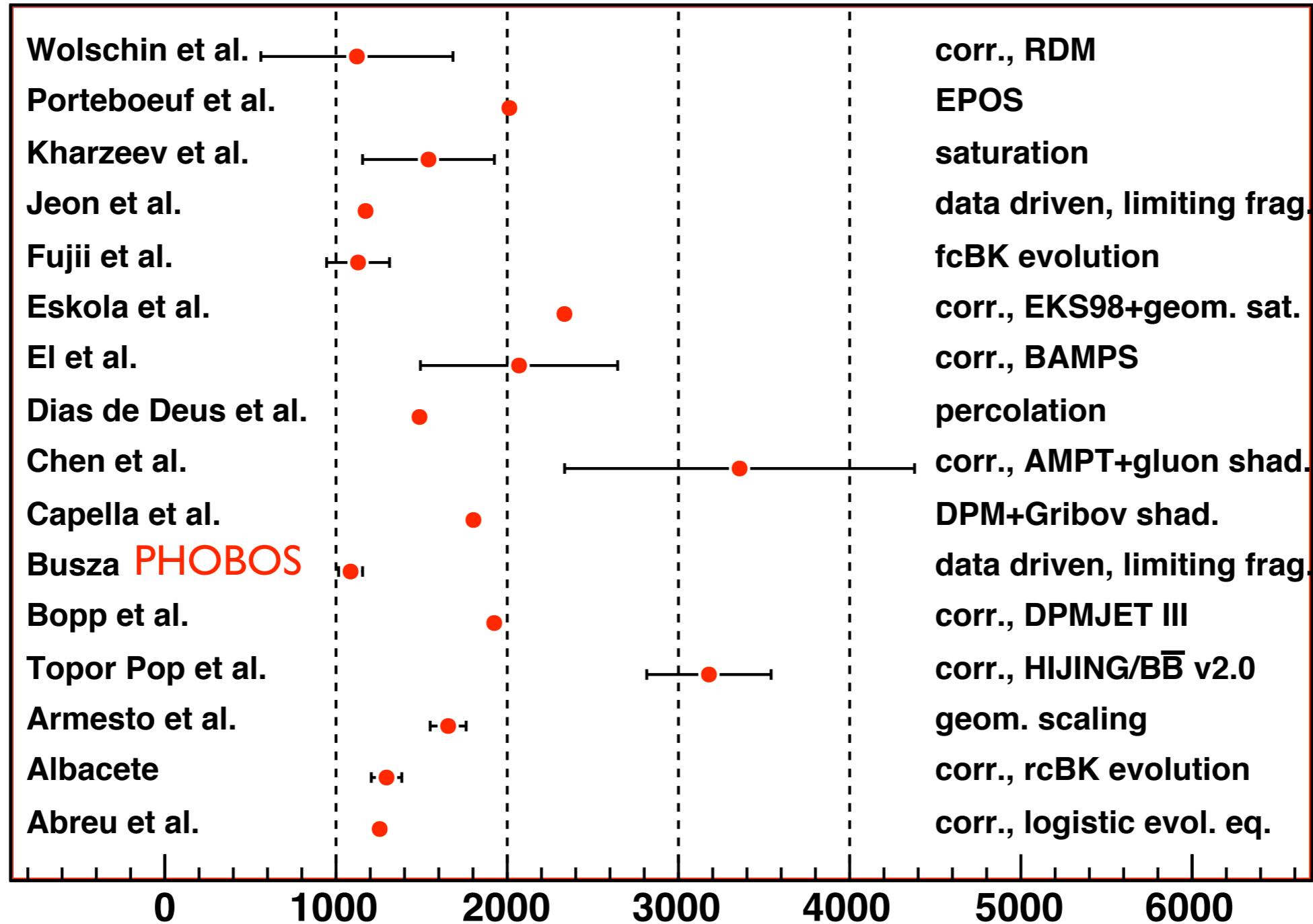
I **st-day**  
observable,  
key input in  
almost all  
other  
predictions.

To unify the discussion, I ‘rescale’  
to  $dN_{ch}/d\eta|_{\eta=0}$  for  $N_{part}=350$   
using a Monte Carlo  
(Amelin et al., EPJC22(2001)149):

%	$\langle b \rangle$ (fm)	$\langle N_{part} \rangle$	$\langle N_{coll} \rangle$	$dN_{ch}/dy _{y=0}$	$dN_{ch}/d\eta _{\eta=0}$
0-3	1.9	390	1584	3149	2633
0-5	2.4	375	1490	2956	2472
0-6	2.7	367	1447	2872	2402
0-7.5	3.0	357	1390	2759	2306
0-8.5	3.1	350	1354	2686	2245
0-9	3.2	347	1336	2649	2214
0-10	3.4	340	1303	2583	2159

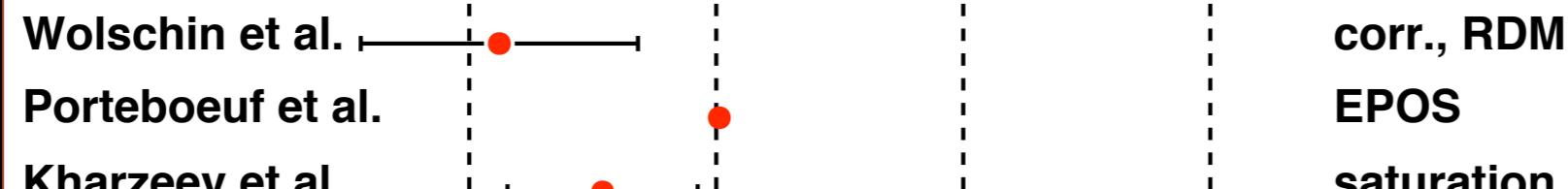
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$dN_{ch}/d\eta|_{\eta=0}$  in Pb+Pb at  $\sqrt{s_{NN}}=5.5$  TeV for  $N_{part}=350$

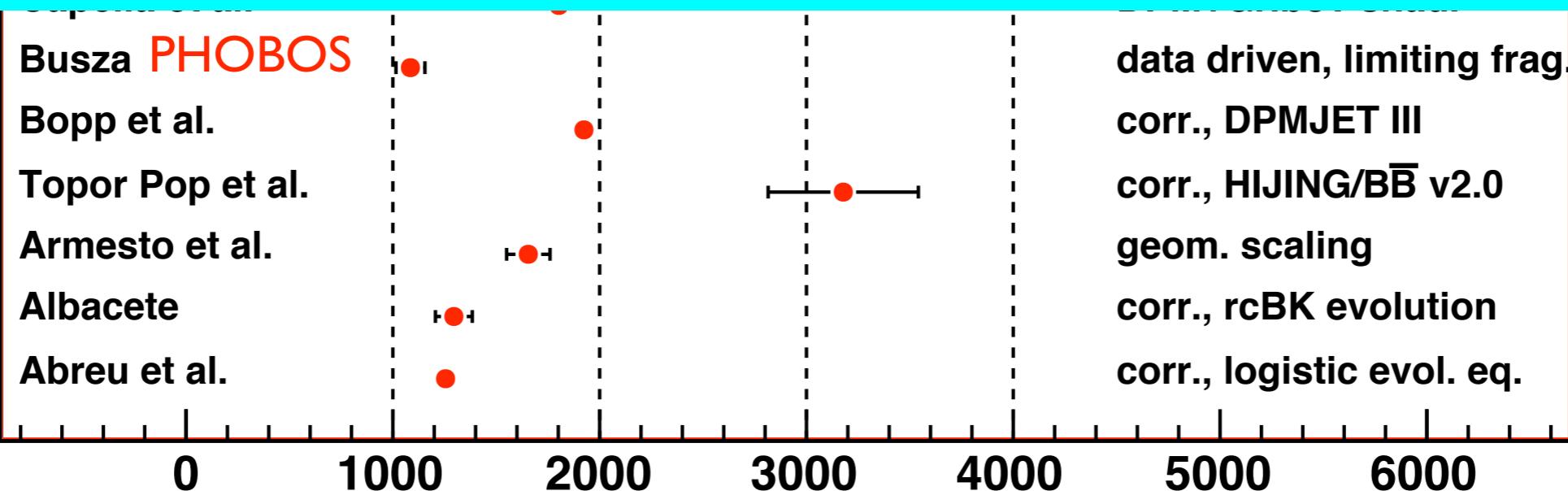


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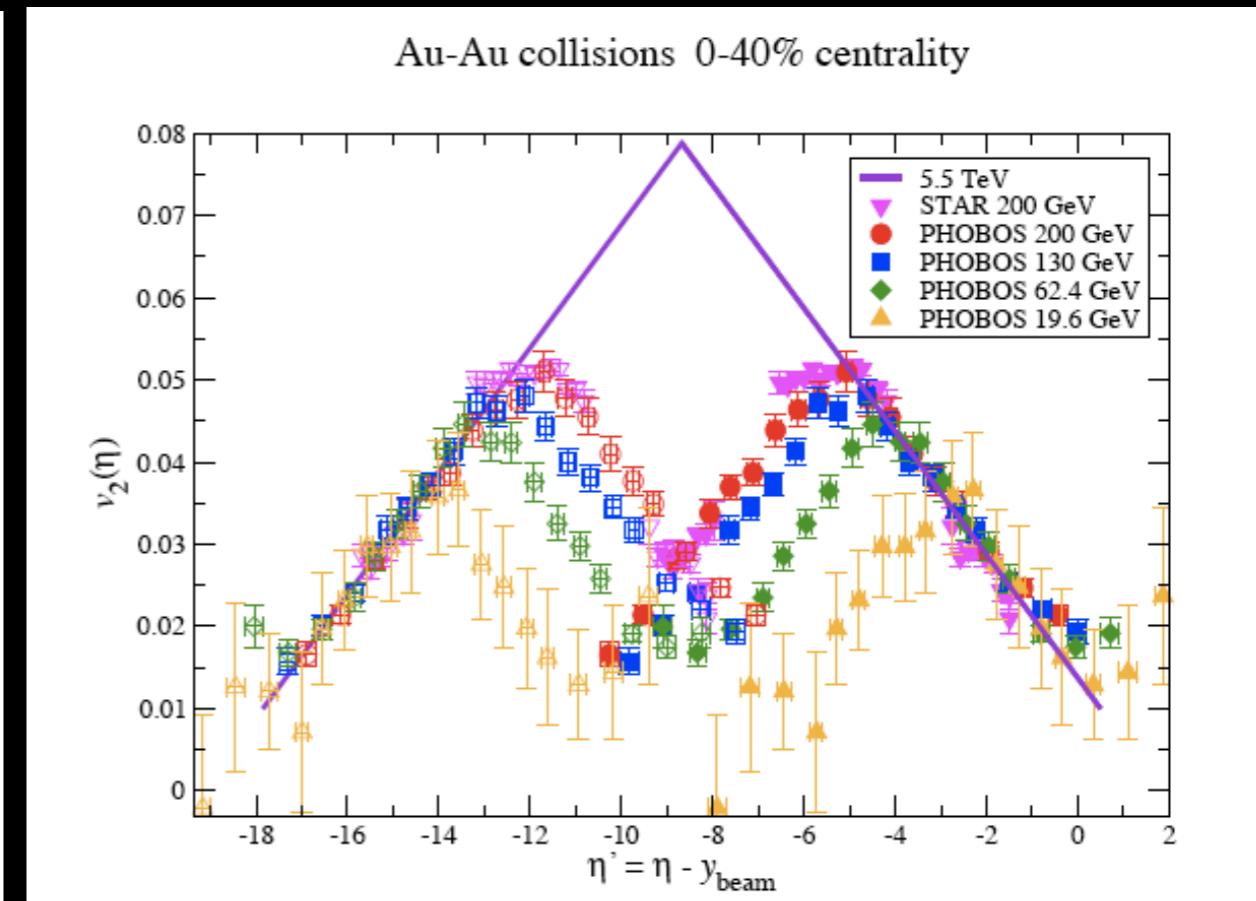
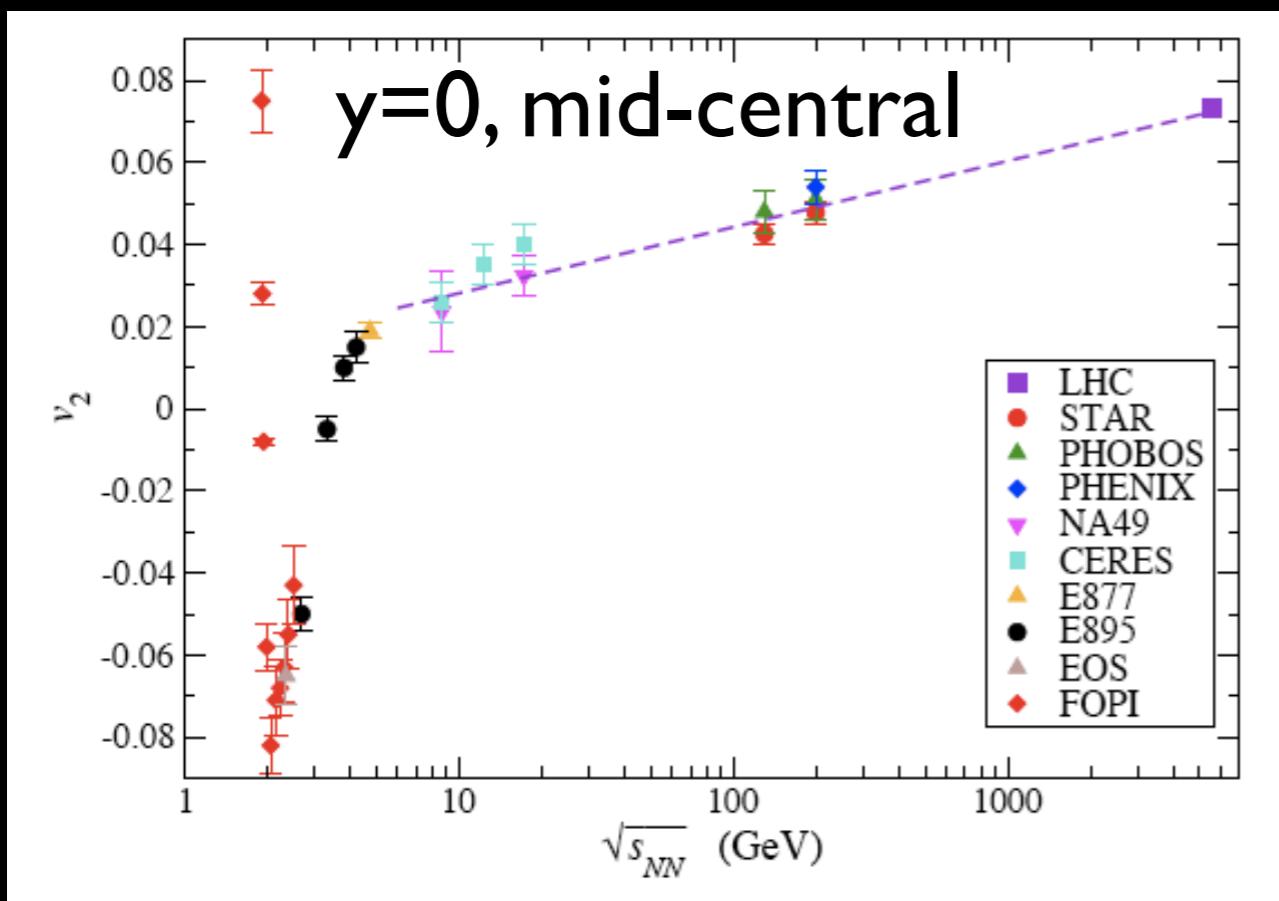


- ⇨ Generic expectation: less than 2000.
- ⇨ Most models include now a large degree of collectivity: strong gluon shadowing, strong color fields, saturation,...



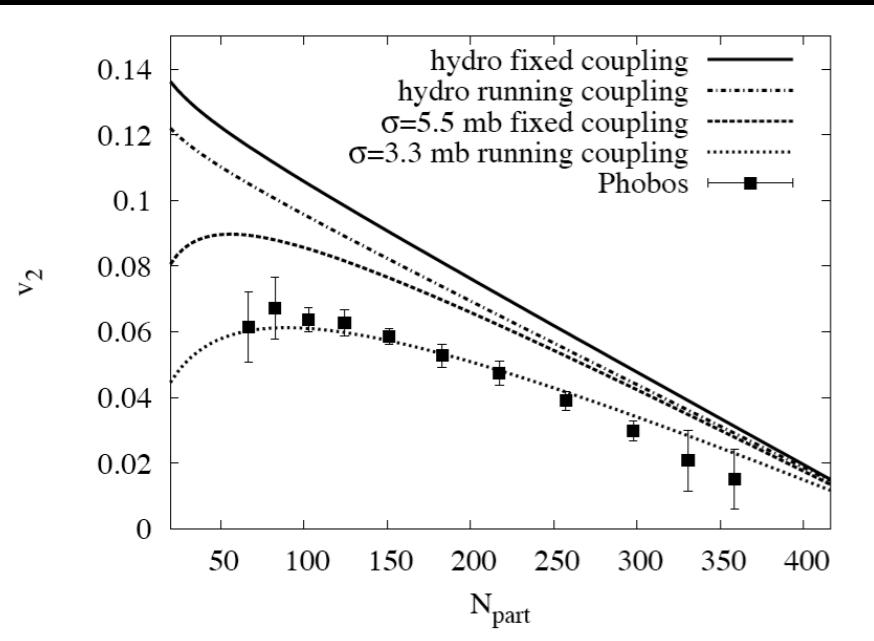
## A.2. Azimuthal asymmetries: $v_2$

$p_T$ -integrated  $v_2$  is expected to increase from RHIC to the LHC: origin? Hydro predictions do not coincide with naive extrapolations.

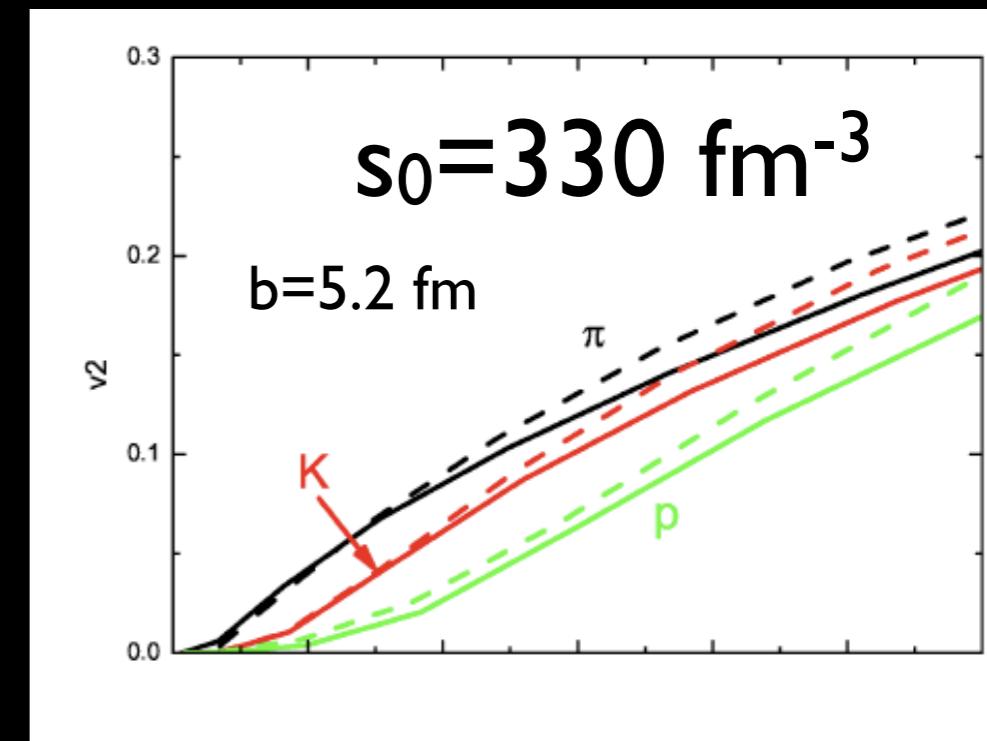


From BW.

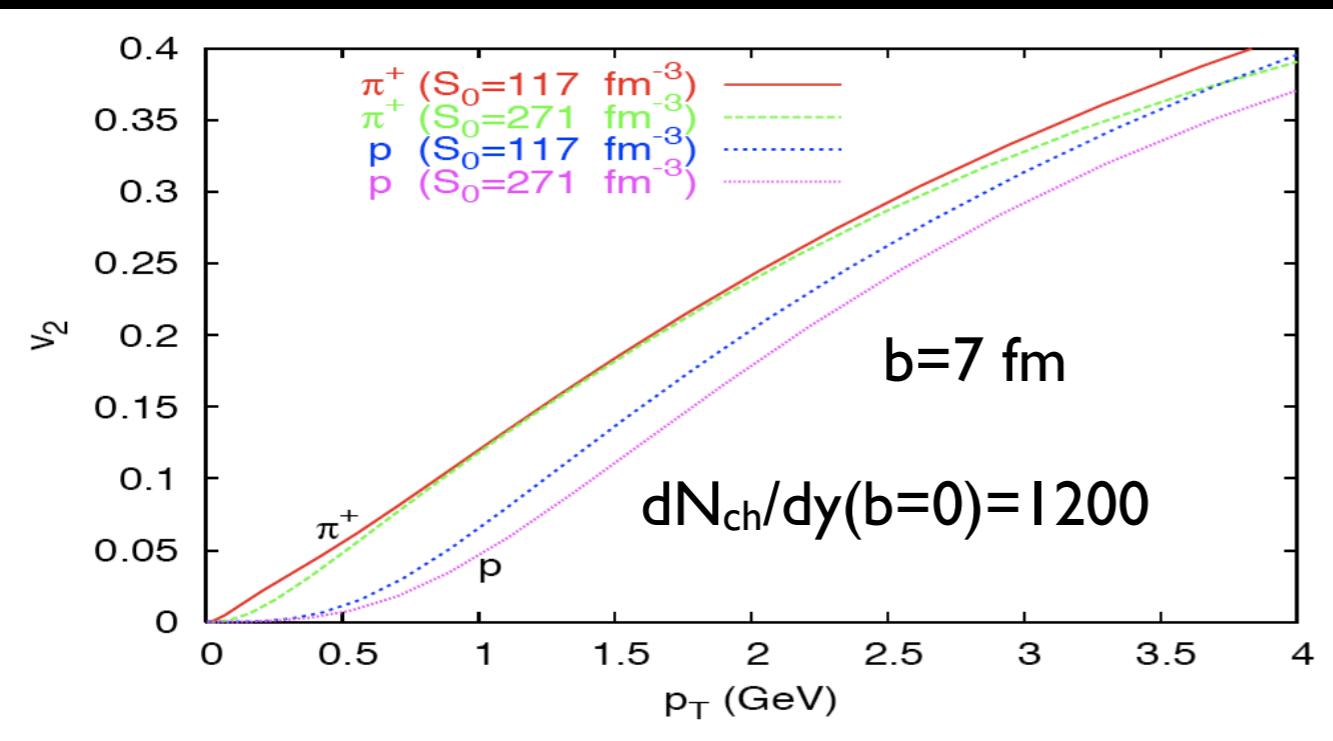
# A.2. Azimuthal asymmetries: $v_2$ in hydro



Drescher et al.,  
smaller  
deviation  
from **ideal** hydro  
limit than at  
RHIC.

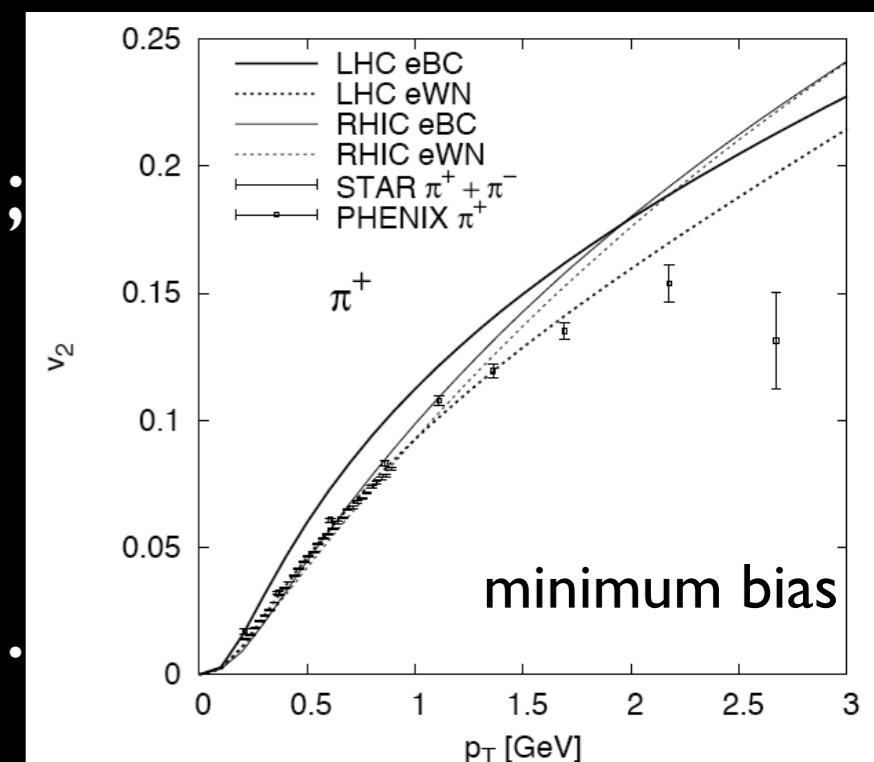


Bluhm et al., QPM EOS,  
results < RHIC at low  $p_T$ .

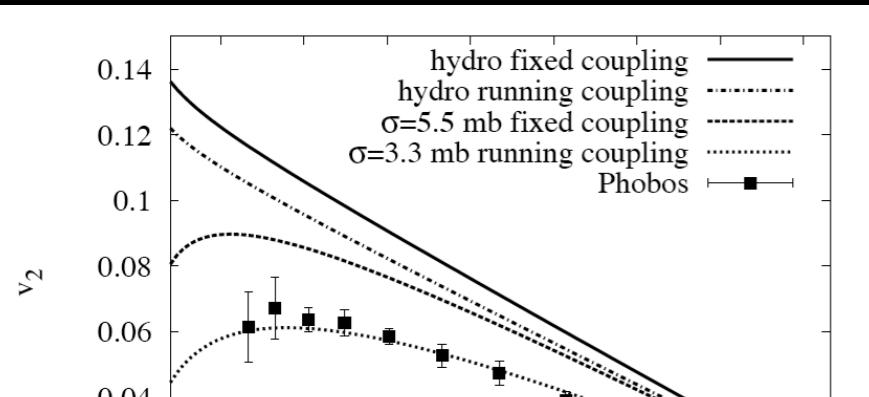


Kestin et al.

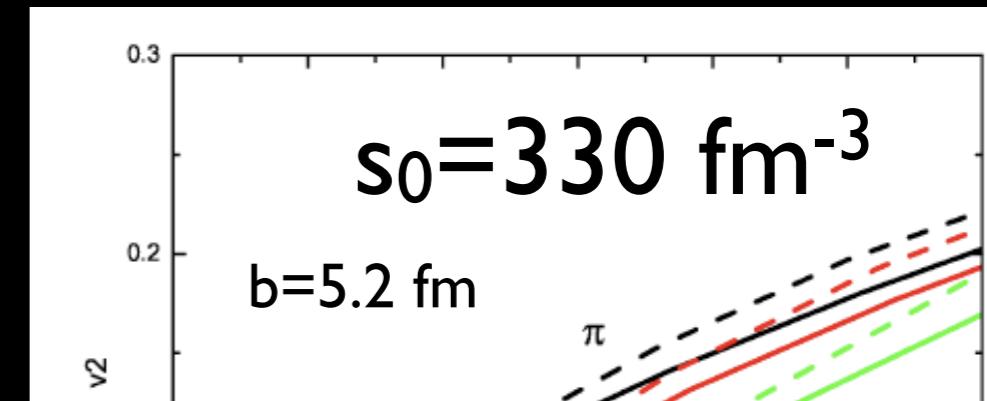
Eskola et al.;  
hydro valid  
for  $p_T < 4$   
GeV (also  
Arleo et al.).



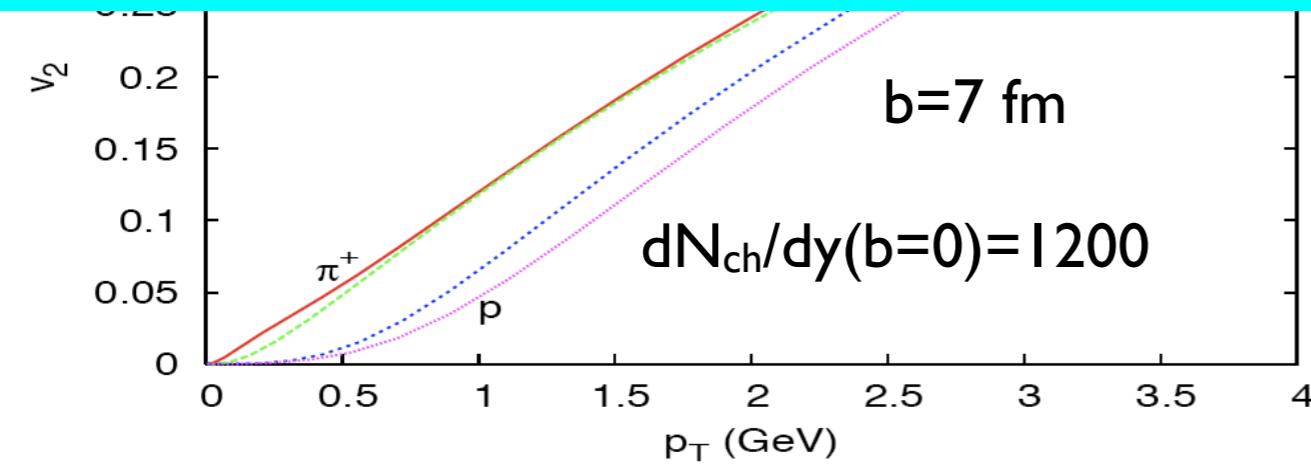
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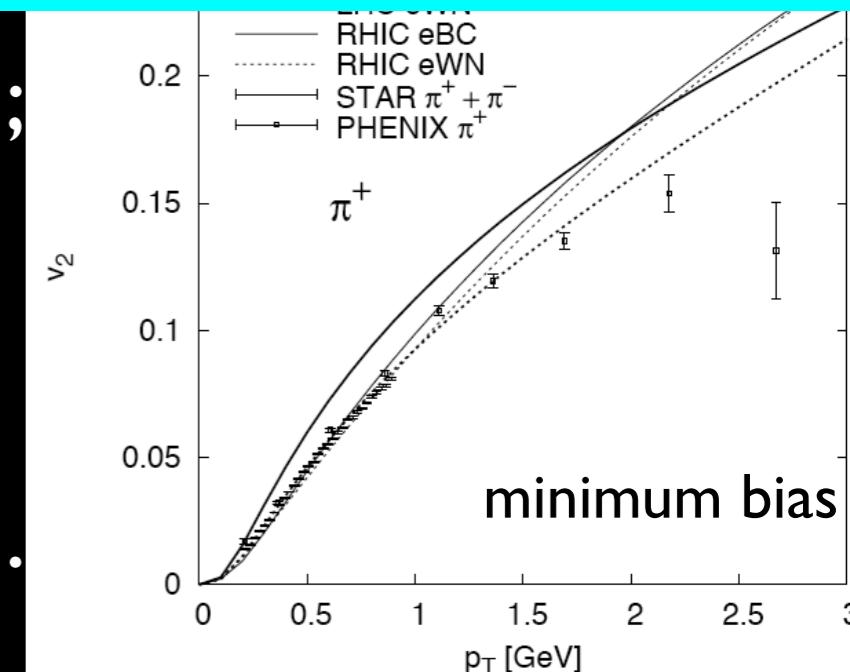


- ⇨ **Generic expectation:  $v_2$  similar or slightly decreasing at low  $p_T$ .**
- ⇨ **A strong decrease would probably signal an increase in  $\eta/s$ , but initial conditions have to be settled.**

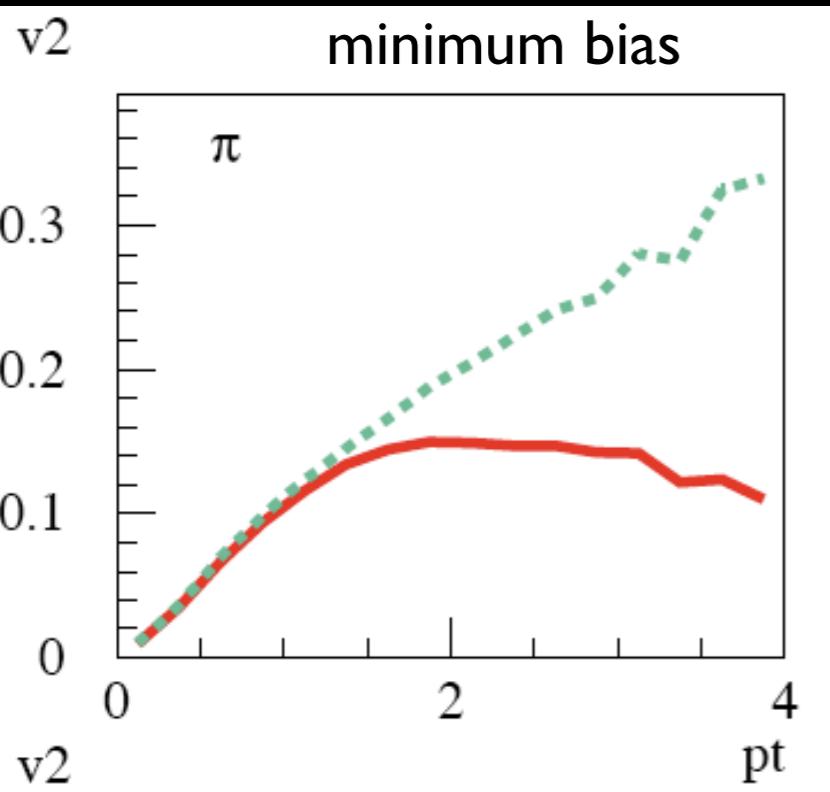


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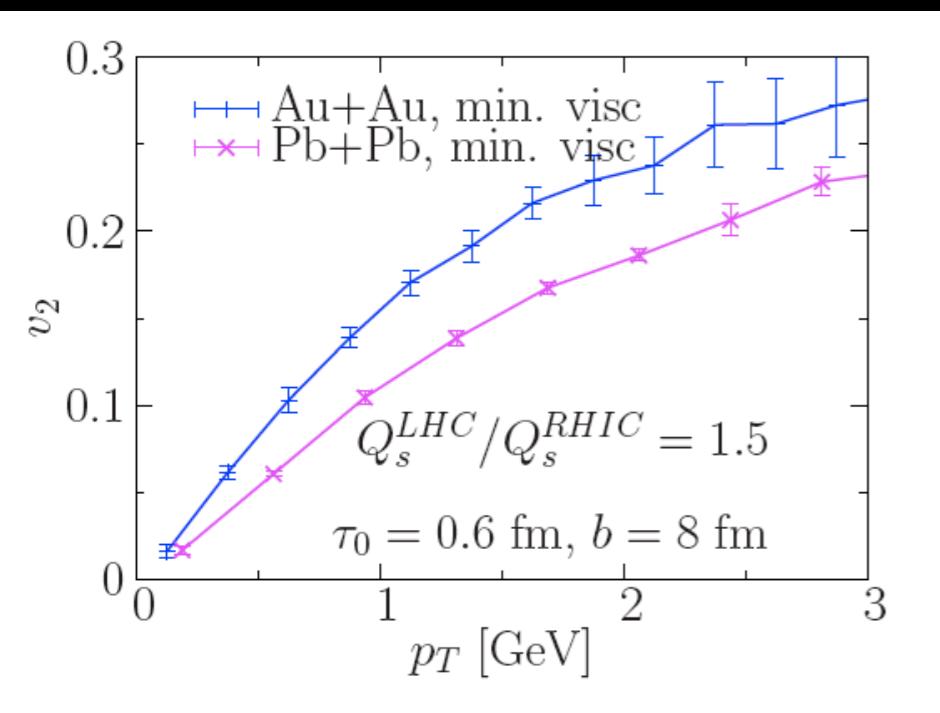
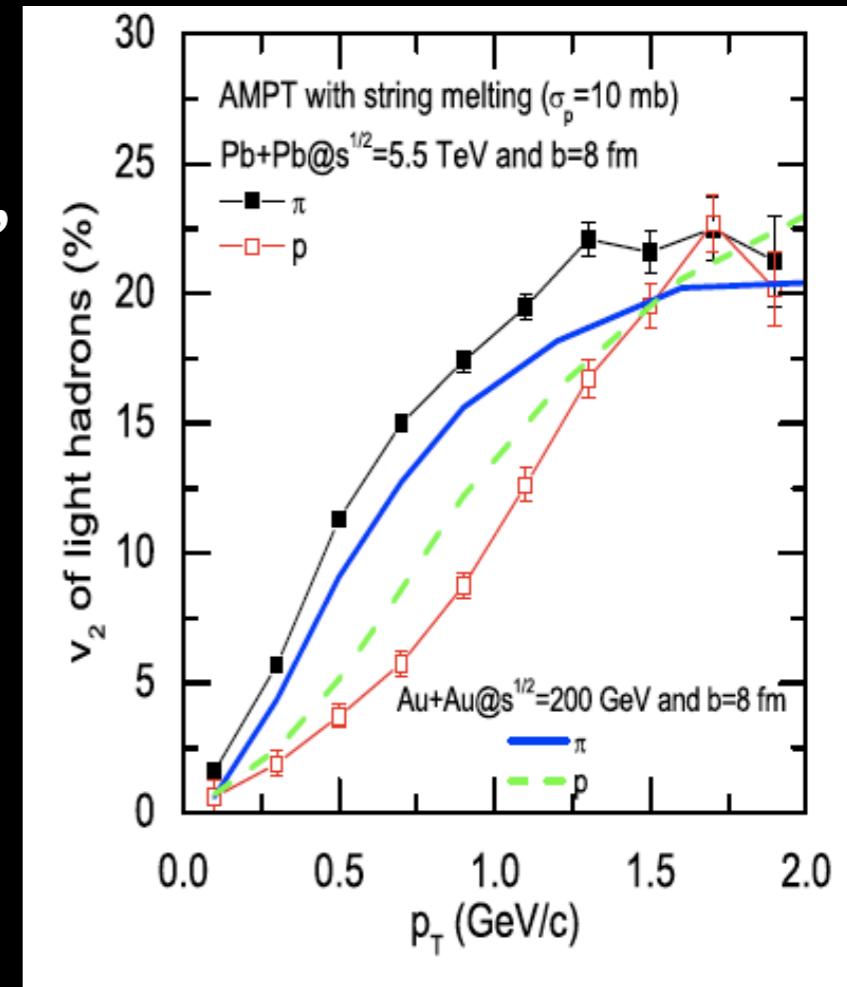


# A.2. Azimuthal asymmetries: $v_2$ in others

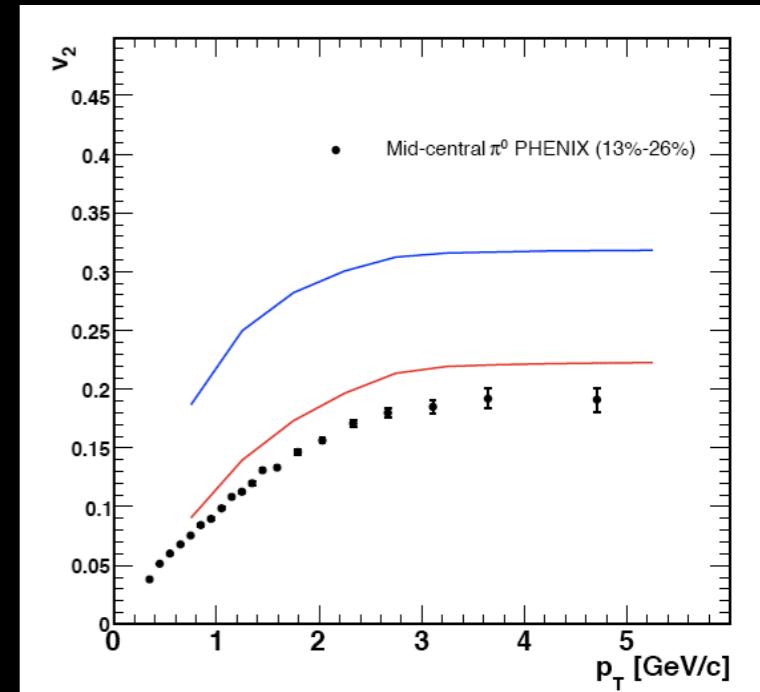


Porteboeuf  
et al.,  
EPOS,  
hydro  
core,  
 $v_2 \sim$  RHIC.

Chen et al.,  
AMPT,  
parton +  
hadron  
transport.

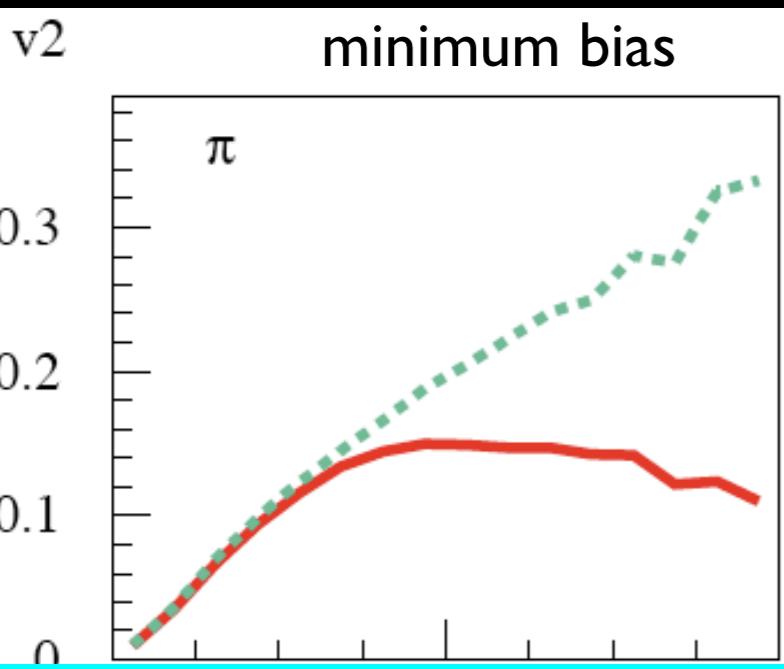


Molnar, MPC  
parton cascade,  
fixed  $\eta/s=0.08$ .



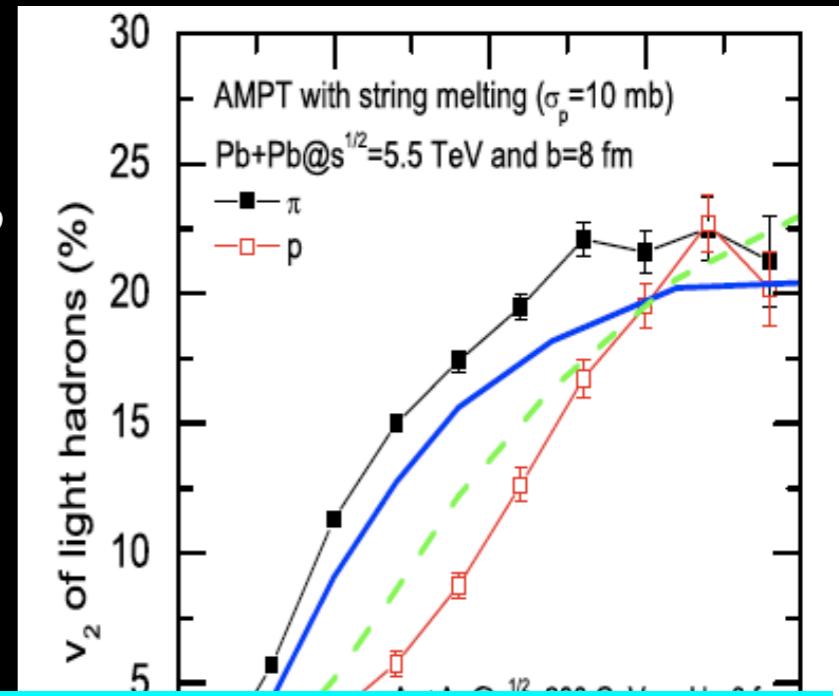
Capella et al., comovers.

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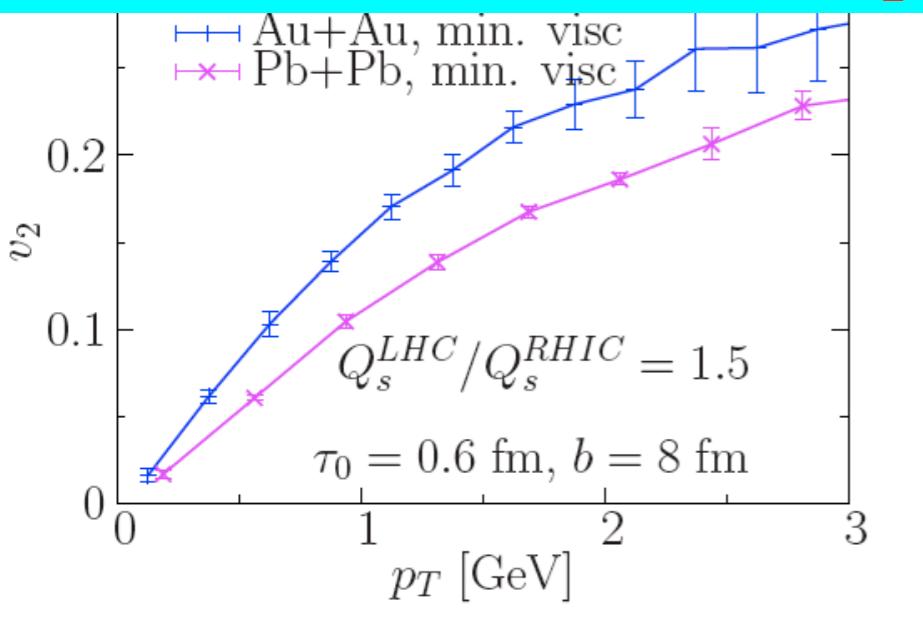


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et al.,  
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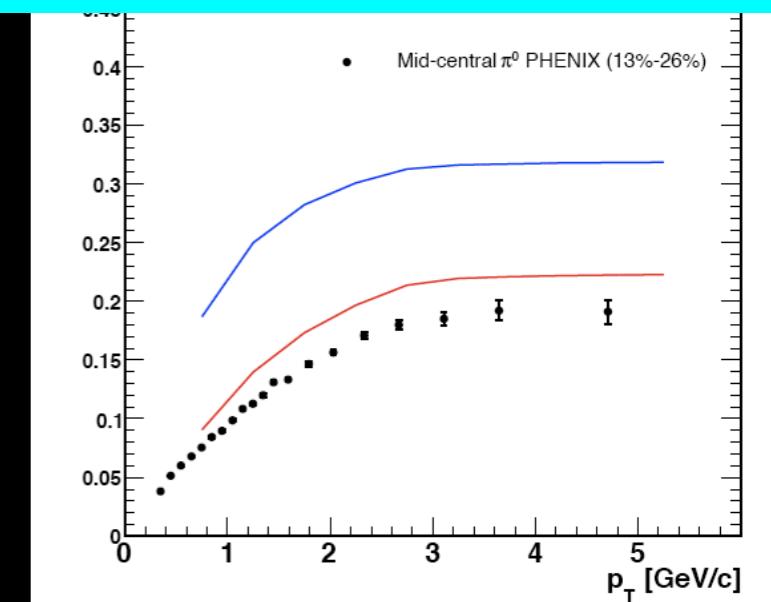
Chen et al.,  
AMPT,  
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- **Generic expectation:  $v_2$  increases at low  $p_T$ .**
- **A strong increase is not expected in any hydro description.**



Molnar, MPC  
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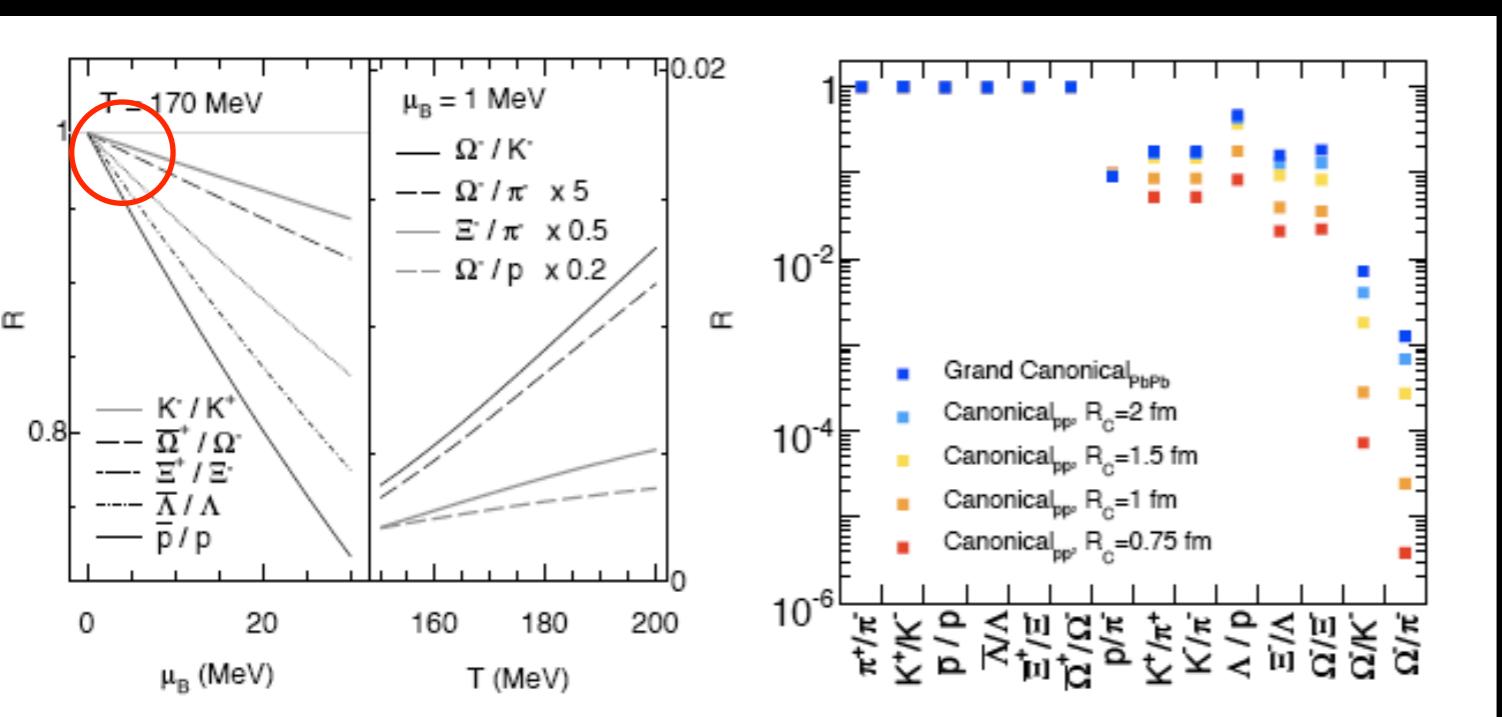
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# A.3. Hadr. flavor observ.: statistical

$\pi^-/\pi^+$	$K^-/K^+$	$\bar{p}/p$	$\bar{\Lambda}/\Lambda$	$\bar{\Xi}/\Xi$	$\bar{\Omega}/\Omega$
1.001(0)	0.993(4)	$0.948^{-0.015}_{+0.008}$	$0.997^{-0.011}_{+0.004}$	$1.005^{-0.007}_{+0.001}$	1.013(4)

Andronic et al., equilibrium values for  
 $\mu_b=0.8 \text{ MeV}, T=161 \text{ MeV}$ .

Kraus et al., (grand-)canonical, T and  
 $R_c$  may be determined.



$T[\text{MeV}]$	140*	140*	162*
$dV/dy [\text{fm}^3]$	2036	4187	6200*
$dS/dy$	7517	15262	18021
$dh_{\text{ch}}/dy$	1150*	2351	2430
$dh_{\text{ch}}^{\text{vis}}/dy$	1351	2797*	2797
$1000 \cdot (\lambda_{q,s} - 1)$	5.6*, 2.1	5.6*, 2.1	5.6*, 2.0
$\mu_{B,S}[\text{MeV}]$	2.4, 0.5	2.3, 0.5	2.7, 0.6
$\gamma_{q,s}$	1.62, 2.42	1.6*, 2.6	1*, 1*
$s/S$	0.034*	0.037*	0.025
$E/b$	420*	428	408
$E/TS$	1.02	1.05	0.86
$P/E$	0.165	0.164	0.162
$E/V[\text{MeV/fm}^3]$	530	538	400
$P[\text{MeV}]$	87	88	65
$p$	25/45	49/95	66/104
$b - \bar{b}$	2.6	5.3	6.1
$(b + \bar{b})/h^-$	0.335	0.345	0.363
$0.1 \cdot \pi^\pm$	49/67	99/126	103/126
$K^\pm$	94	207	175
$\phi$	14	33	23
$\Lambda$	19/28	41/62	37/50
$\Xi^-$	4	9.5	5.8
$\Omega^-$	0.82	2.08	0.98

Rafelski et al., non-equilibrium scenarios; non-strange resonances reduced.

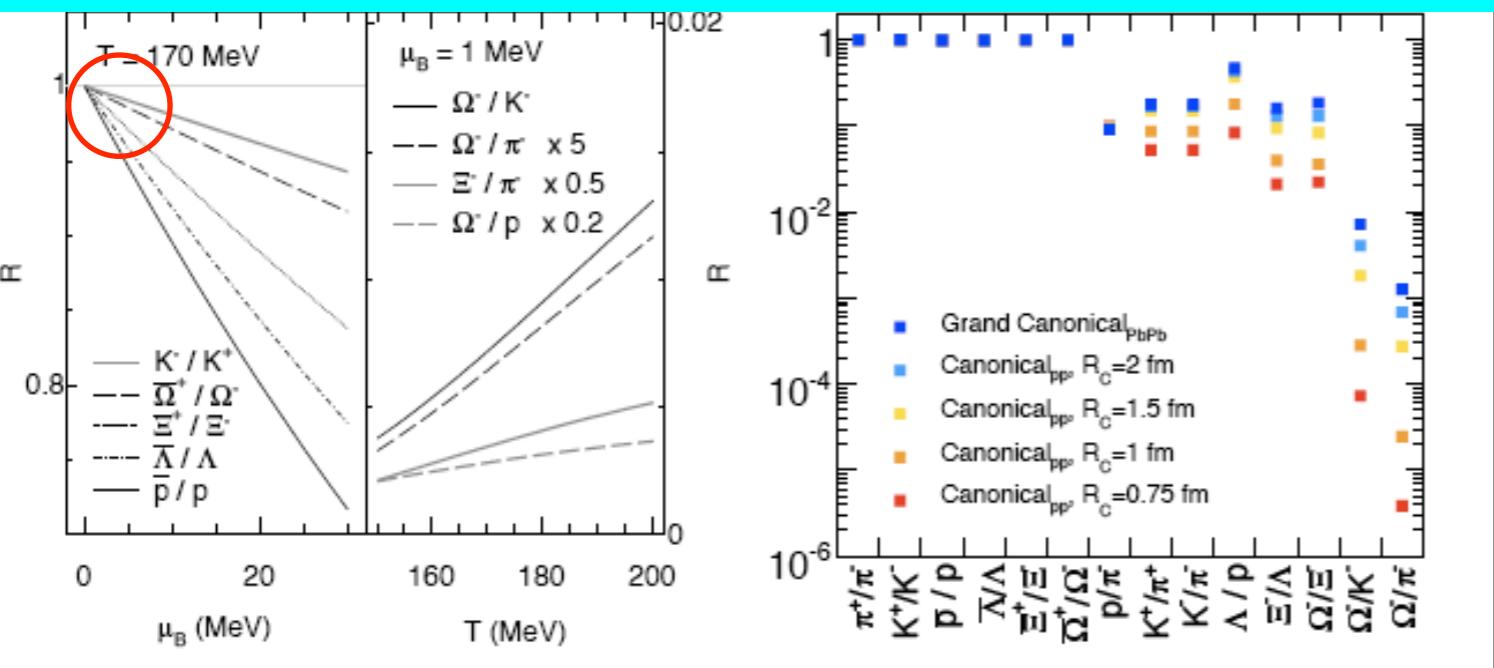
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$1000 \cdot (\lambda_{q,s} - 1)$	5.6*, 2.1	5.6*, 2.1	5.6*, 2.0

- ▷ Different statistical scenarios may be distinguished.
- ▷ This becomes of great importance for open charm and charmonium: different scenarios lead to marked differences in production.

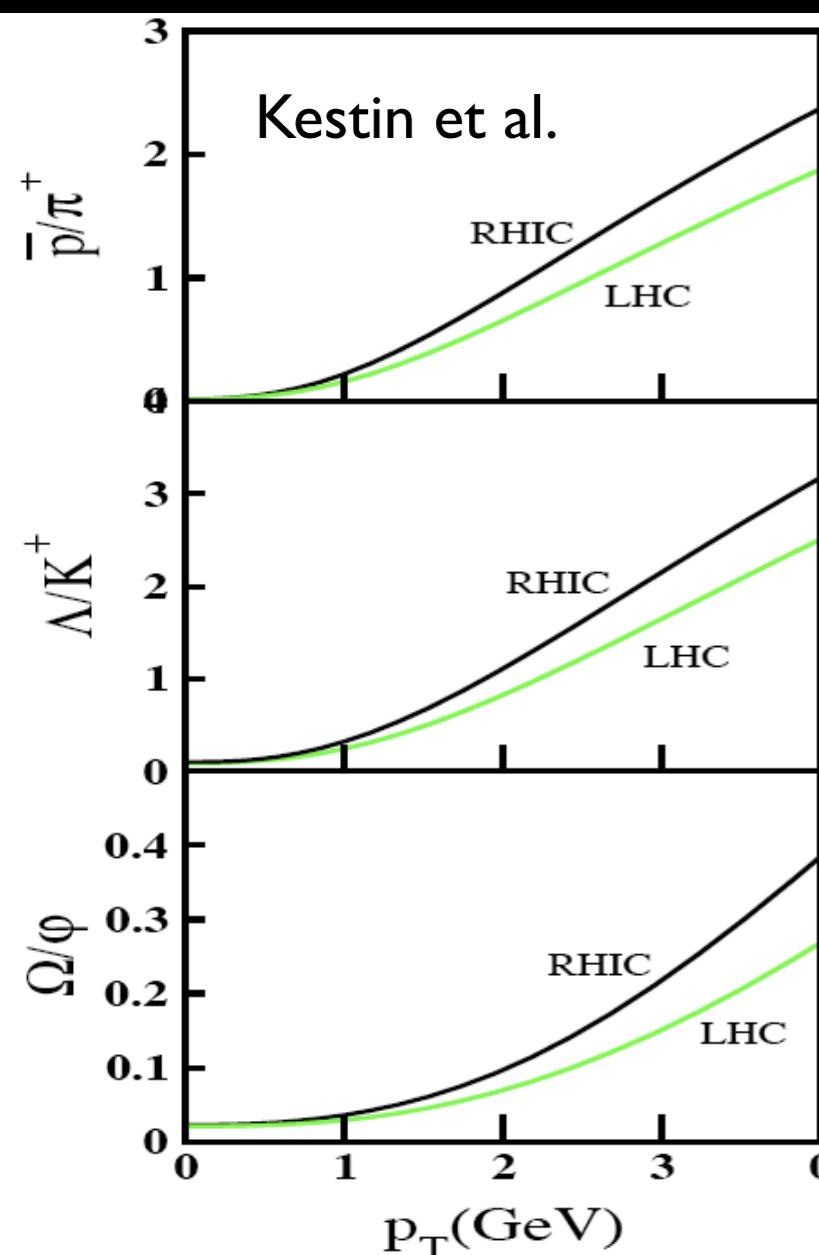


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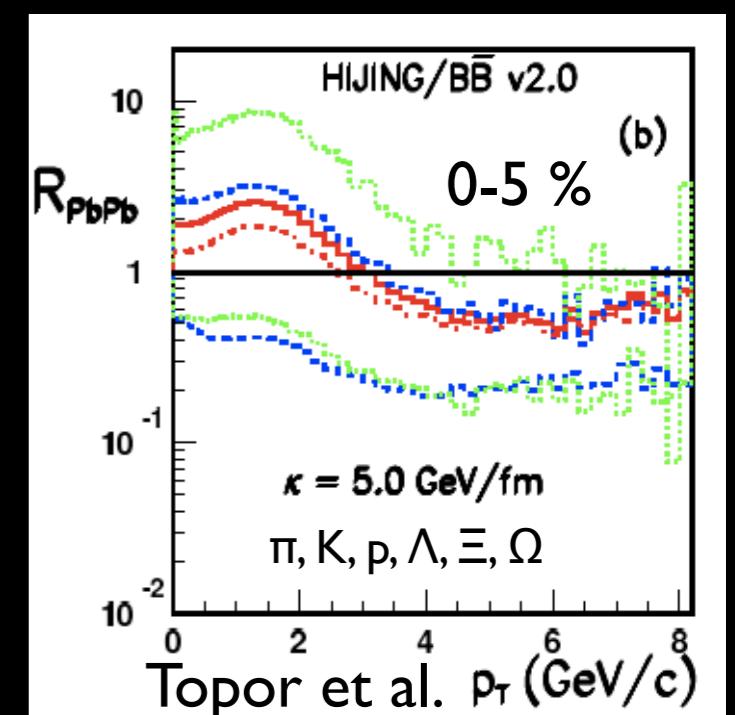
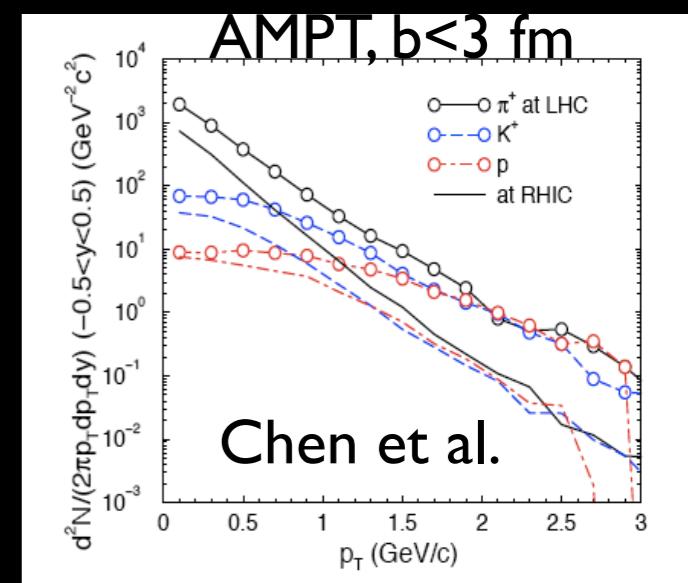
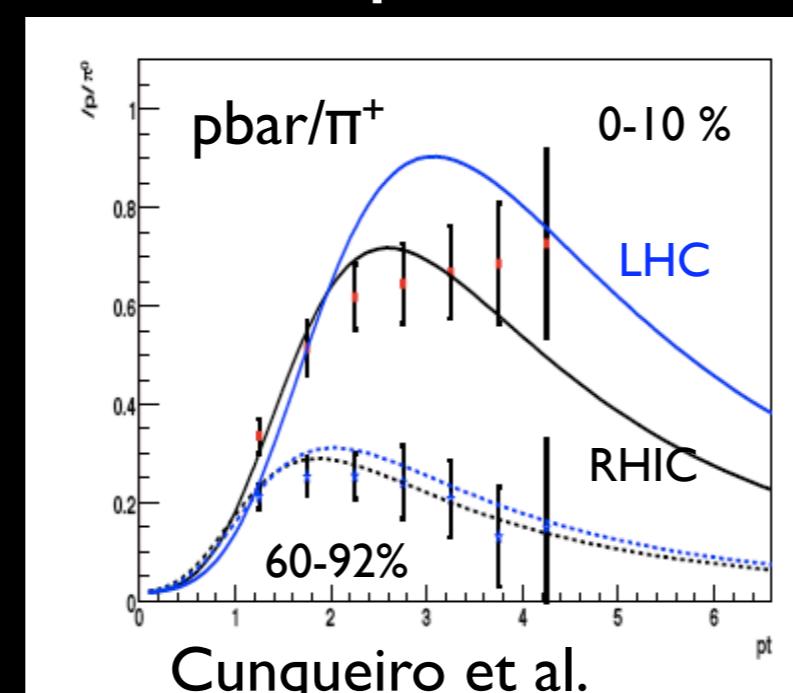
Rafelski et al., non-equilibrium scenarios; non-strange resonances reduced.

# A.3. Hadr. flavor observ.: baryons at low p<sub>T</sub>

p-pbar<4 at  $\eta=0$  (Bj: Topor Pop et al., Bopp et al.; hydro: Eskola et al.; EPOS: Porteboeuf et al., RDM: Wolschin et al.).



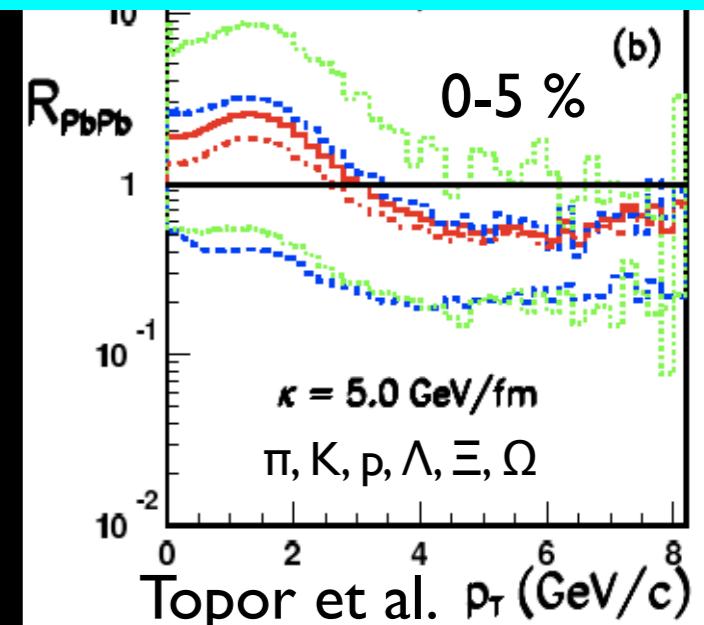
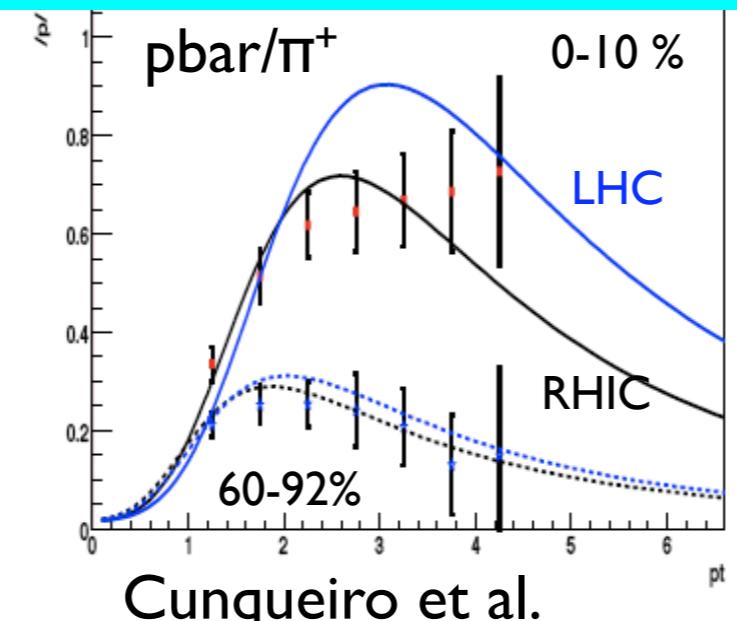
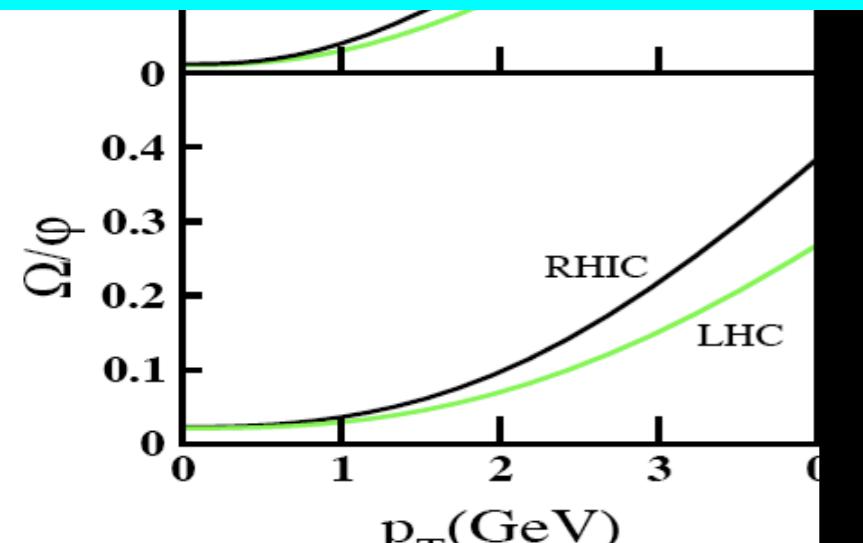
**Hydro and recombination predict larger baryon/meson ratios than models with higher string tension; the latter predicts Cronin for protons.**



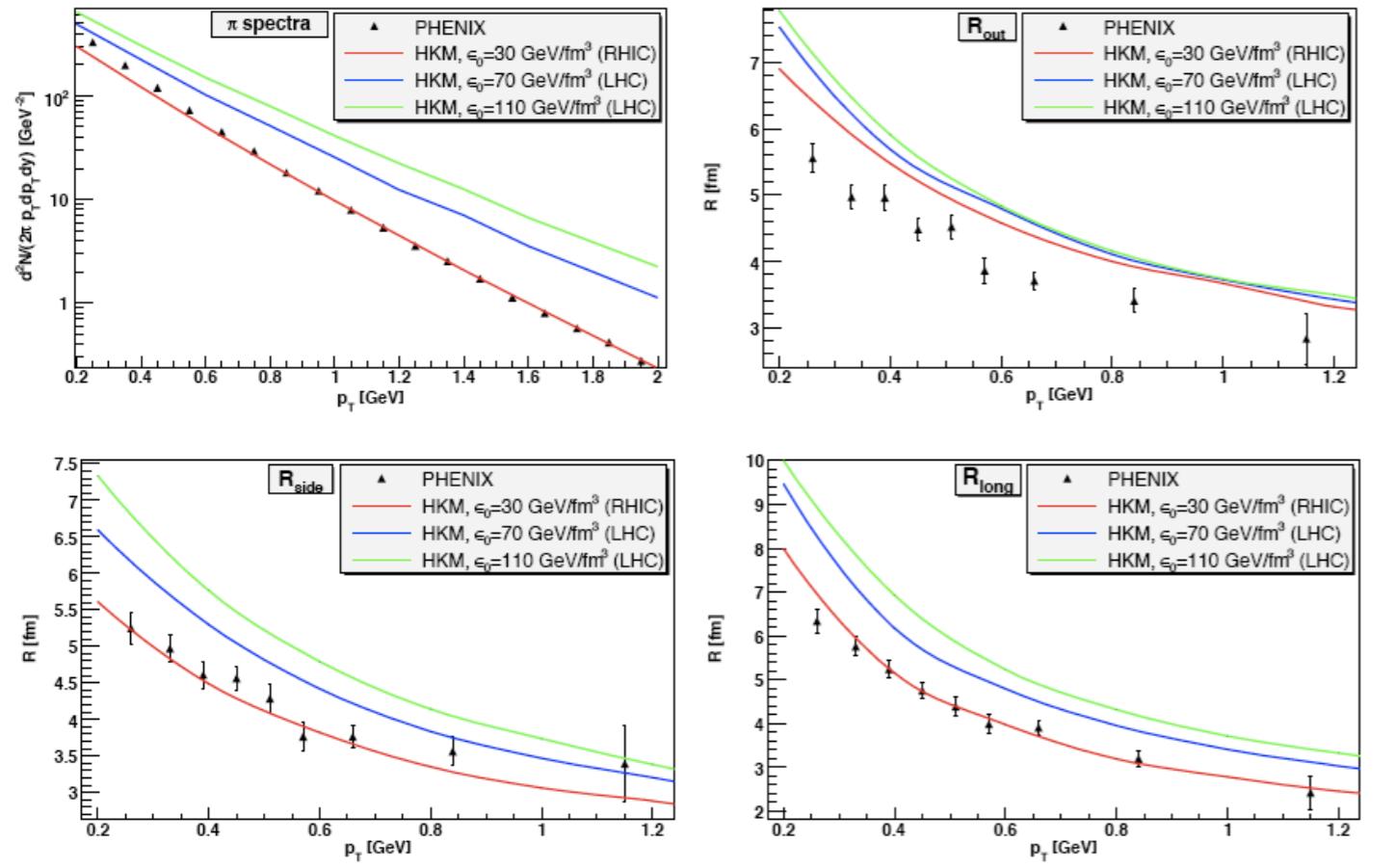
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p-pbar<4 at  $\eta=0$  (Bj: Topor Pop et al., Bopp et al.; hydro: Eskola et al.; EPOS: Porteboeuf et al., RDM: Wolschin et al.).

- ⇒ **Cronin effect for protons will strongly constrain models.**
- ⇒ **Ratios will further clarify the hadronization mechanism.**
- ⇒ **A large b-bbar at  $\eta=0$  would be a real surprise.**



# A.4. Correlations at low p<sub>T</sub>: HBT

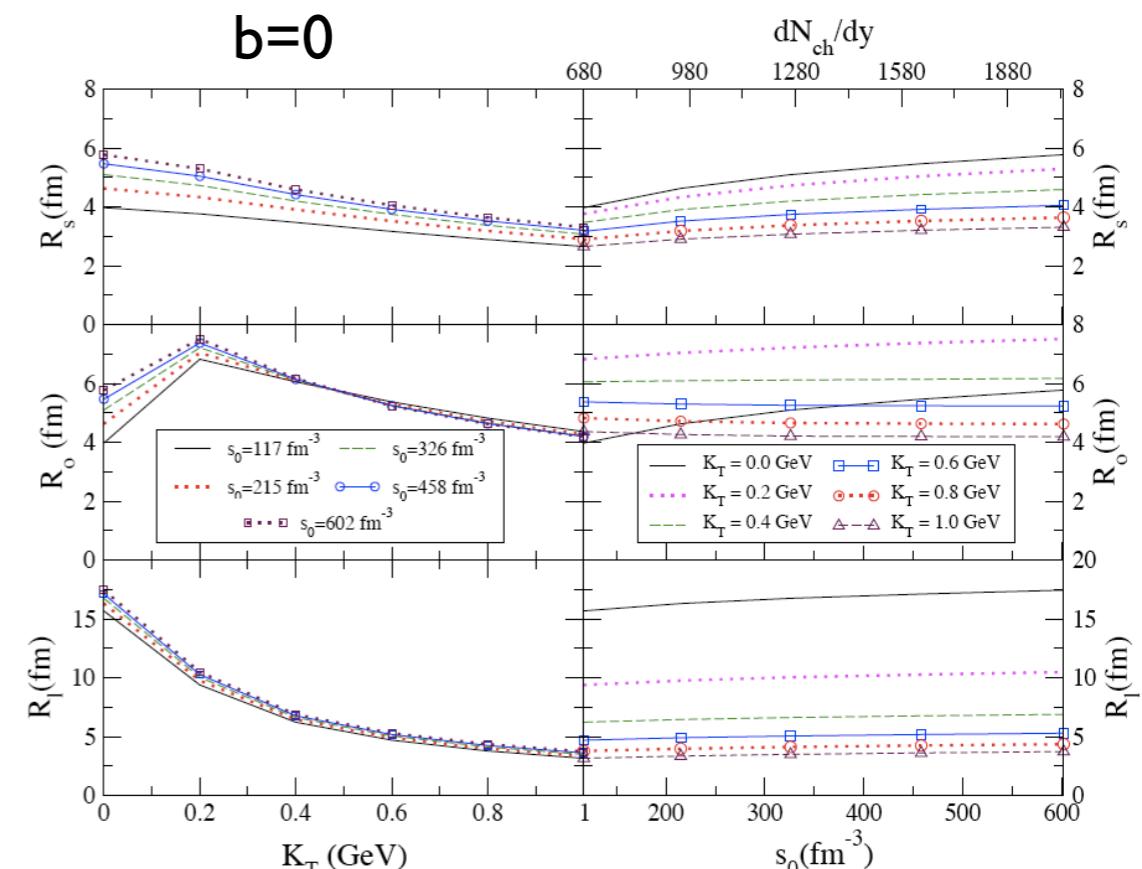


Sinyukov et al., HKM; also  
Karpenko et al., FASTMC.

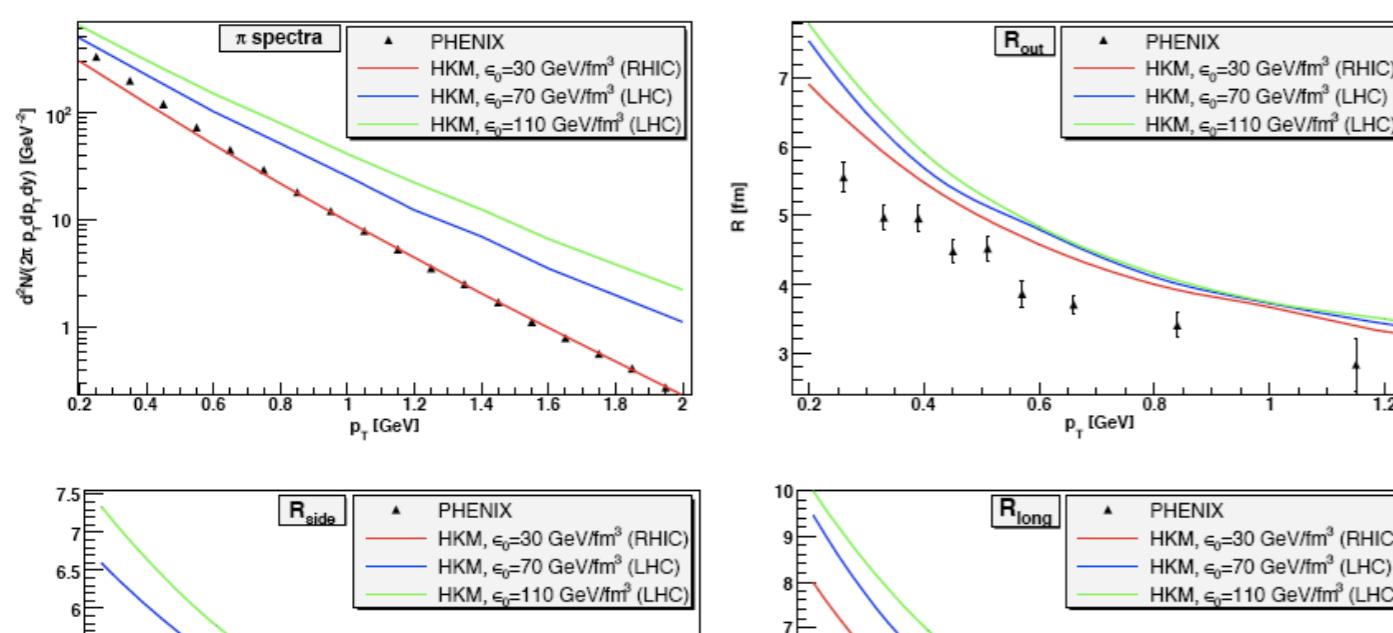
RHIC/LHC for <math>\pi</math>'s	Chen et al., AMPT, $b=0$ , $0.3 < k_T < 1.5$	Chojnacki et al., 0712.0947, hydro+stat., $b=1$ fm, $k_T=0.3$ GeV
$R_{out}$	3.60/4.23	5.4/6.0-6.5
$R_{side}$	3.52/4.70	4.3/5.3-6.3
$R_{long}$	3.23/4.86	6.1/7.6-8.6

Hydro: same problems as at RHIC -  $R_{out}(k_T)$ ,  $R_{side}(k_T)$ ,  $R_{out} \gg R_{side}$ ;  
out-  $\rightarrow$  in-plane shape.

Frodermann et al.



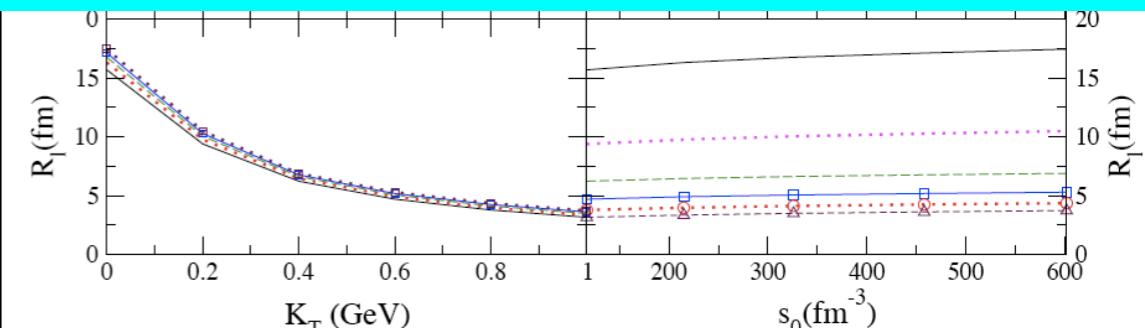
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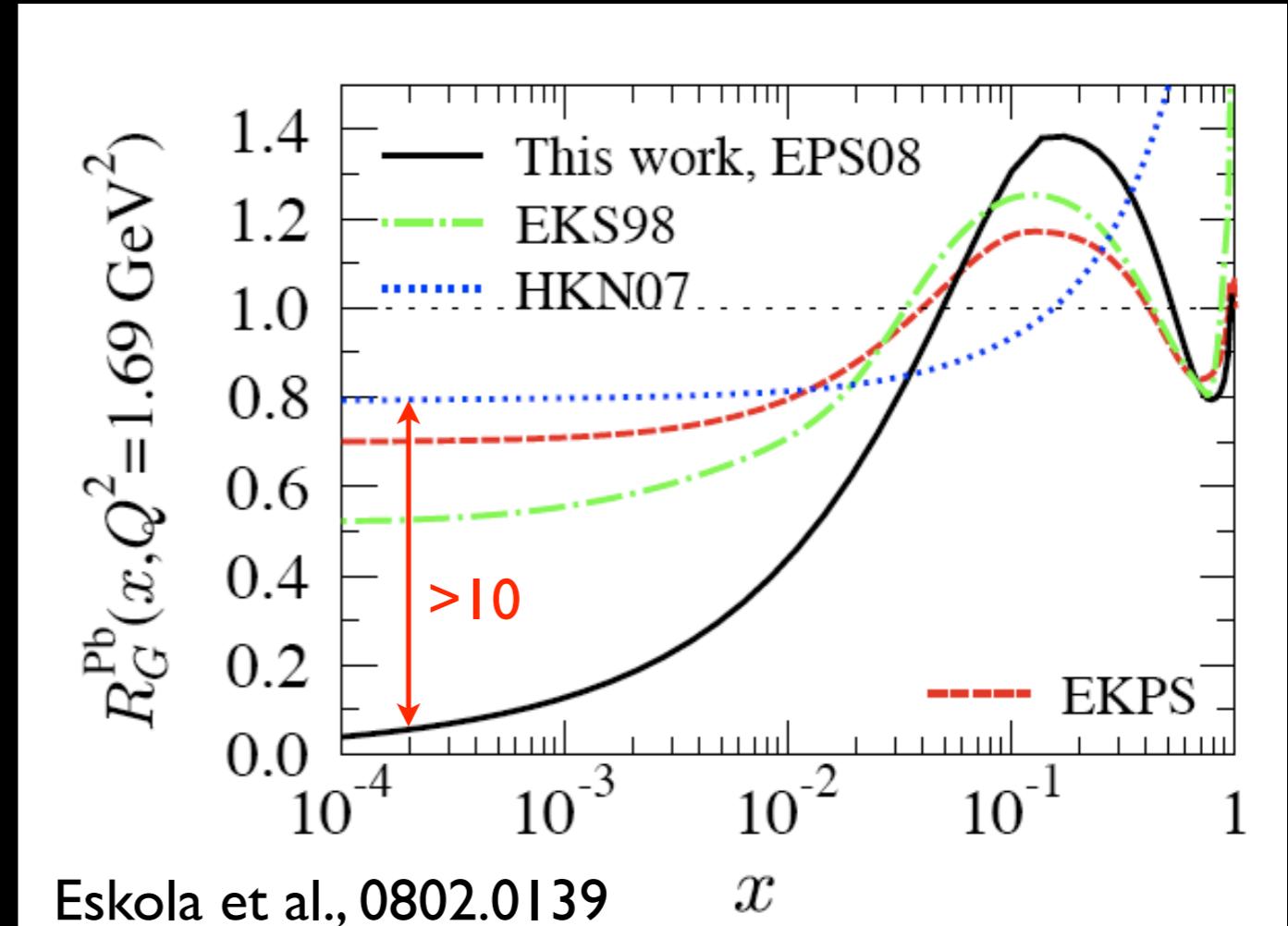
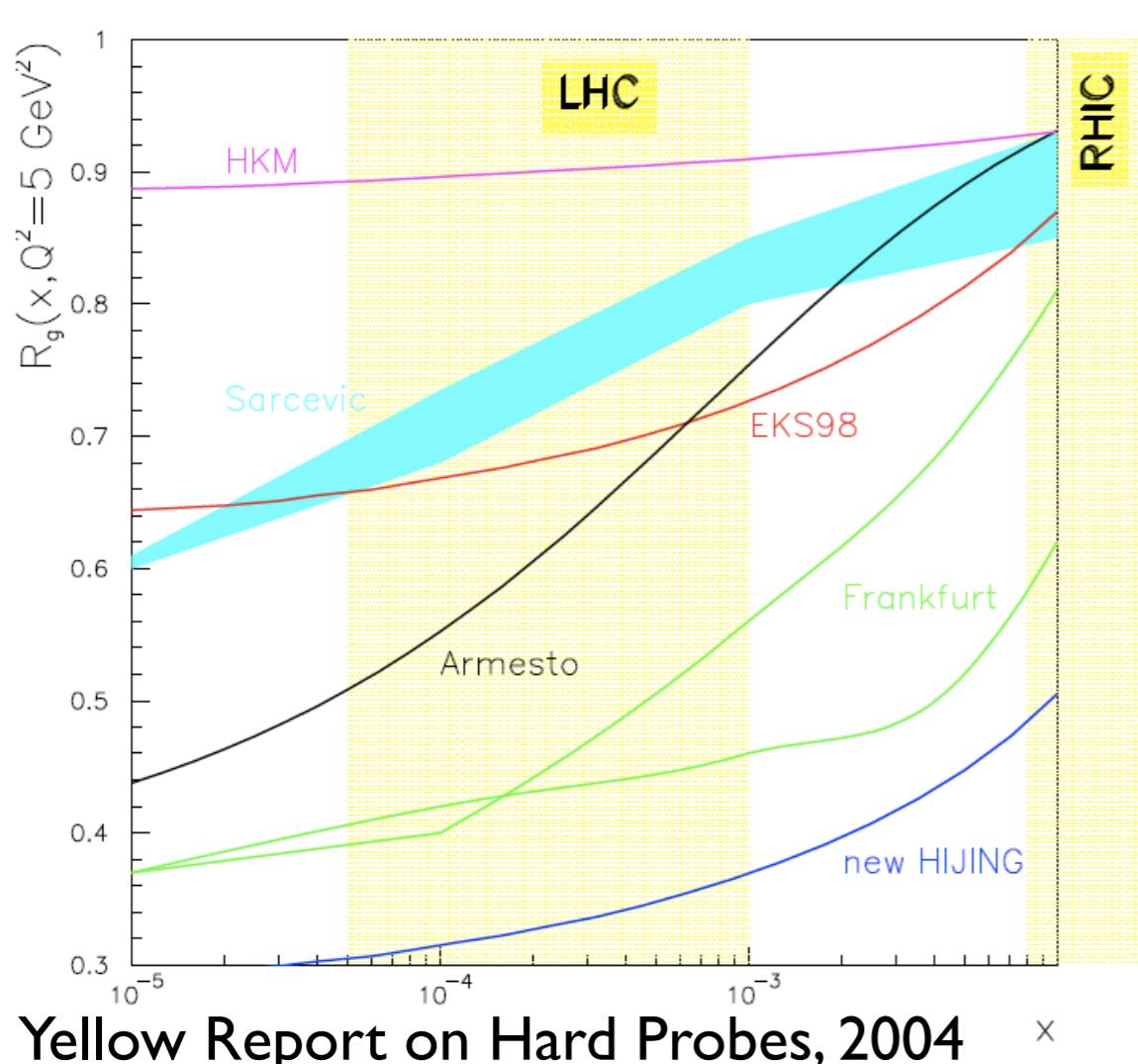
Hydro: same problems as at RHIC -  $R_{out}(k_T)$ ,  $R_{side}(k_T)$ ,  $R_{out} \gg R_{side}$ ; out-  $\rightarrow$  in-plane shape.

- ⇒ R's increase from RHIC to the LHC in all models.
- ⇒ But the predictive power is limited by the problems at RHIC.
- ⇒ Dissipative effects on HBT are not well understood yet.

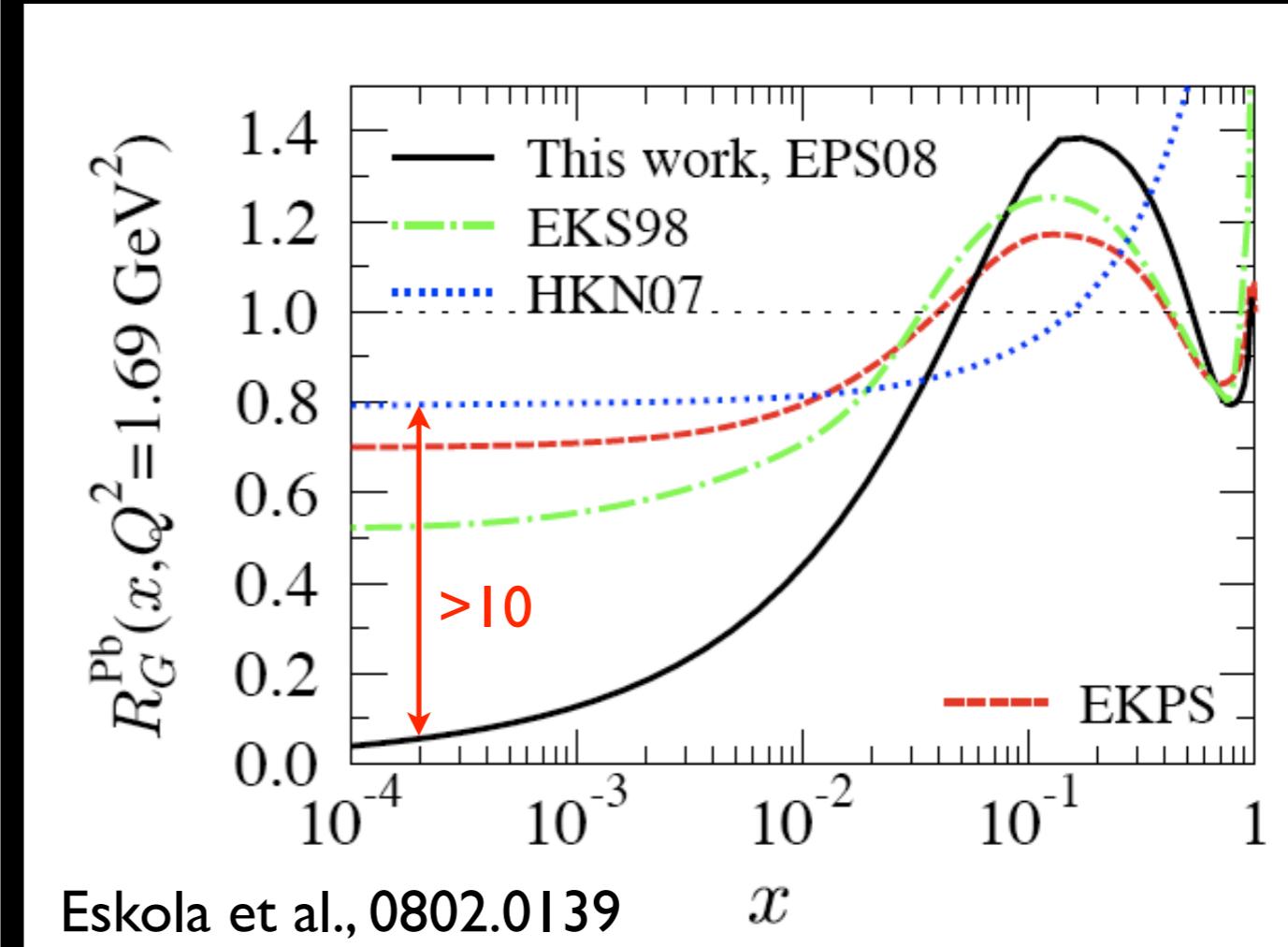
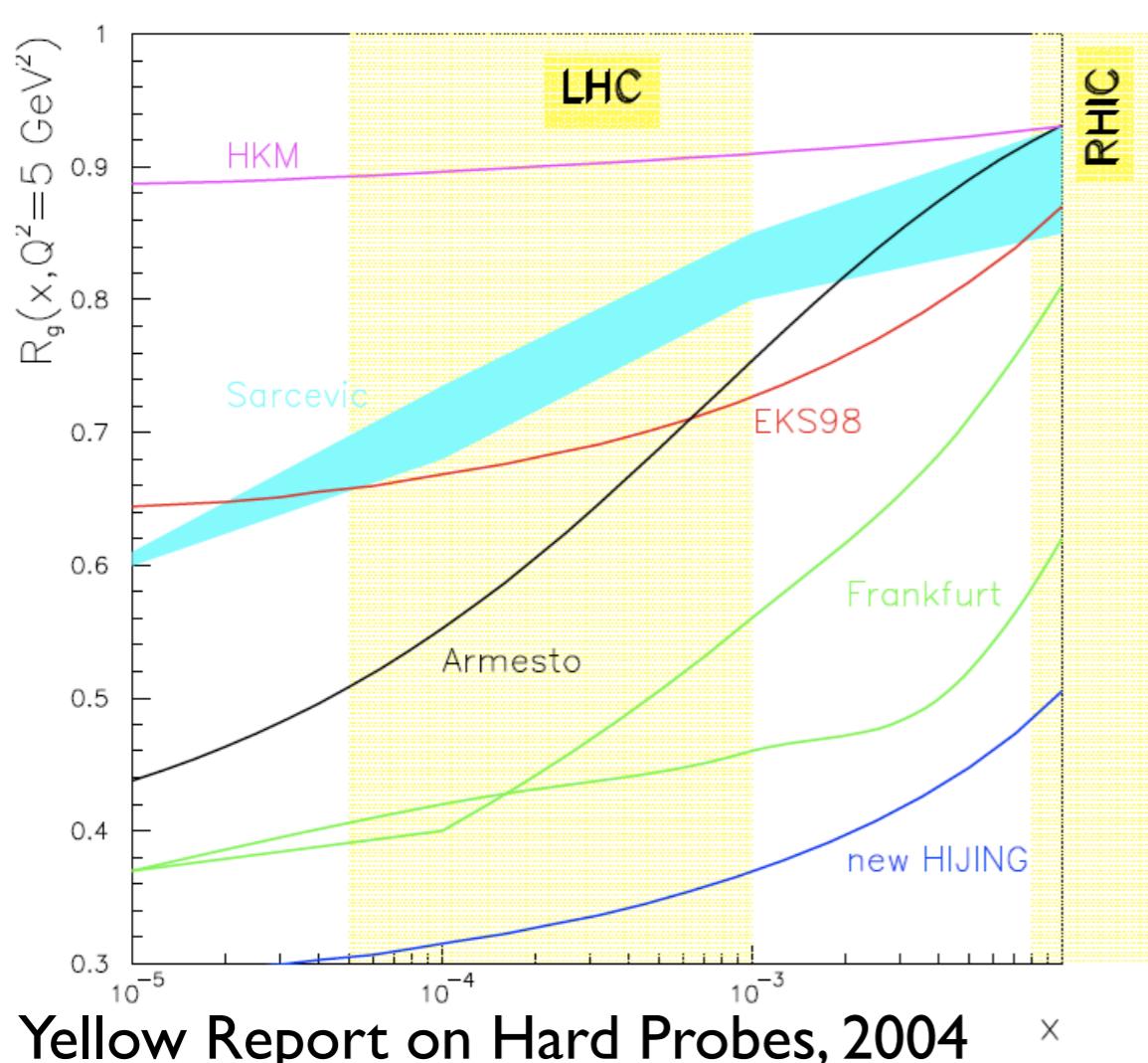
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## B. Hard probes: benchmark

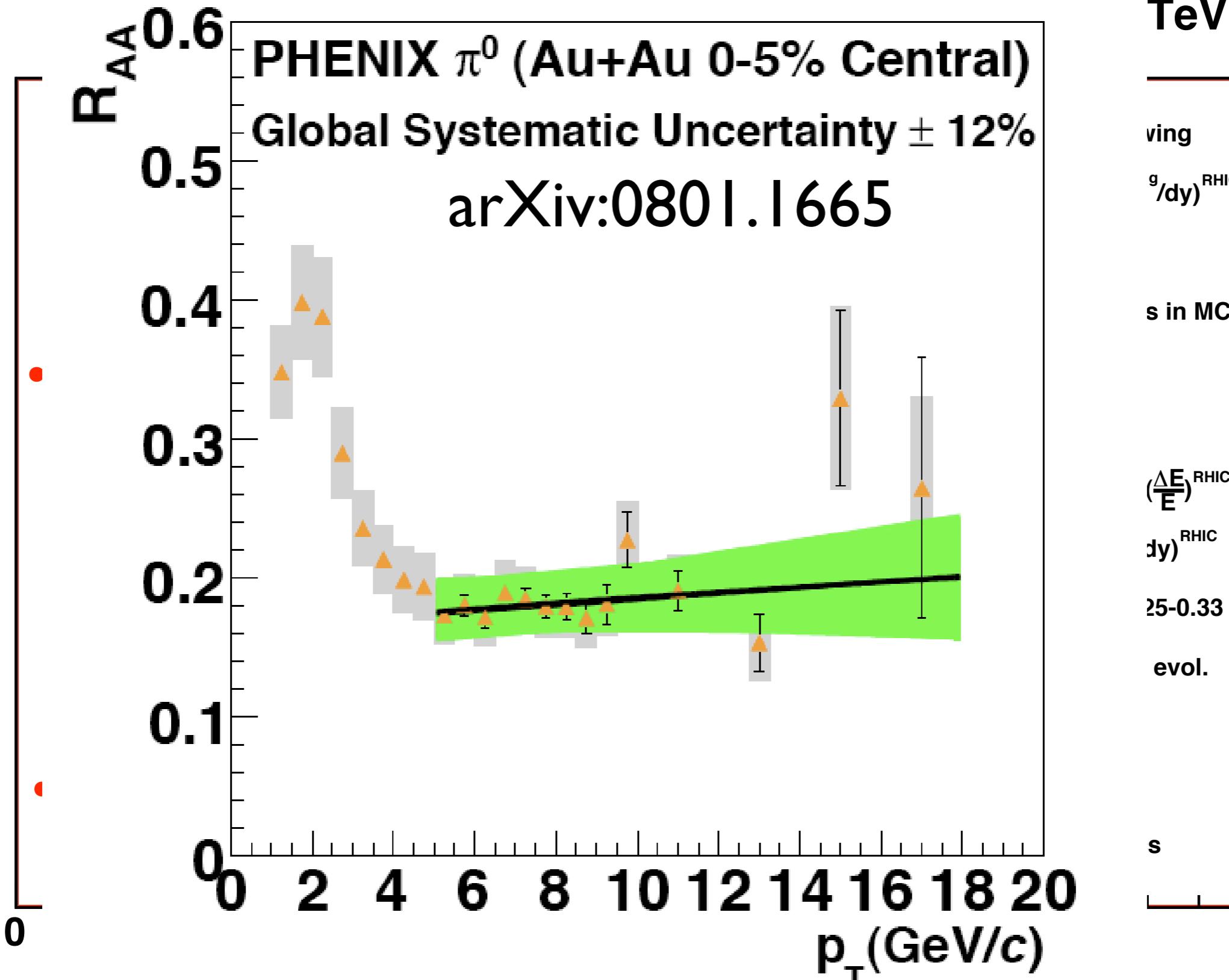


## B. Hard probes: benchmark



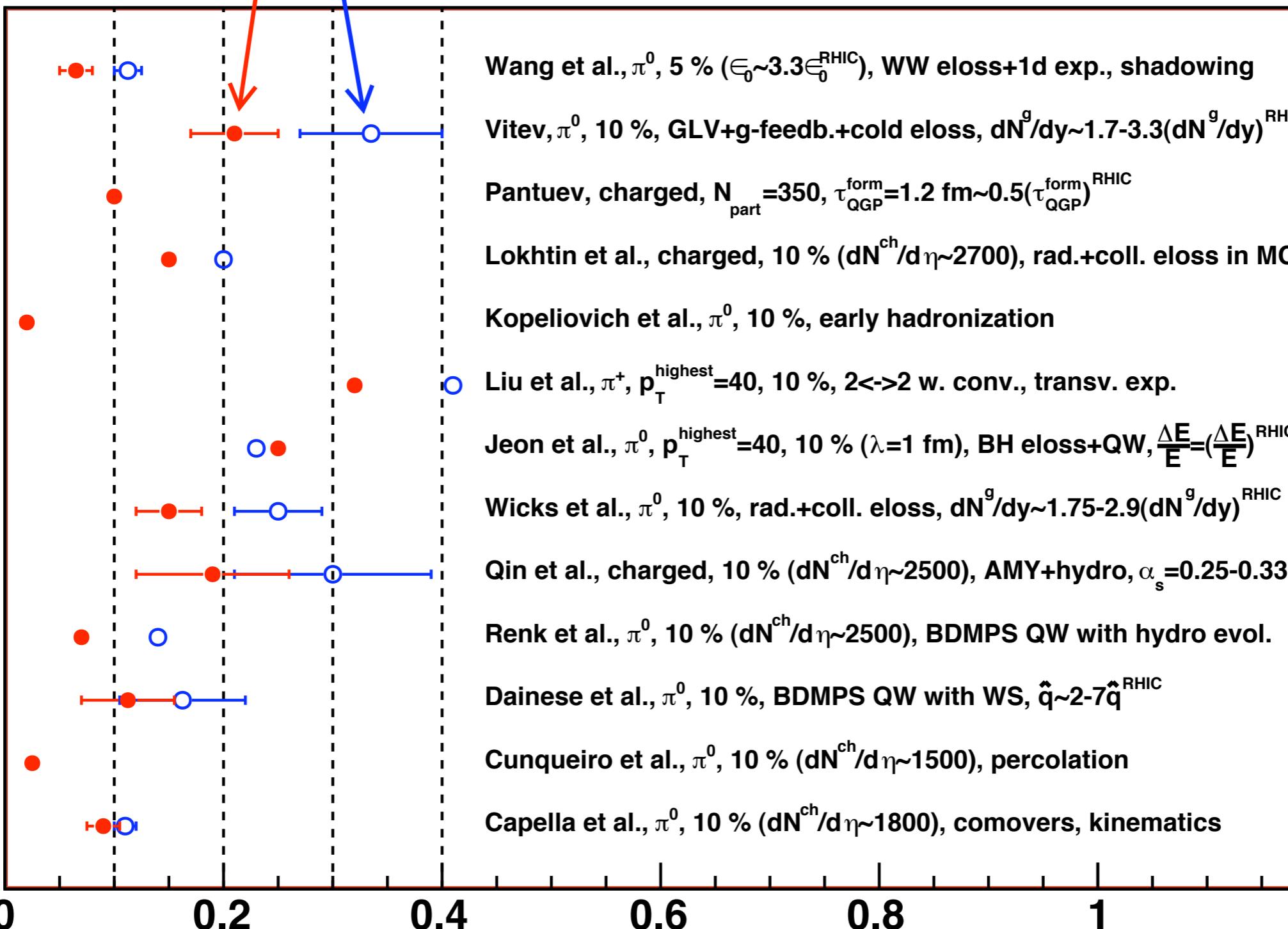
► To avoid this uncertainty within the LHC frame, an accurate control of the benchmark will be required (as it was at RHIC).

# B.I. High-p<sub>T</sub> observ.: R<sub>AA</sub> for light flavors



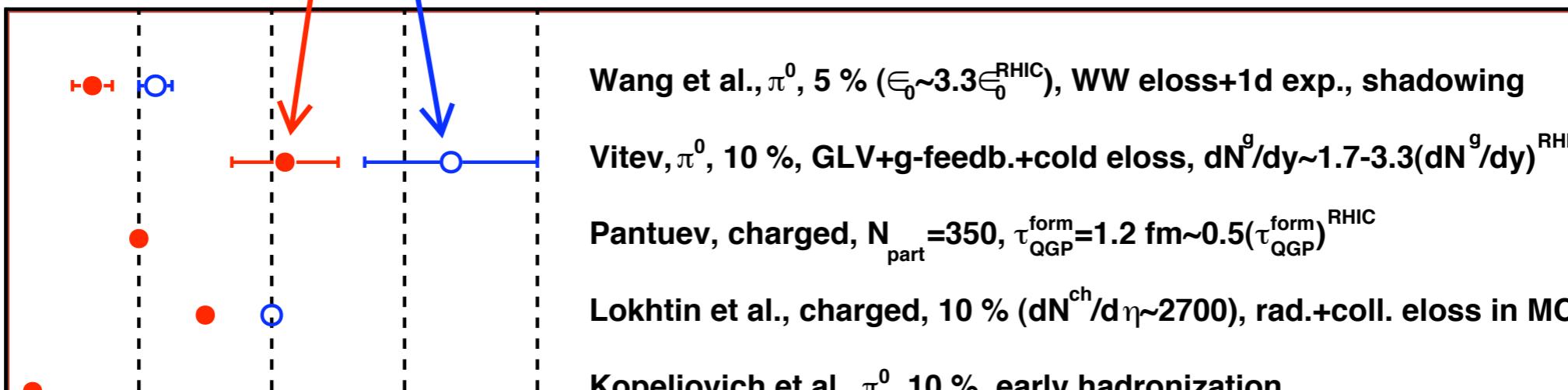
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$R_{\text{PbPb}}(p_T=20,50 \text{ GeV}, \eta=0)$  in central Pb+Pb at  $\sqrt{s_{\text{NN}}}=5.5 \text{ TeV}$

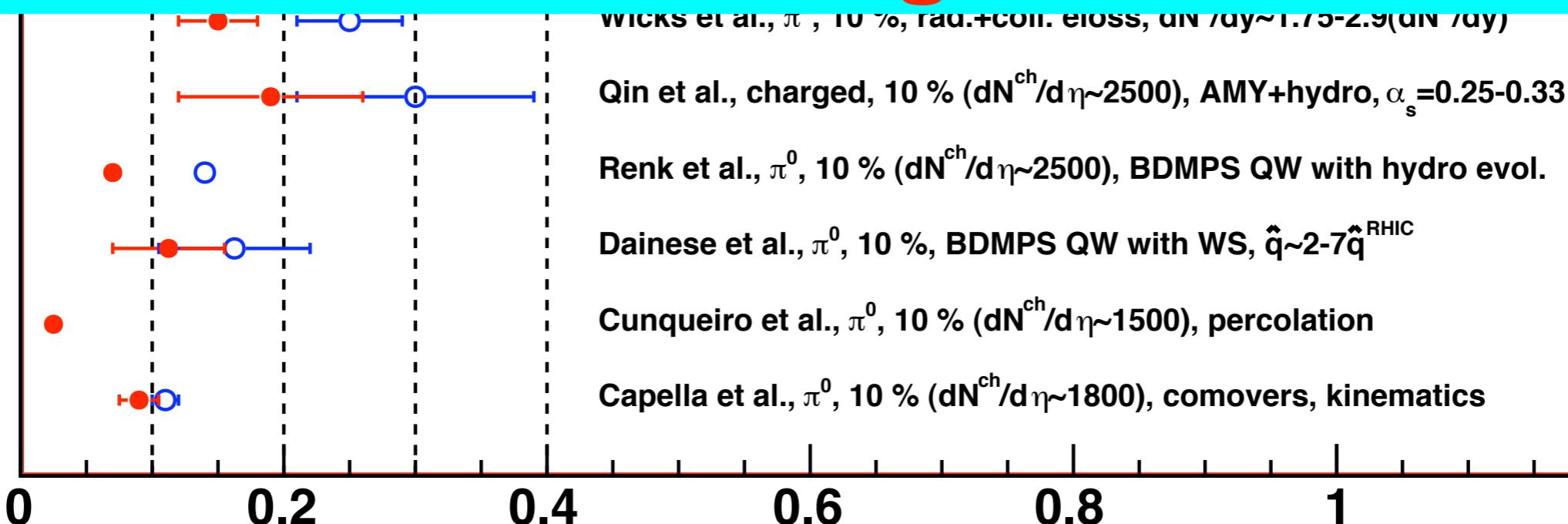


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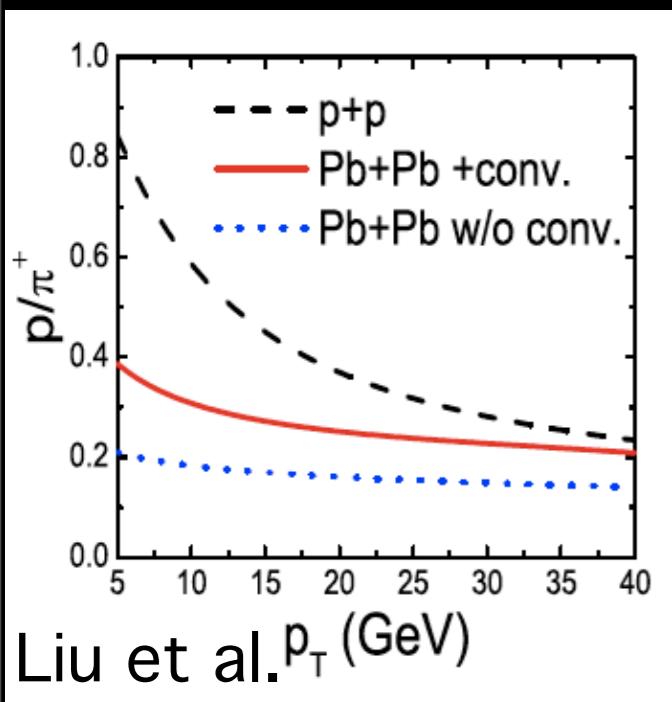
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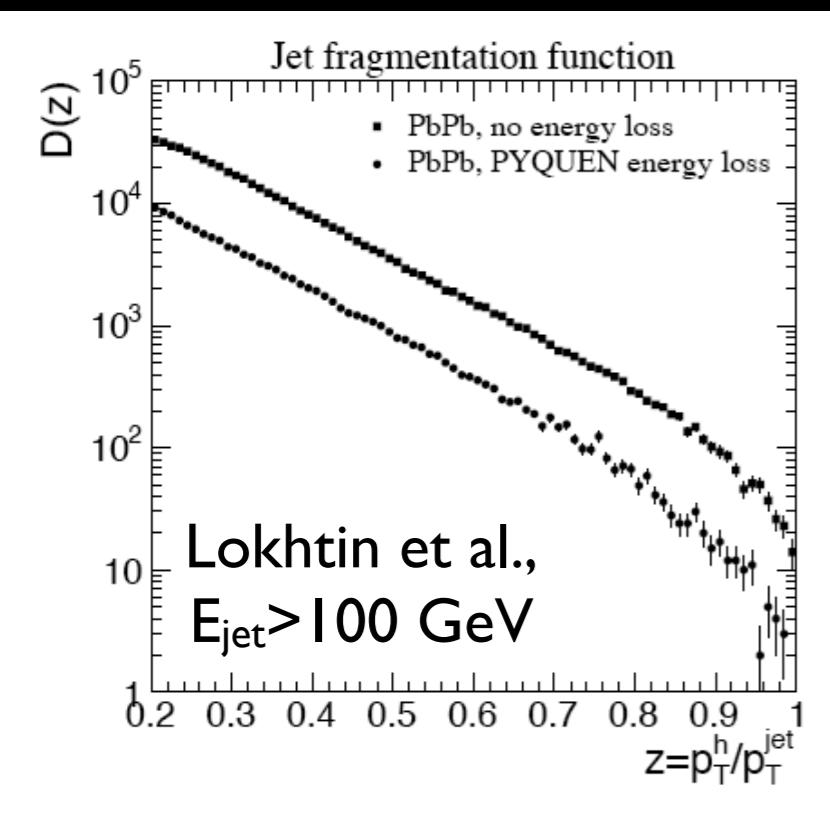
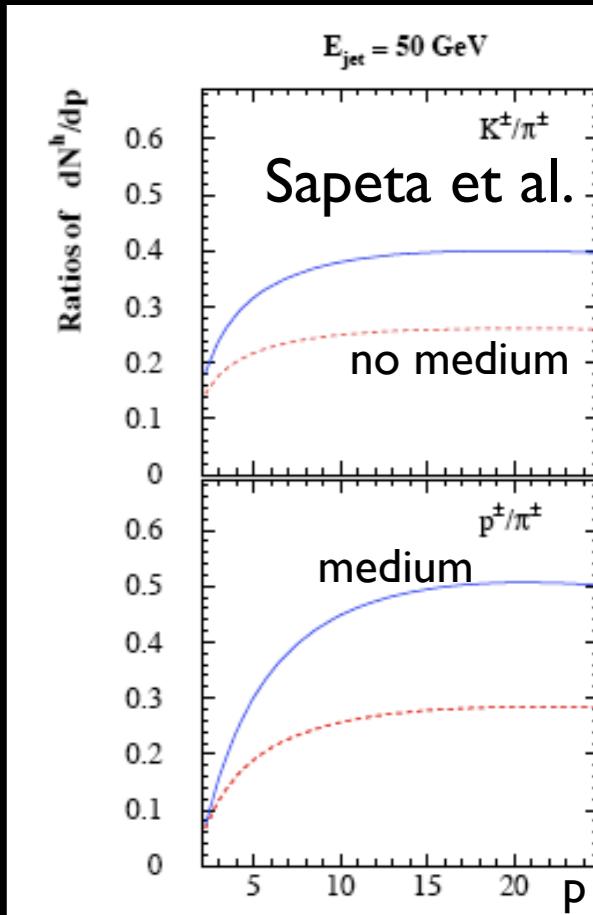
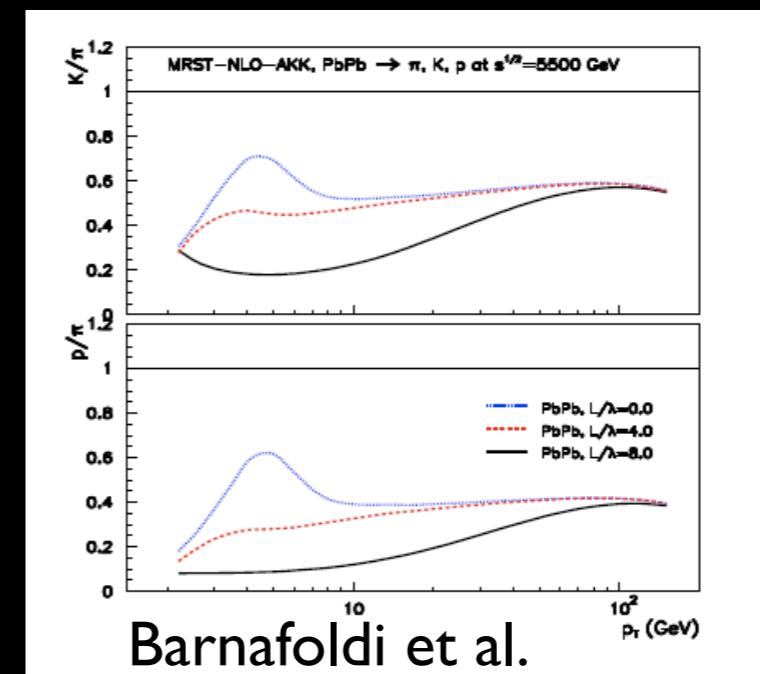
► Radiative energy loss favors  $R_{\text{AA}} \sim 0.1-0.2$  at  $p_T \sim 20 \text{ GeV}$  and increasing with increasing  $p_T$ .



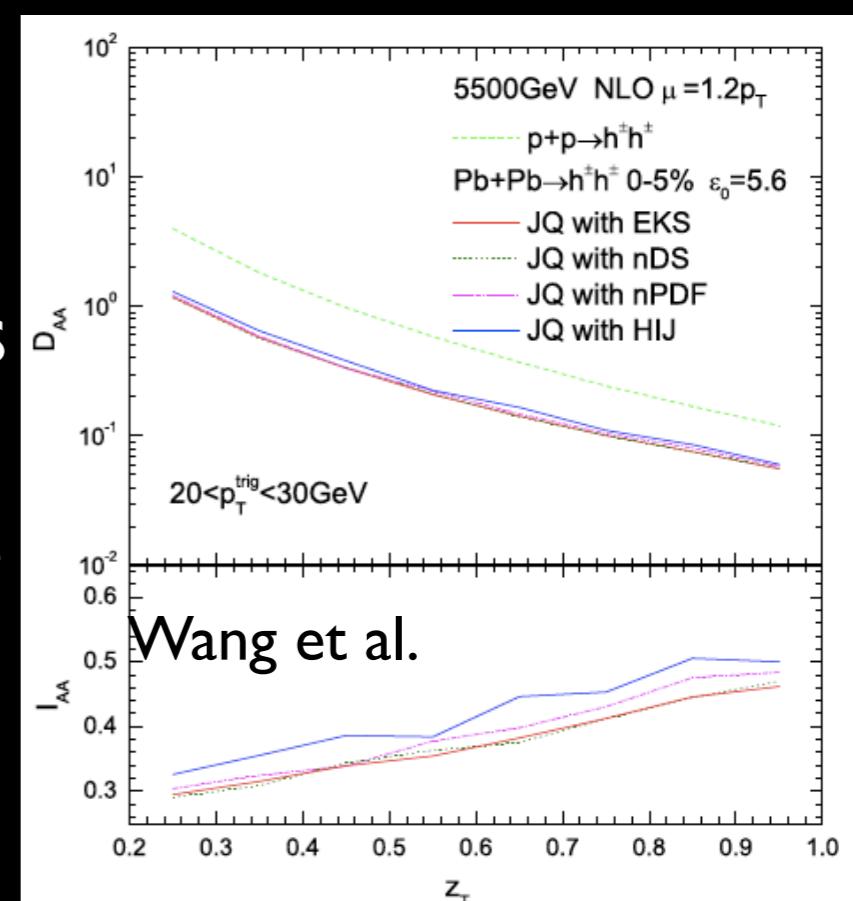
# B.I. High- $p_T$ observ.: hadrochemistry and FF



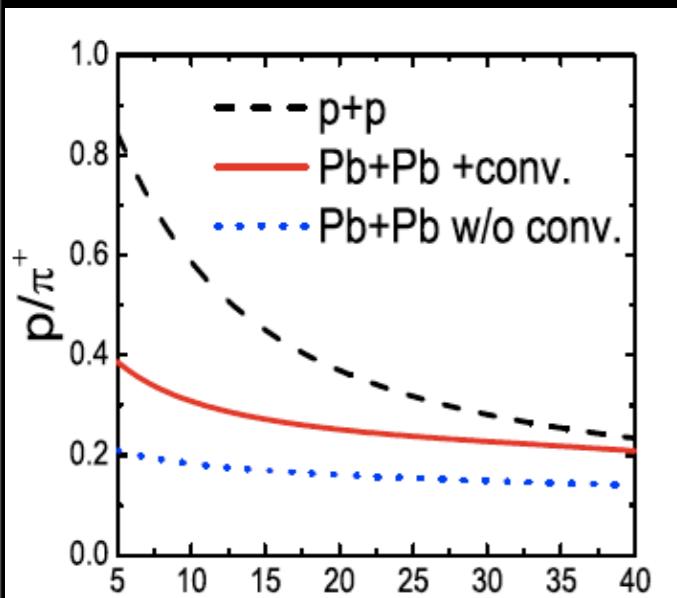
**Modification of hadrochemistry due to elastic +conversions, rad. eloss or modified jet radiation.**



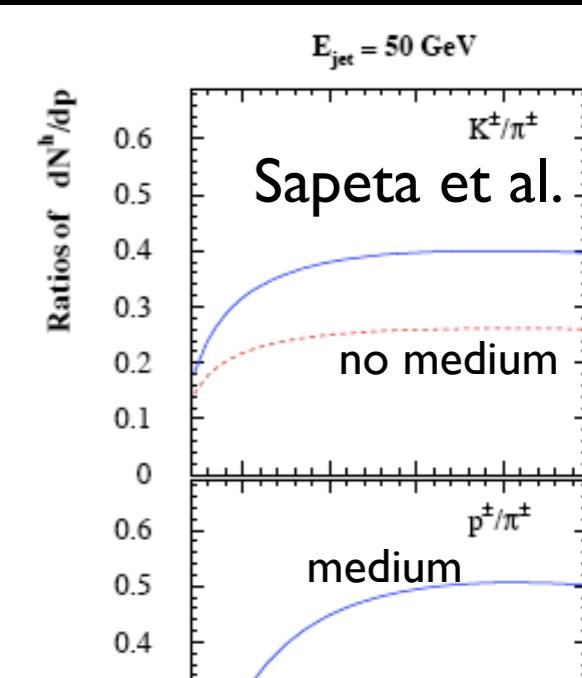
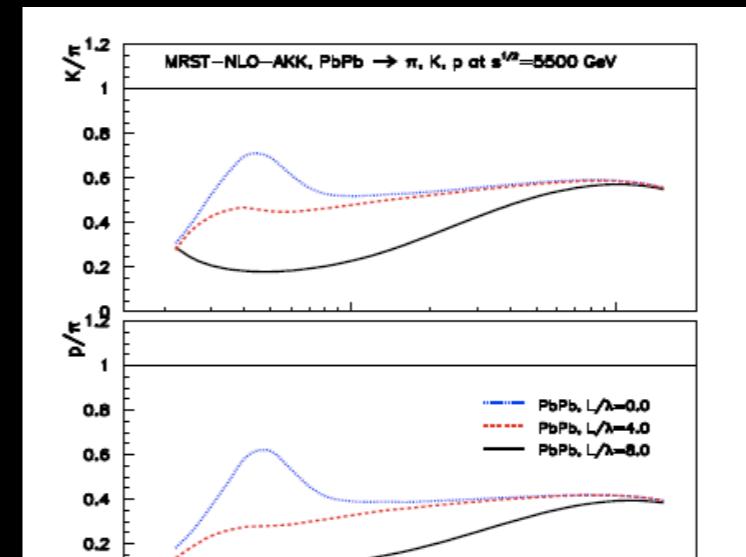
**Modified fragmentation functions, both for jets (elastic+radiative in PYQUEN) and for the hadron-triggered case (WW rad. model).**



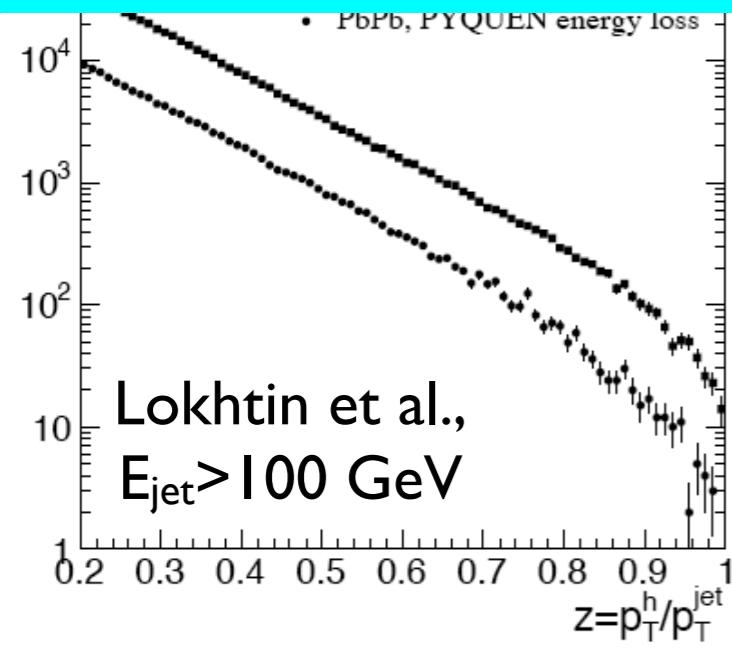
# B.I. High- $p_T$ observ.: hadrochemistry and FF



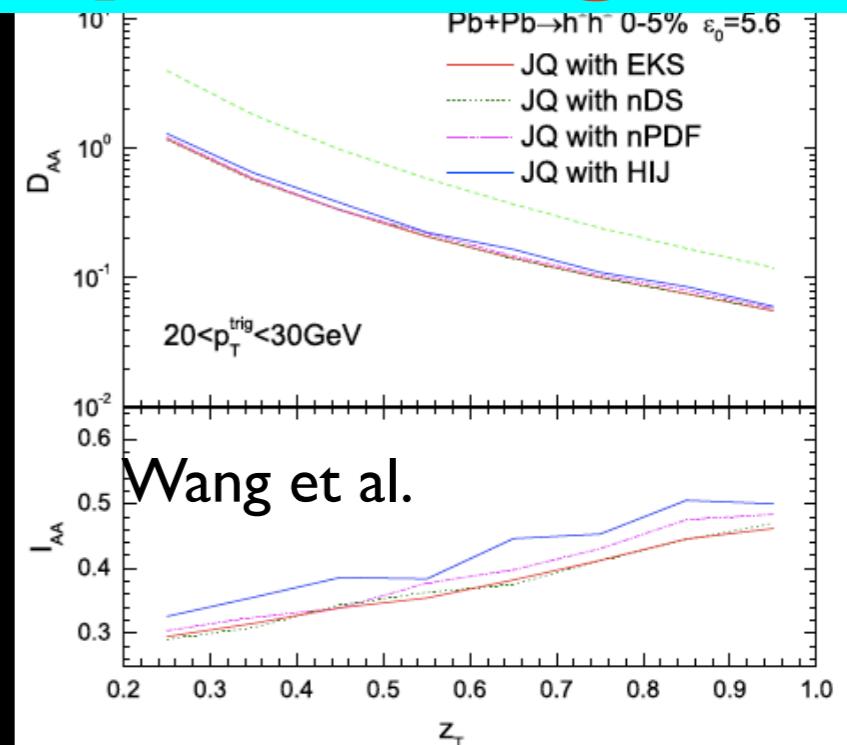
**Modification of hadrochemistry due to elastic +conversions, rad. eloss or modified jet**



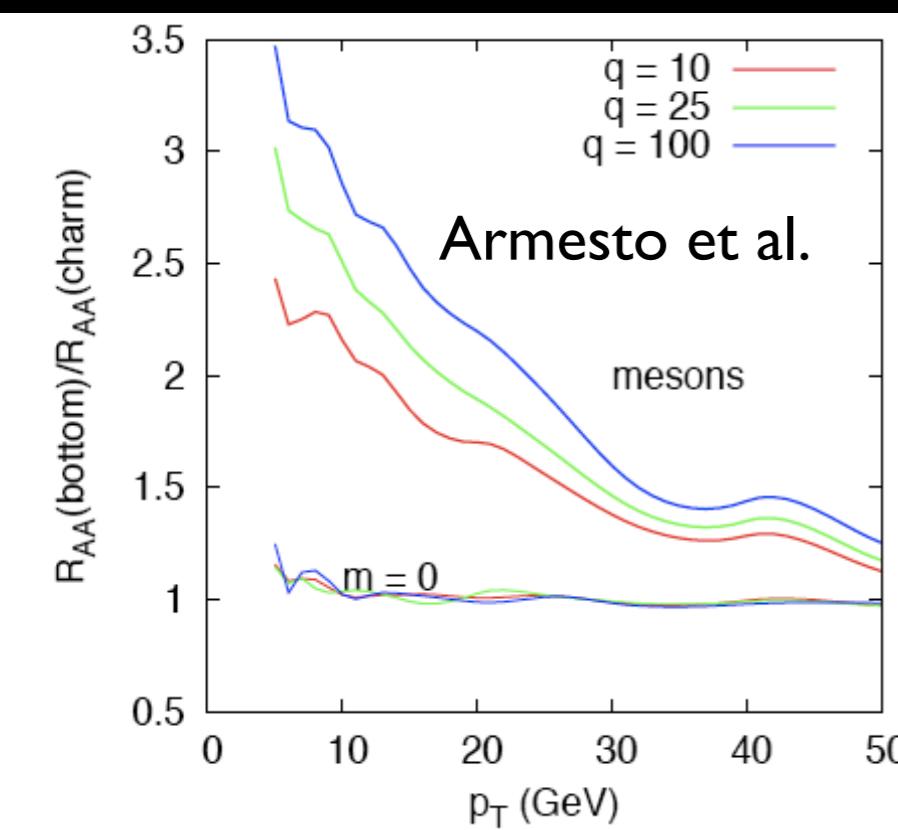
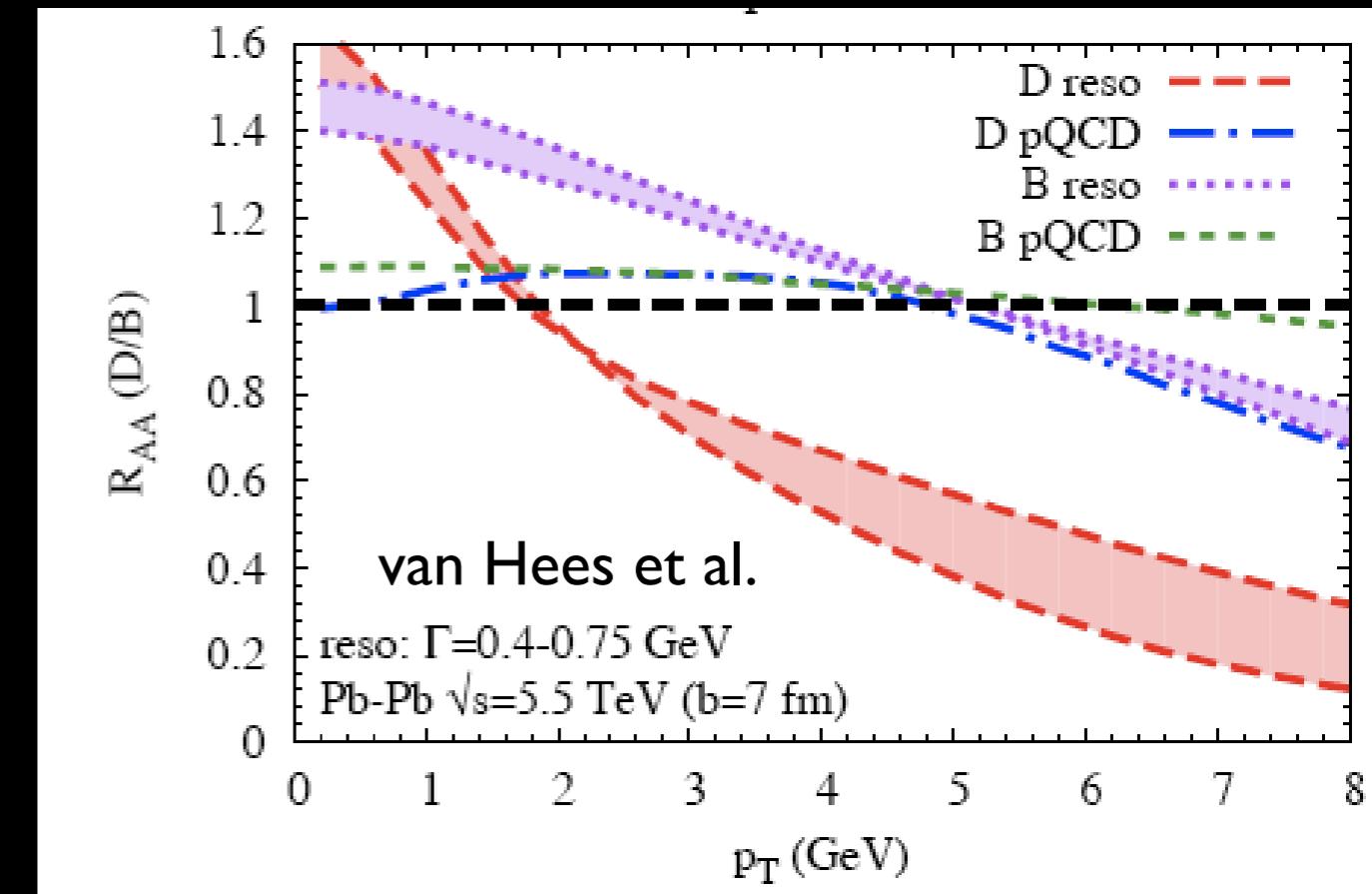
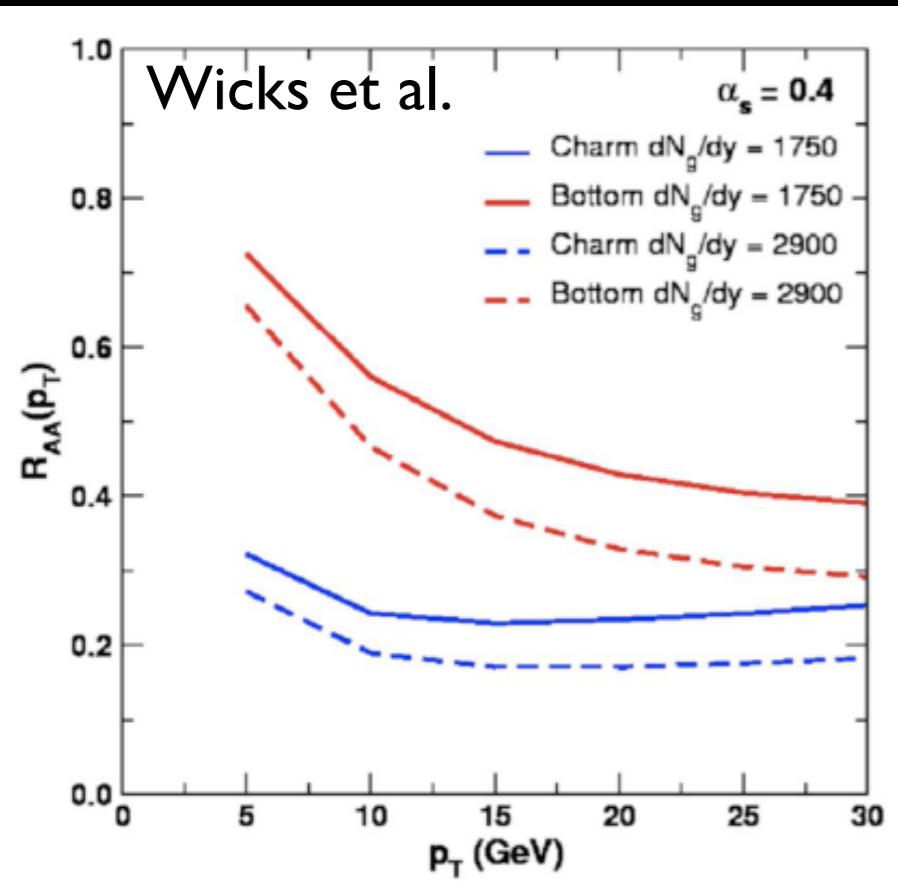
► **Chemical composition and more differential observables will be key to establish the mechanism underlying jet quenching.**



**fragmentation functions, both for jets (elastic+radiative in PYQUEN) and for the hadron-triggered case (WW rad. model).**

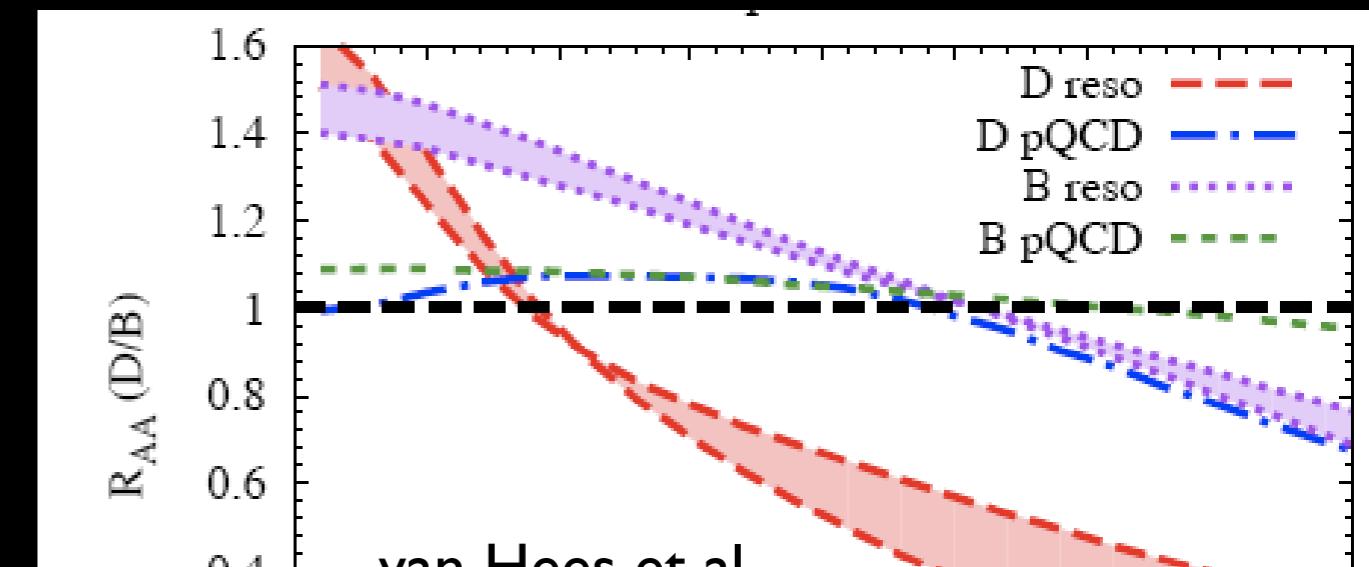
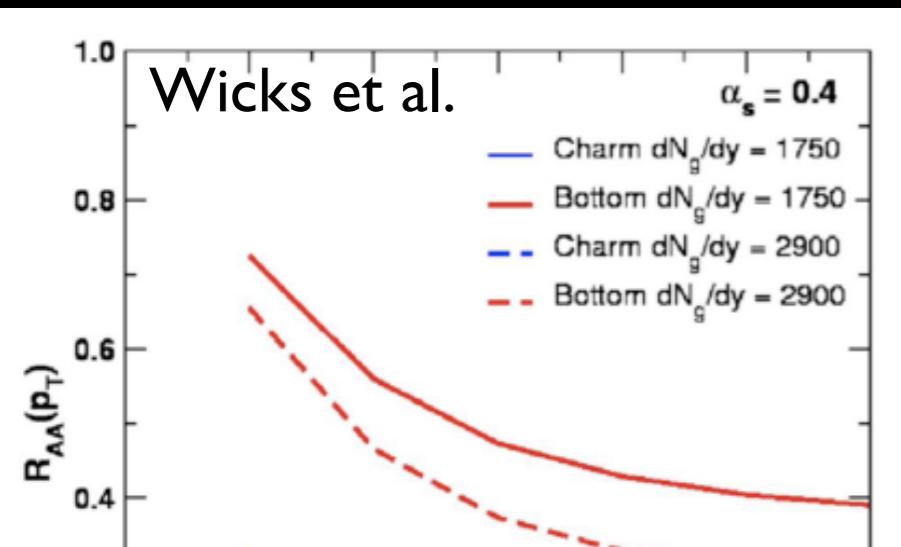


## B.2. HQ and quarkonium: R<sub>AA</sub> for heavy

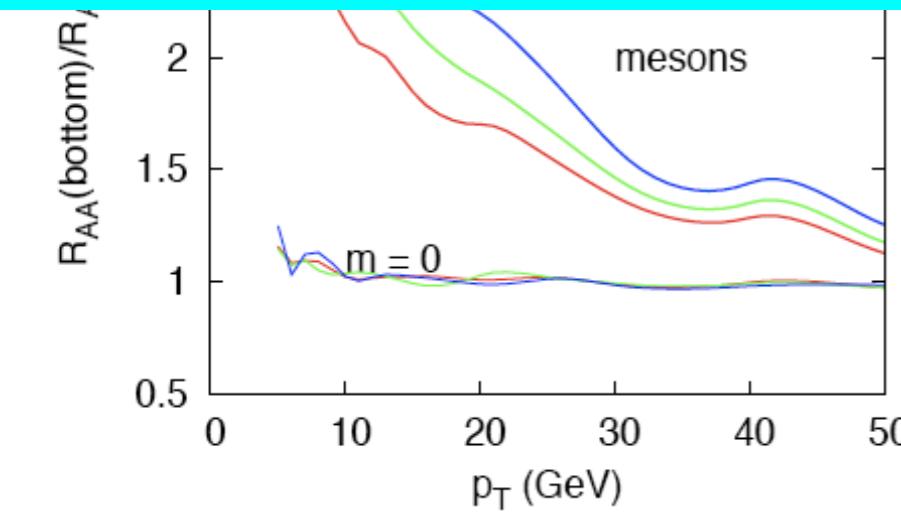


pQCD-based models valid for hadronization outside the medium, so for high  $p_T$  (see also Vitev). Double ratios B/D are sensitive to mass effects until quite high  $p_T$  and offer possibilities to discriminate models for HQ jet quenching (Horowitz et al.).

## B.2. HQ and quarkonium: R<sub>AA</sub> for heavy

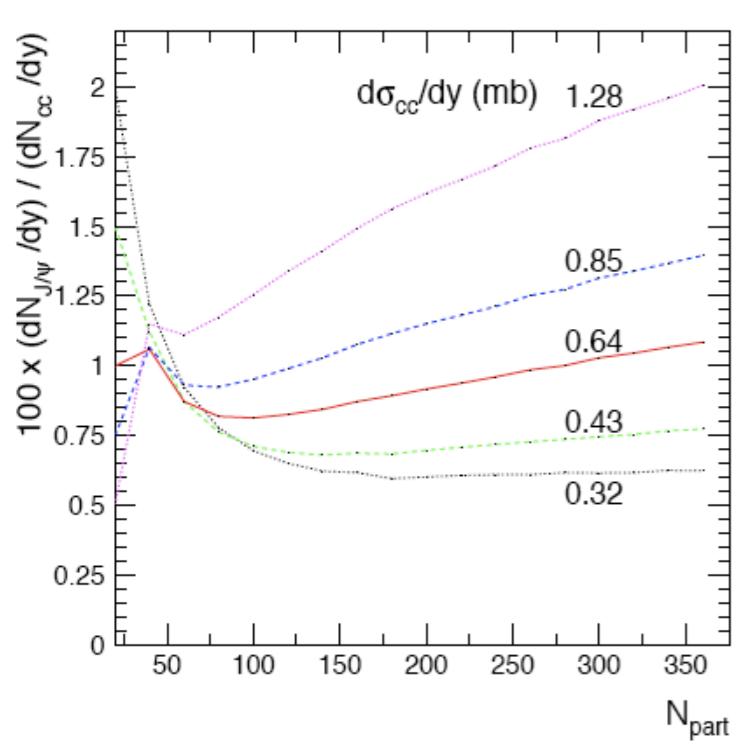


- LHC will provide  $R_{AA}$  for both leptons and mesons.
  - Double ratios will become available.
  - This will clarify the mechanism for HQ jet quenching and hadronization.

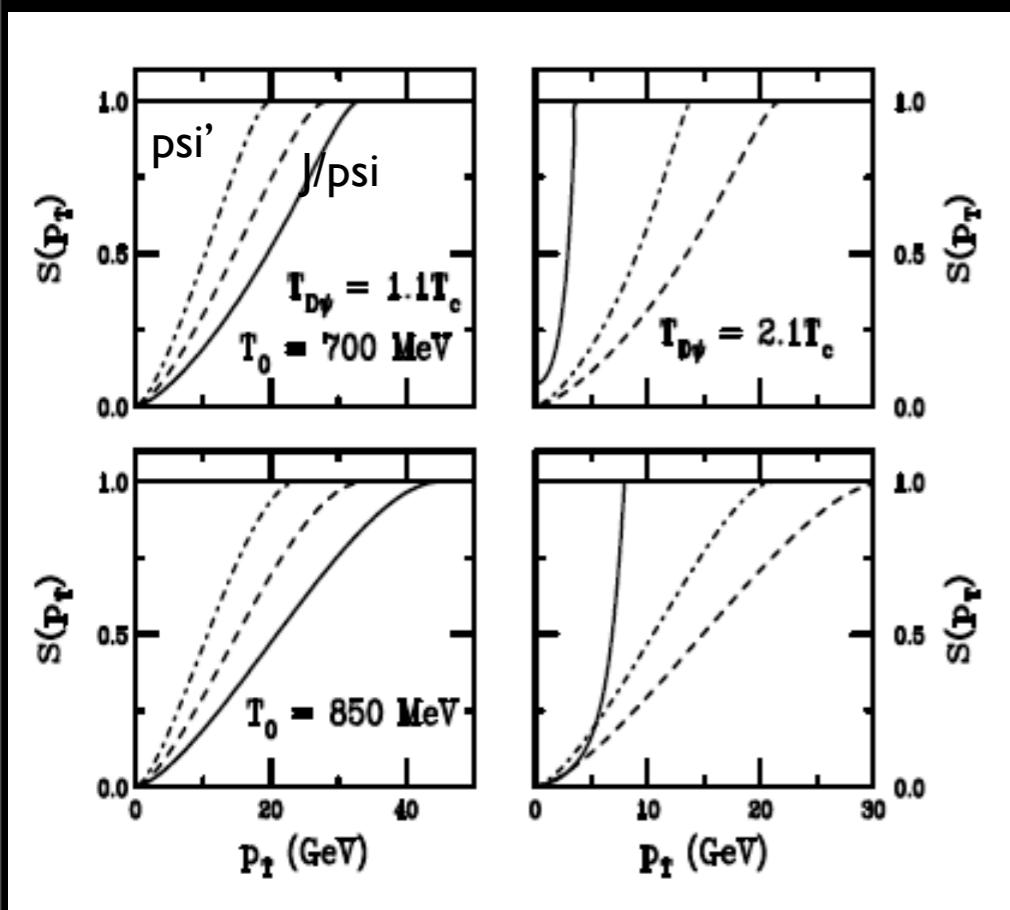
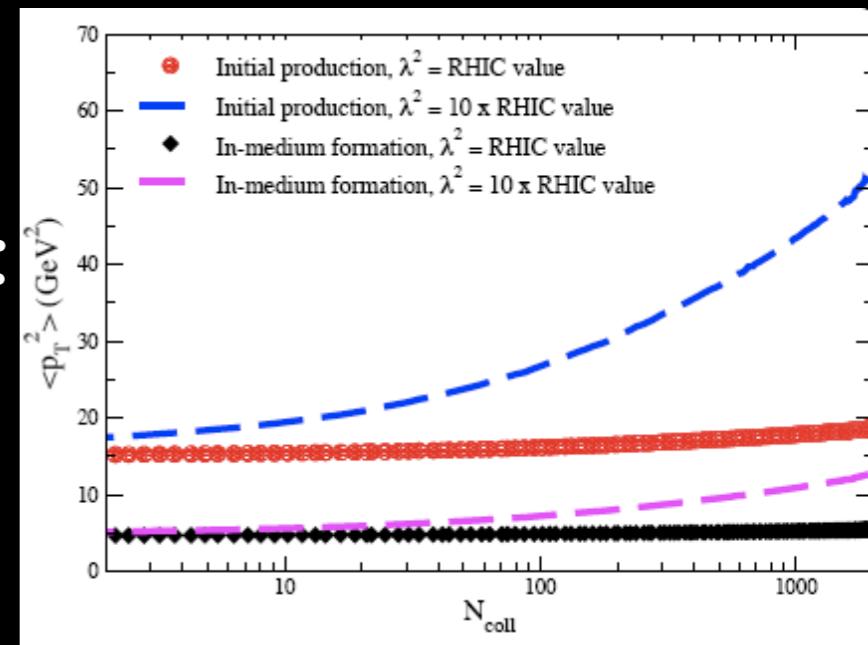


Double ratios B/D are sensitive to mass effects until quite high  $p_T$  and offer possibilities to discriminate models for HQ jet quenching (Horowitz et al.).

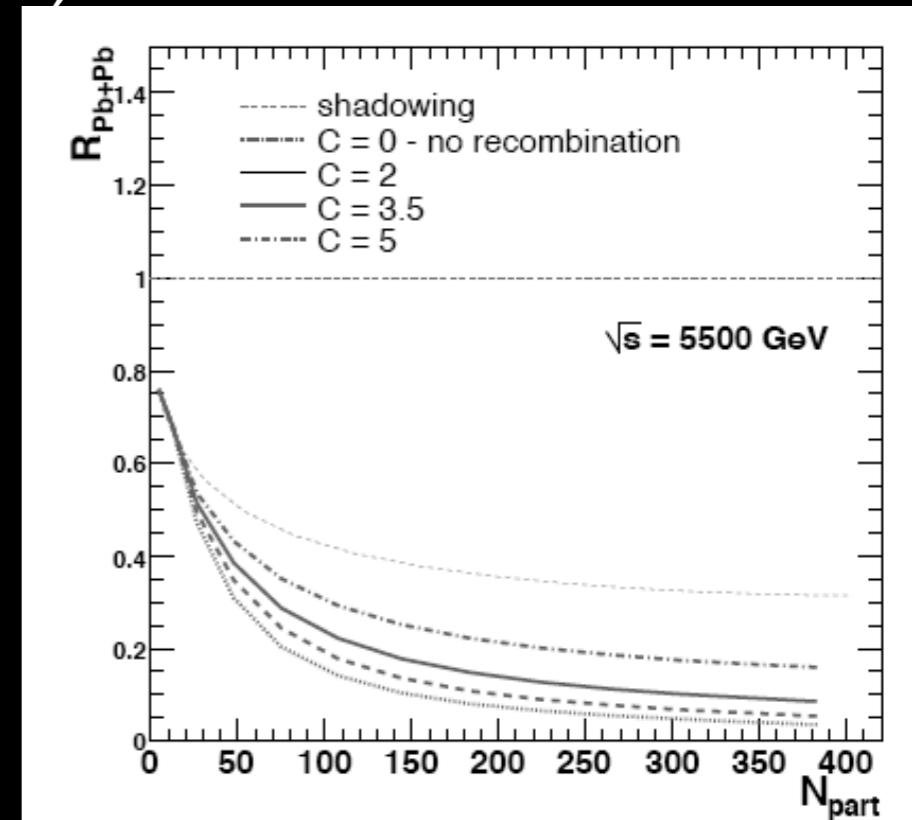
## B.2. HQ and quarkonium: quarkonia suppr.



Andronic et al.: dependence on charm cross section; Thews et al.:  $p_T$ -broadening to verify recombination, with uncertainties from cold matter effects (Kang et al.).

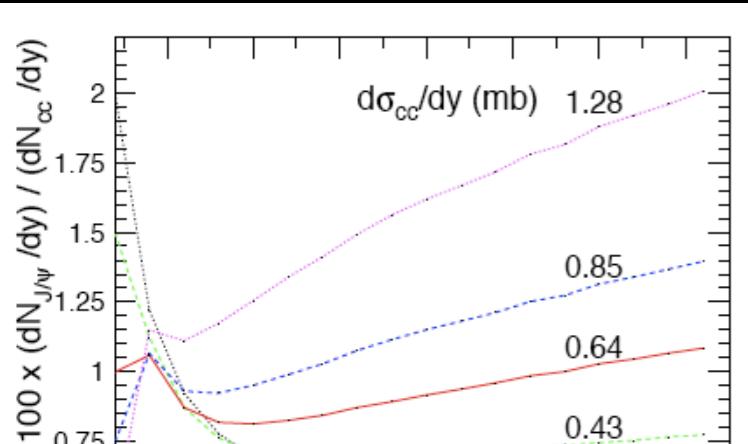


Vogt:  $p_T$ -dependent screening, no regeneration;  
H. Liu et al., suppression at larger  $p_T$ .

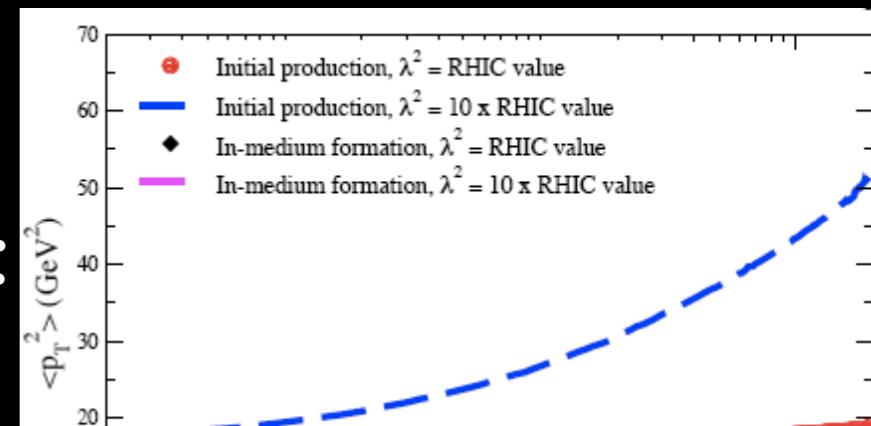


Capella et al., 0712.4331,  
comovers+reco.

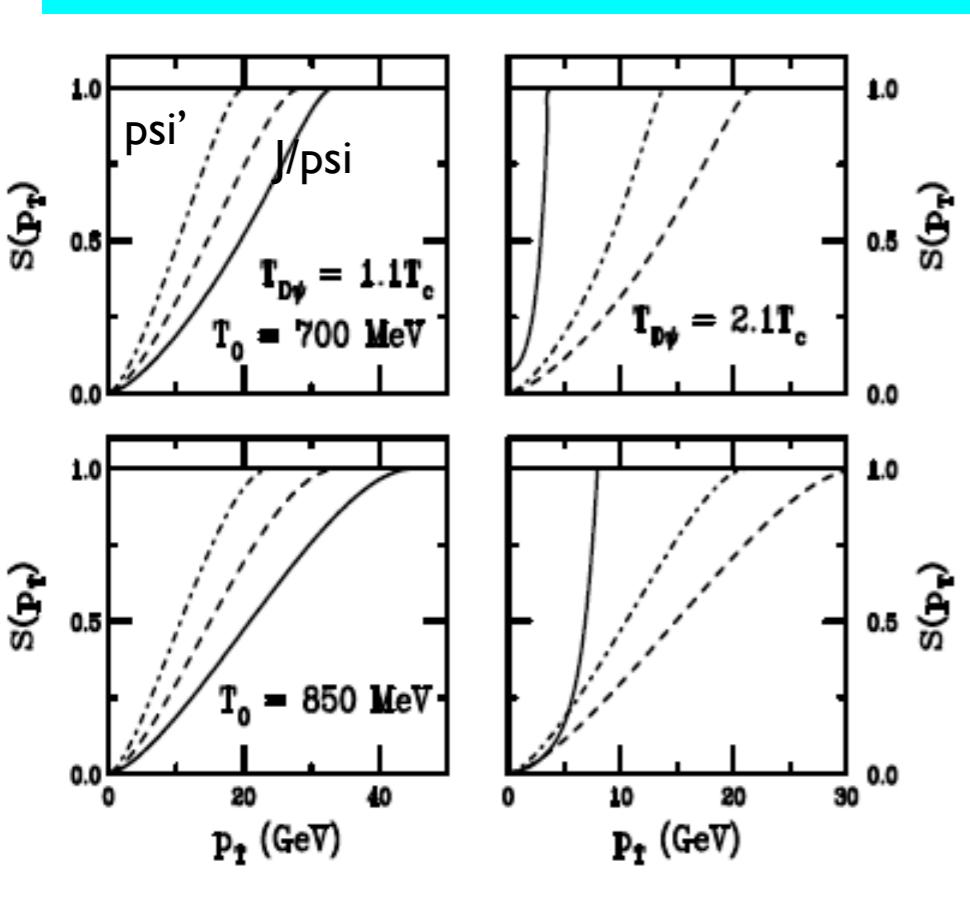
## B.2. HQ and quarkonium: quarkonia suppr.



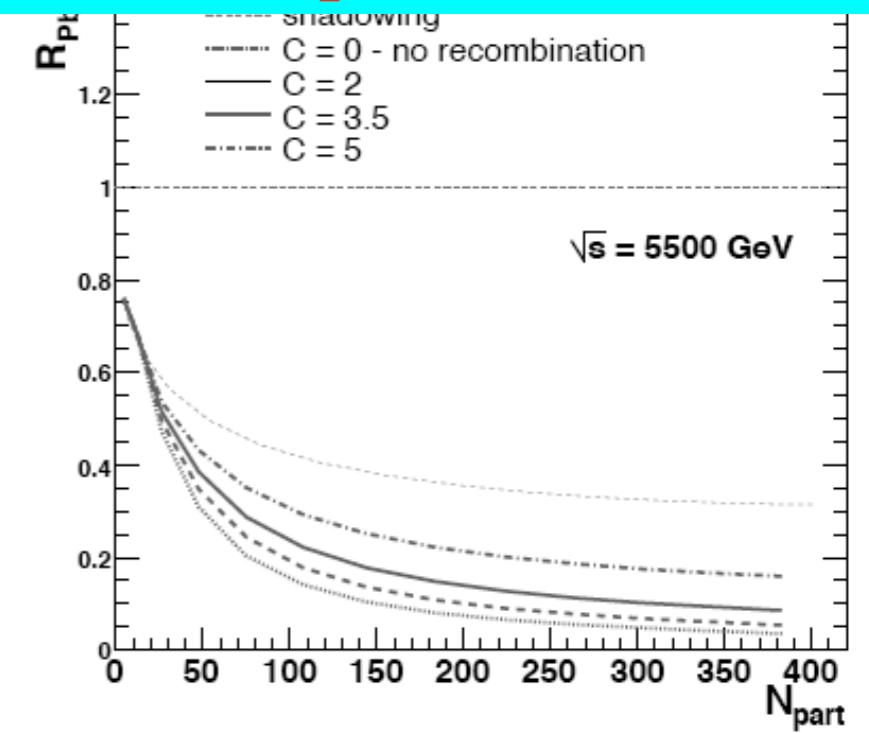
Andronic et al.: dependence on charm cross section; Thews et al.:  $p_T$ -broadening to verify



➡ Considerable uncertainties both in the production and in the suppression mechanisms limit our predictive power.



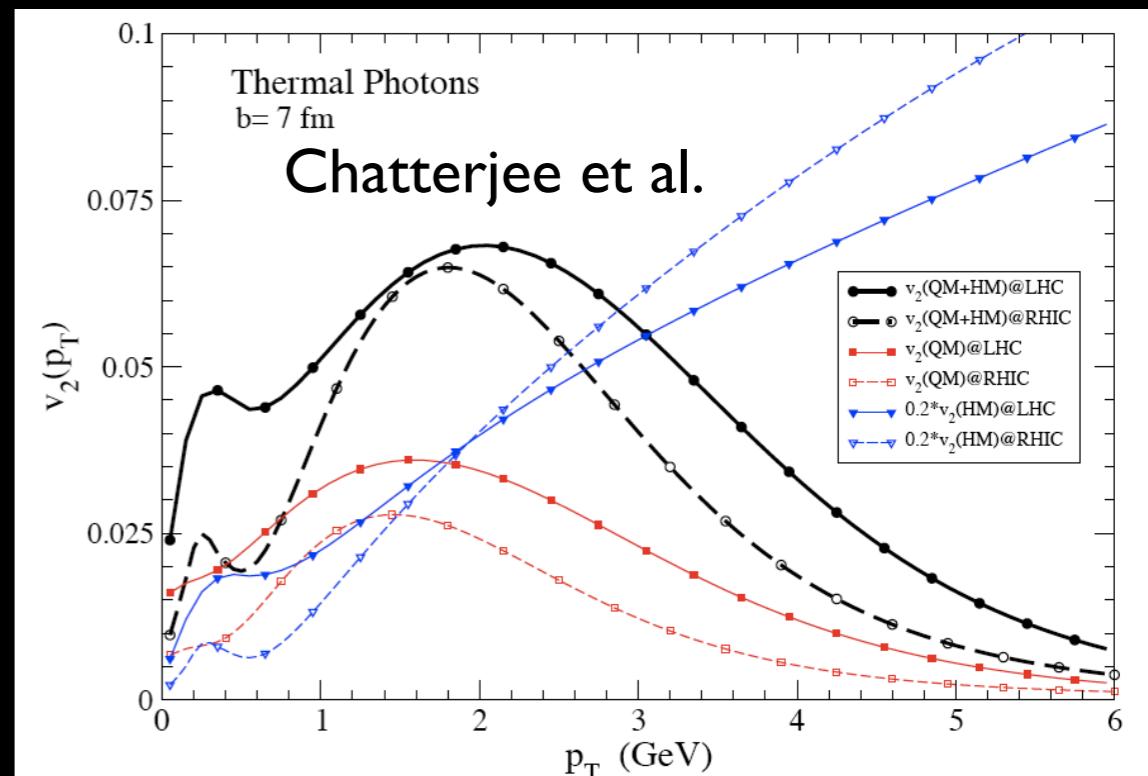
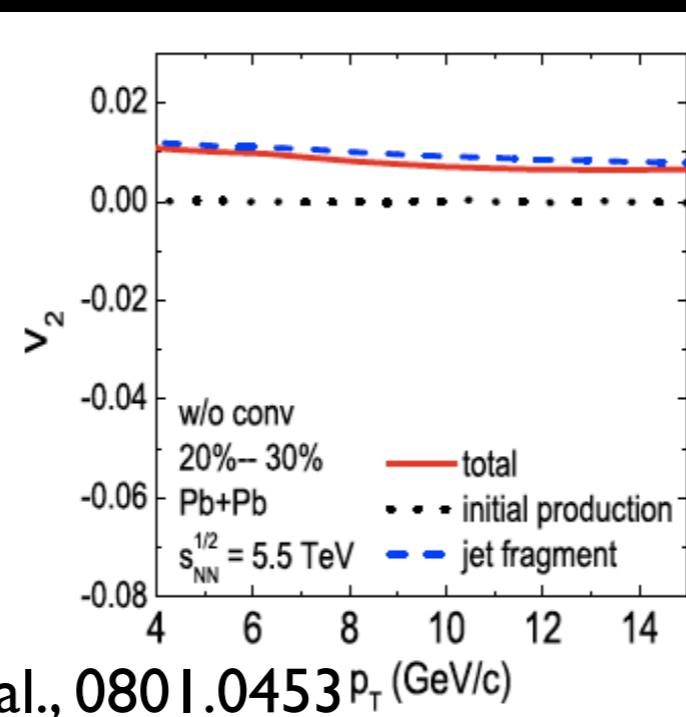
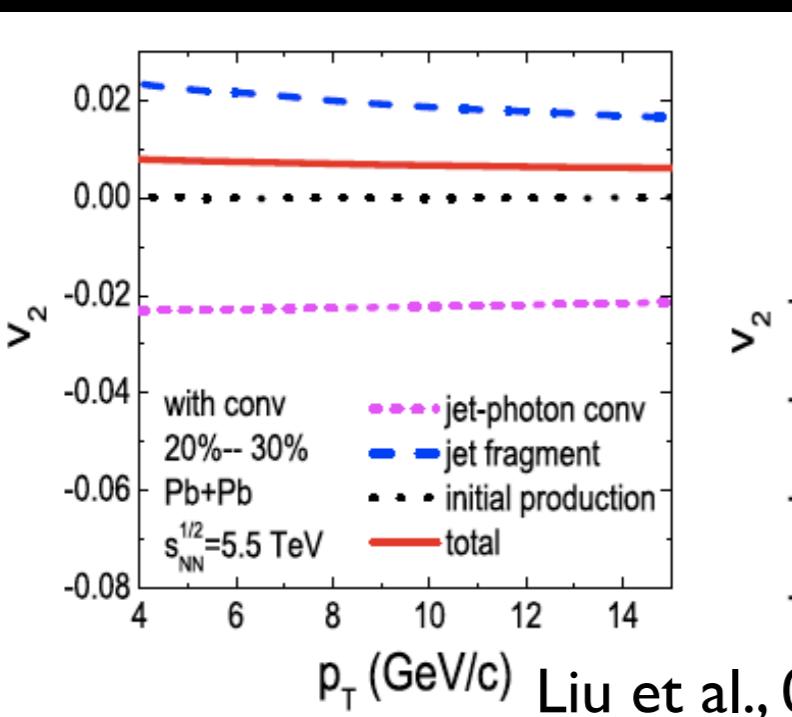
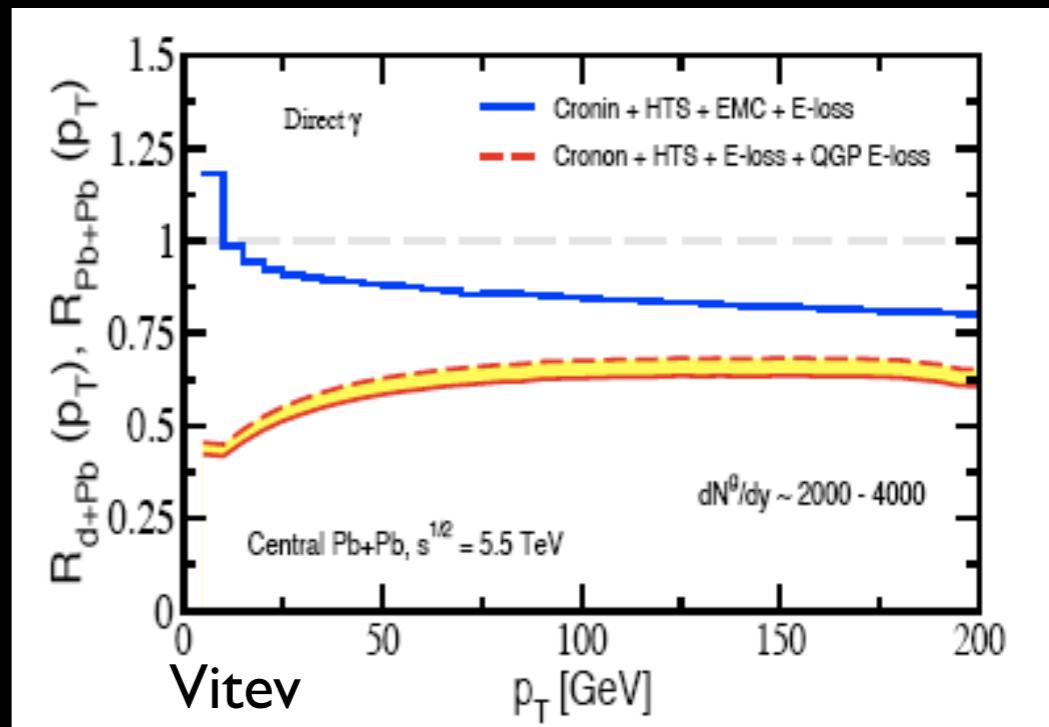
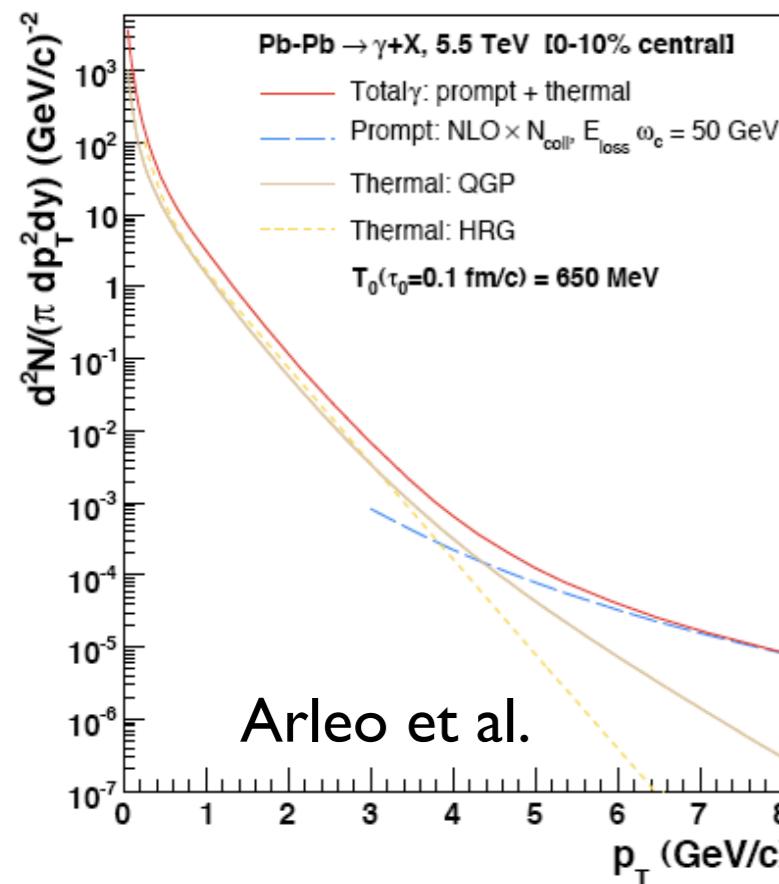
Vogt:  $p_T$ -dependent screening, no regeneration;  
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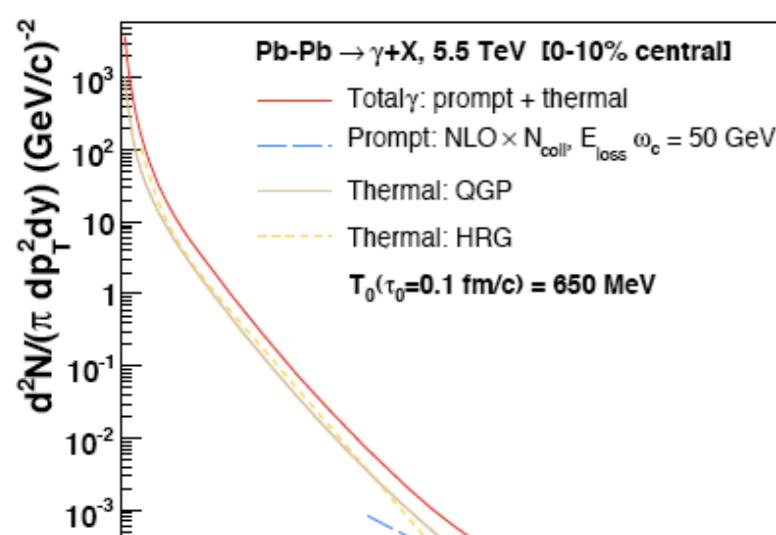
Capella et al., 0712.4331,  
comovers+reco.

# B.3. Leptonic probes and photons: photons

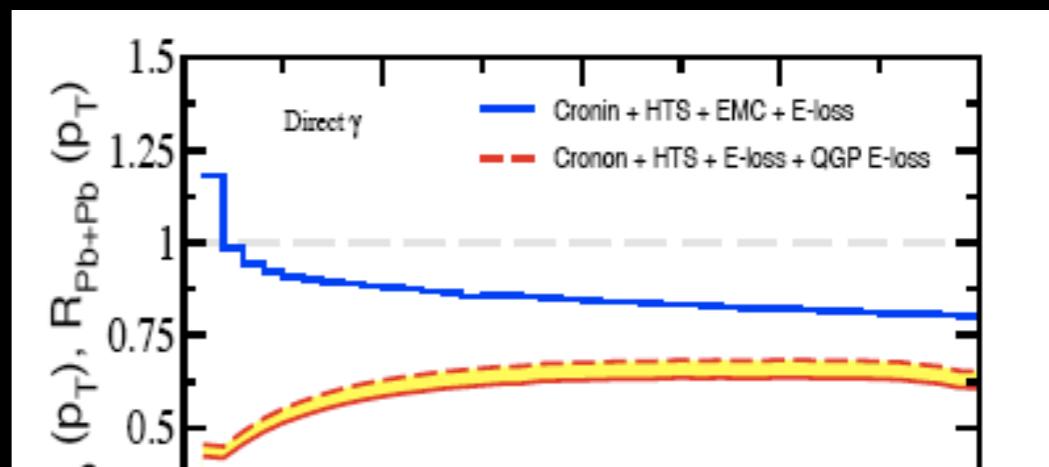
High  $p_T$ : quenching  
in fragmentation  
(also Arleo).  
Low  $p_T$ : new  
effects e.g.  
conversions,  
thermal,...  
expected (also  
Rezaeian et al.).



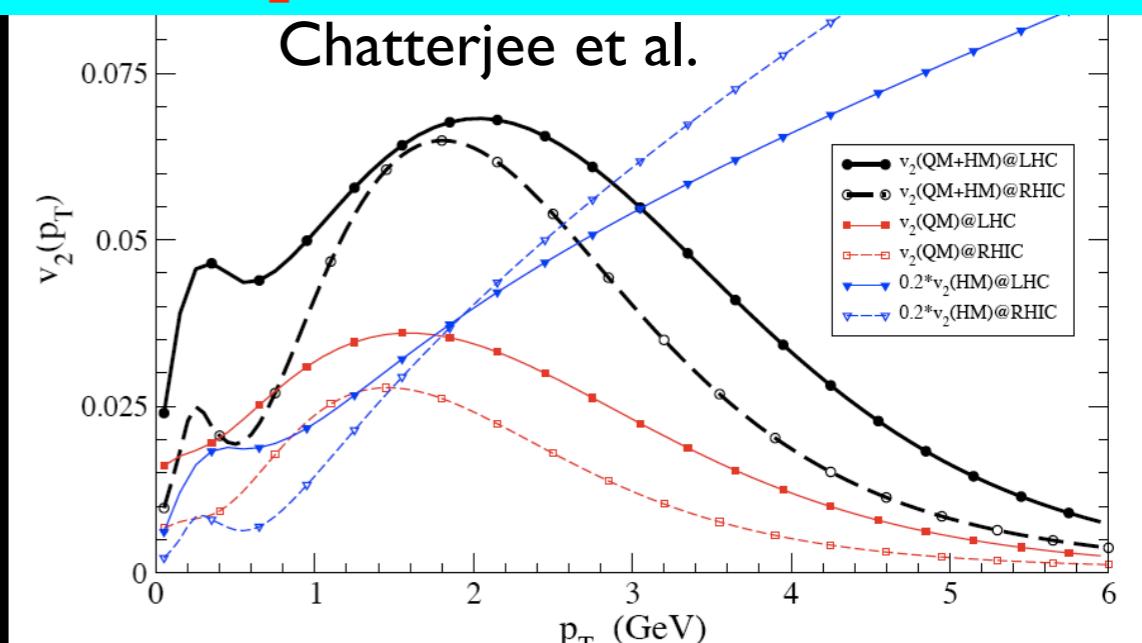
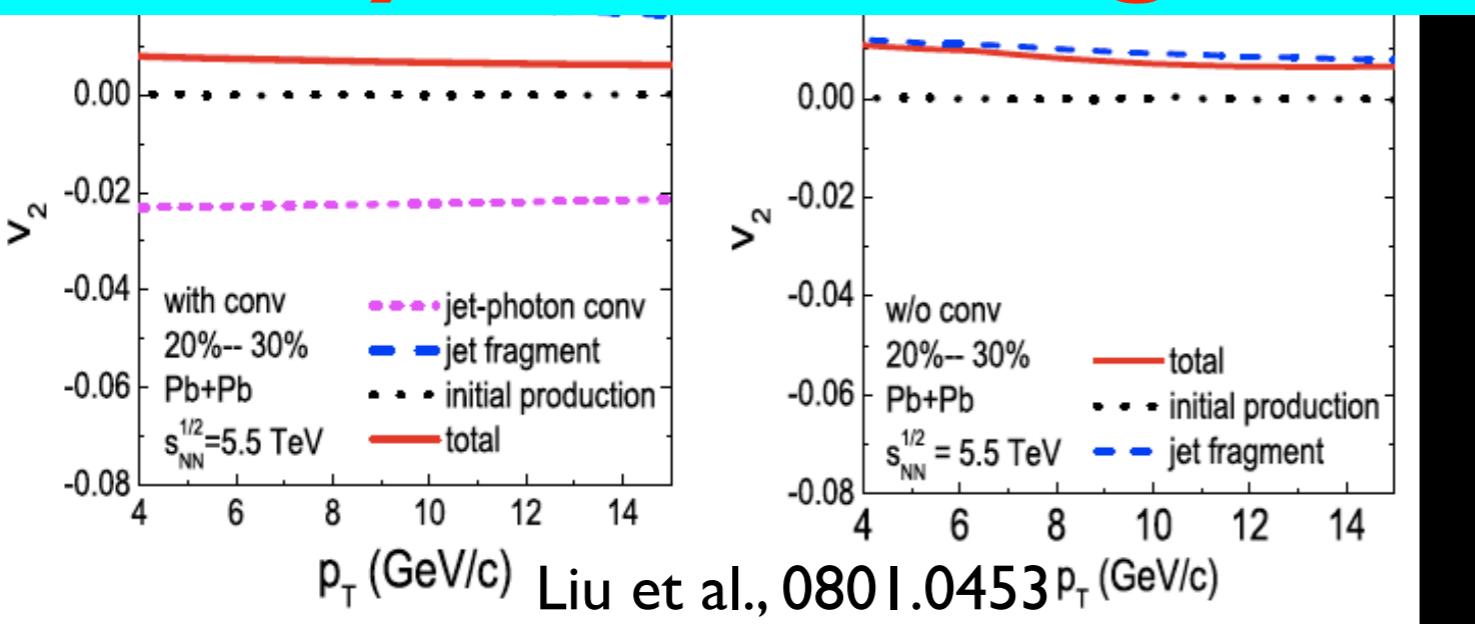
# B.3. Leptonic probes and photons: photons



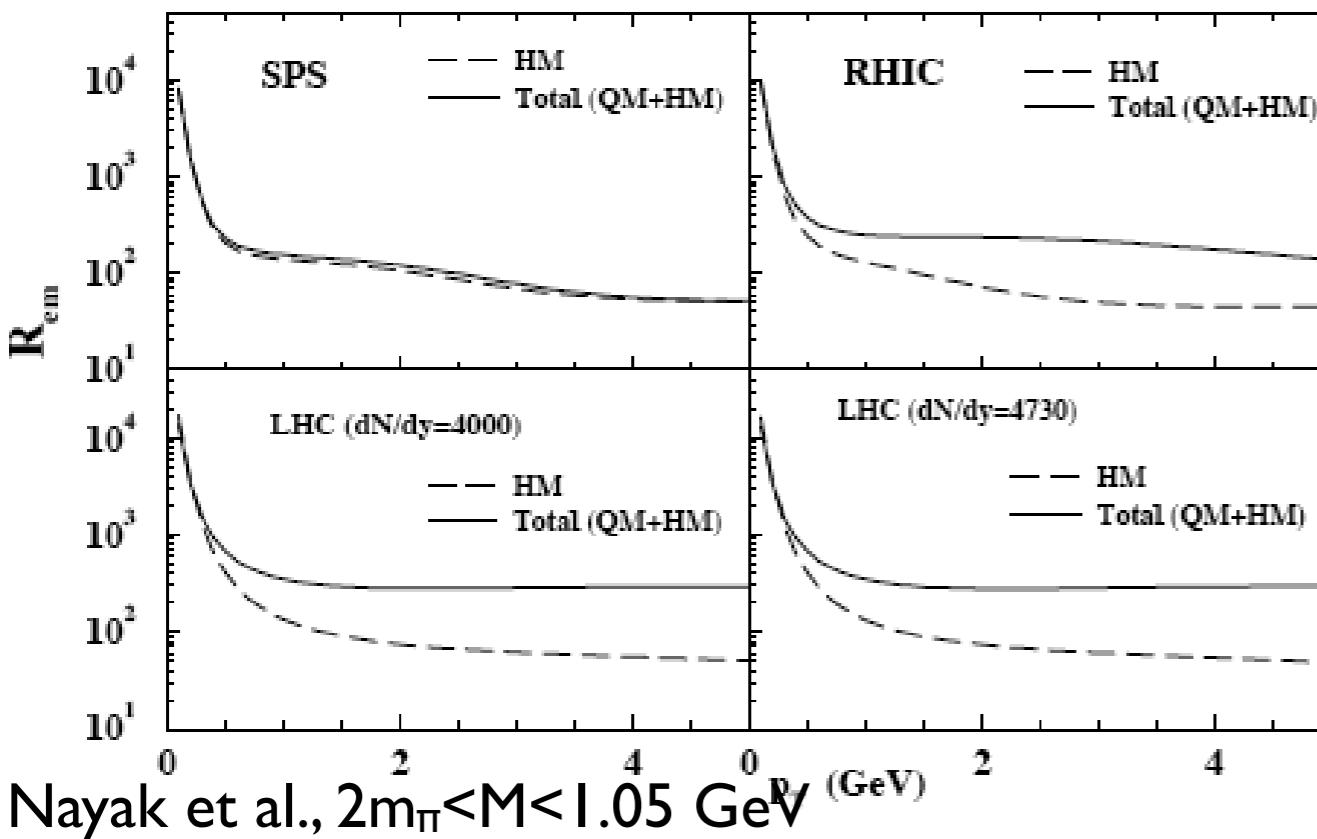
High  $p_T$ : quenching  
in fragmentation  
(also Arleo).  
Low  $p_T$ : new  
effects e.g.



- At the LHC, all the regions up to very high  $p_T$  can be studied.
- Control of the benchmark, from pp to PbPb, key to disentangle thermal production.

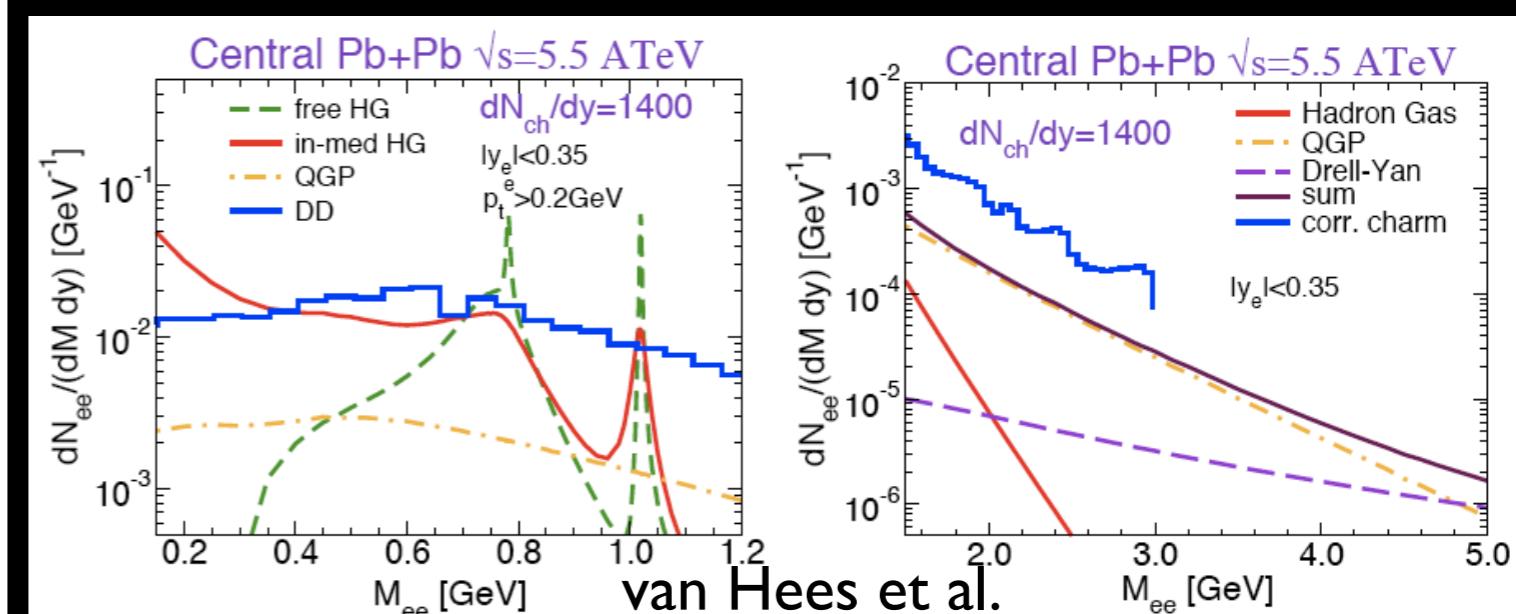
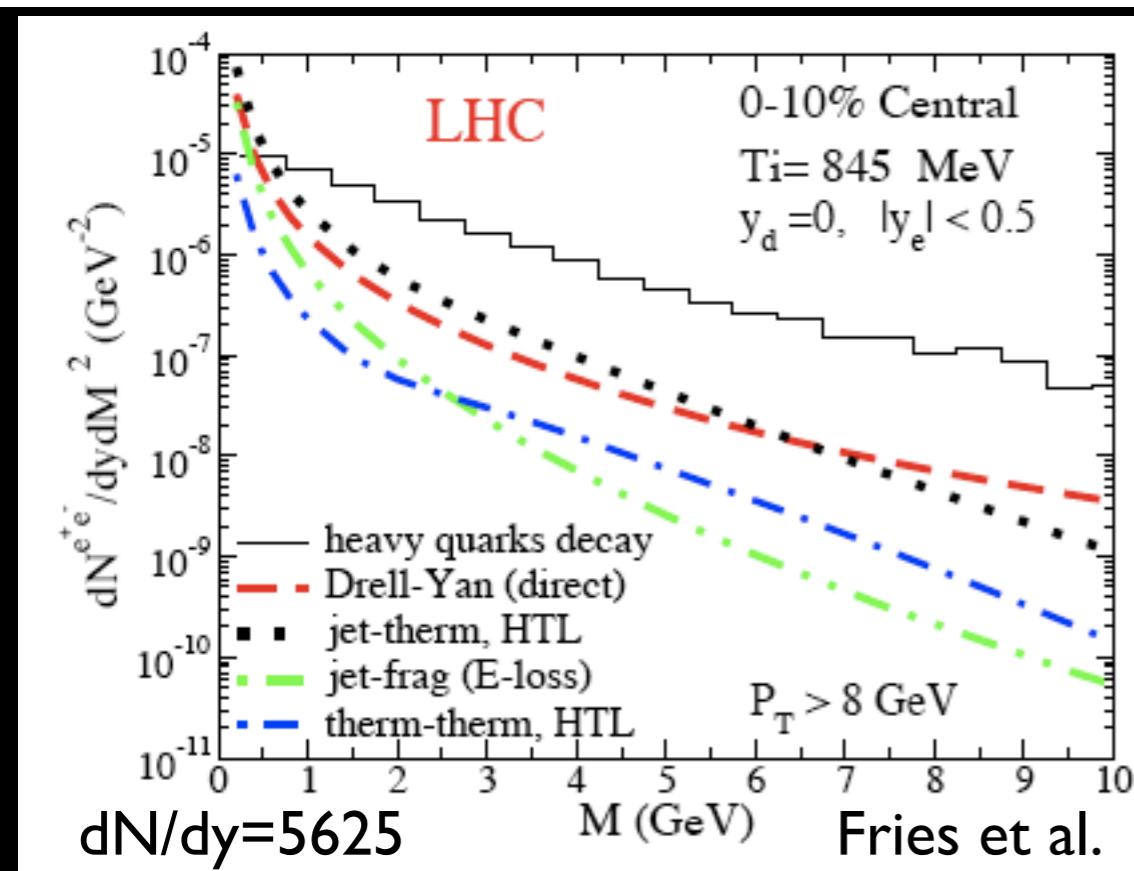


# B.3. Leptonic probes and photons: dileptons

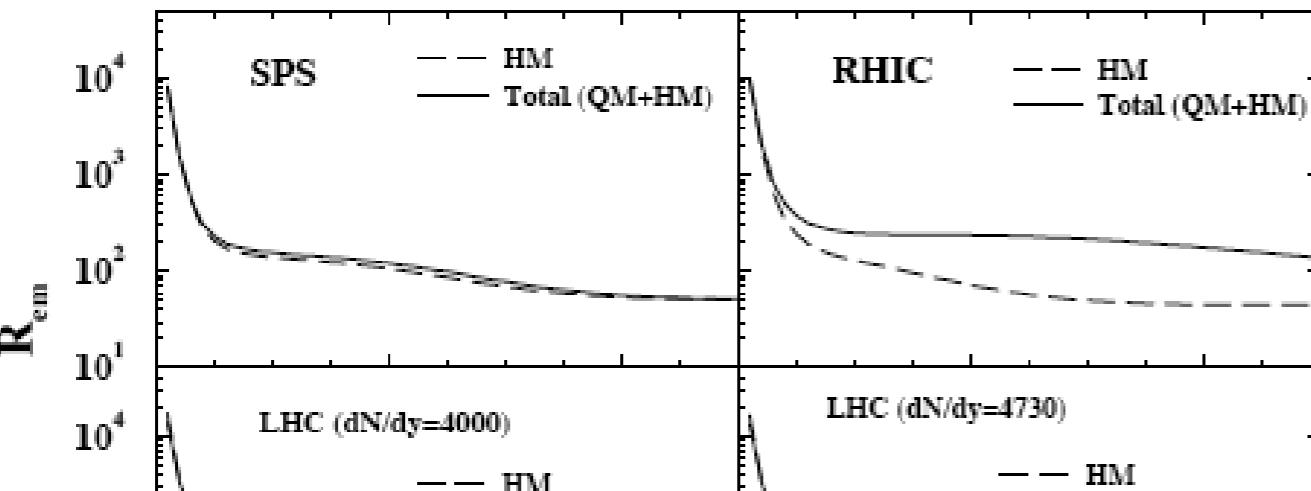


Ratio  $\gamma/\gamma^*$  shows a plateau at large  $p_T$  of thermal origin, independently of details of expansion, EOS,...

Dremin: Cherenkov radiation as the origin of the ‘broadening’ of the  $\rho$ .

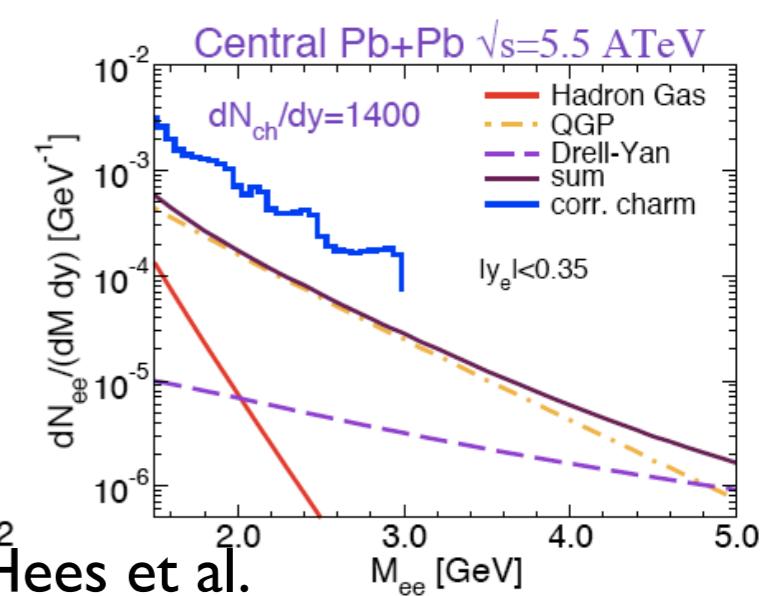
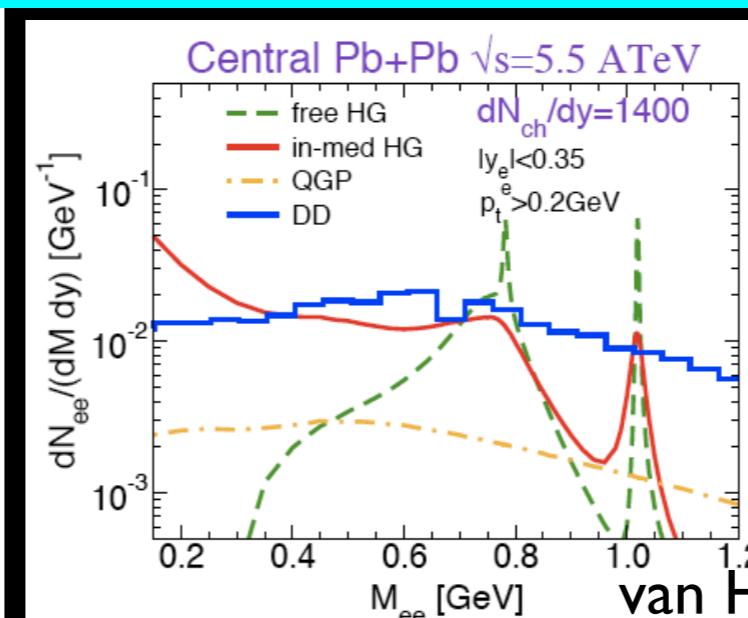
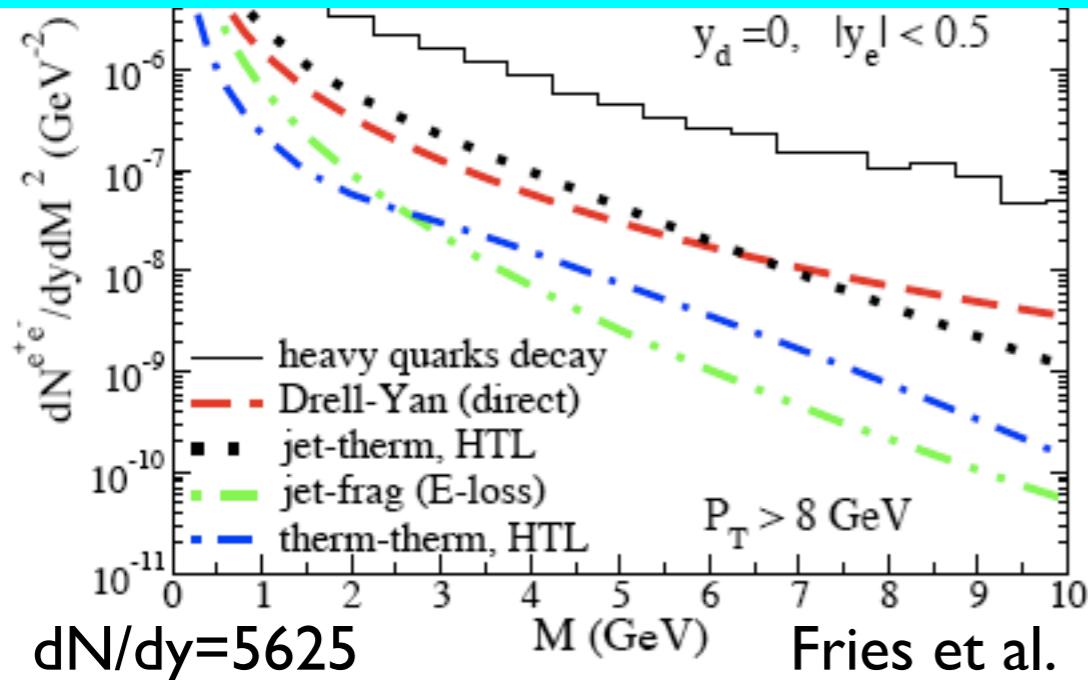


## B.3. Leptonic probes and photons: dileptons



Ratio  $\gamma/\gamma^*$  shows a plateau at large  $p_T$  of thermal origin, independently of details of expansion, EOS,...

- ⇨ Both low and high  $M$  and  $p_T$  required to disentangle different mechanisms.
- ⇨ But the huge HQ contribution looks really challenging.

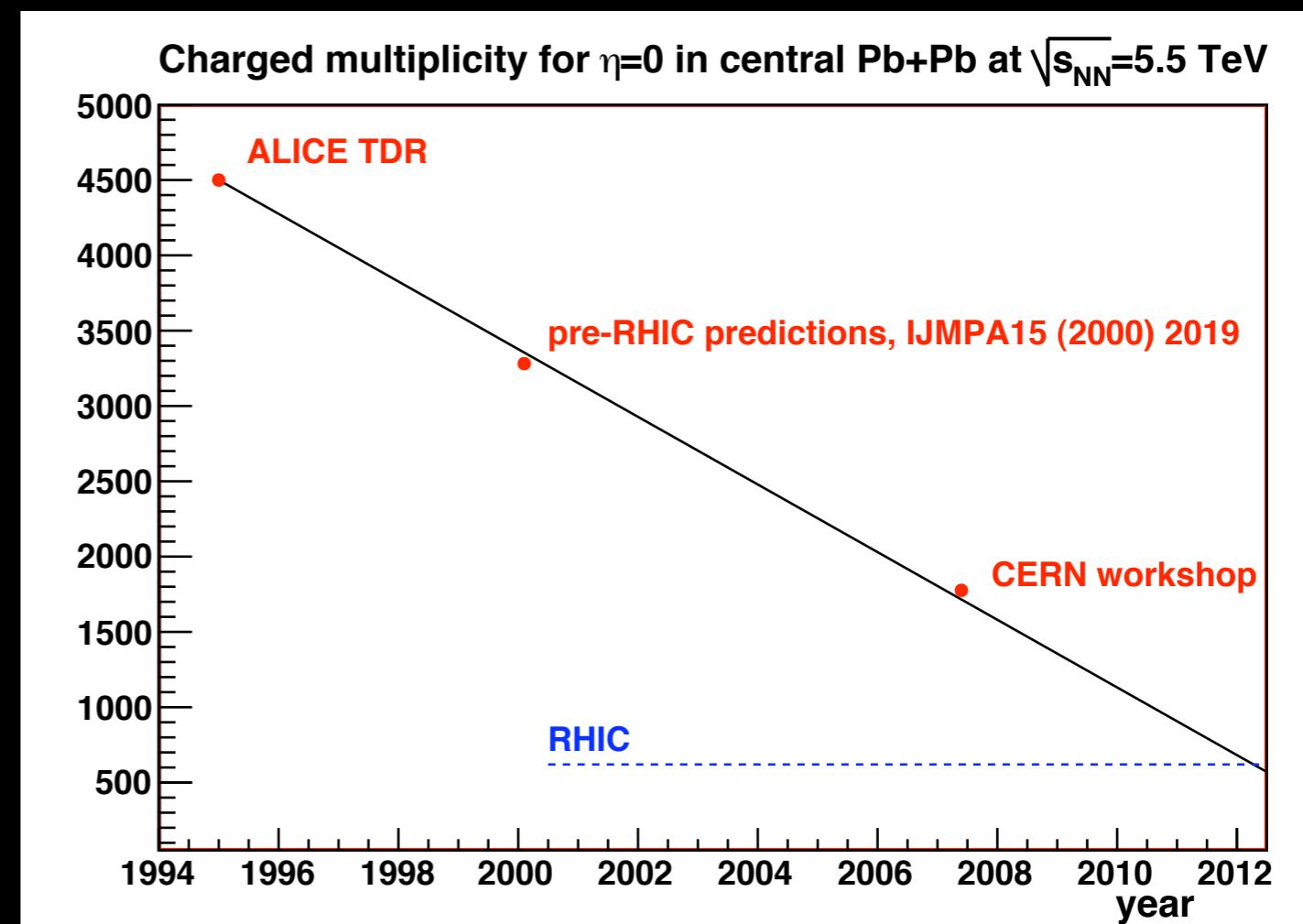
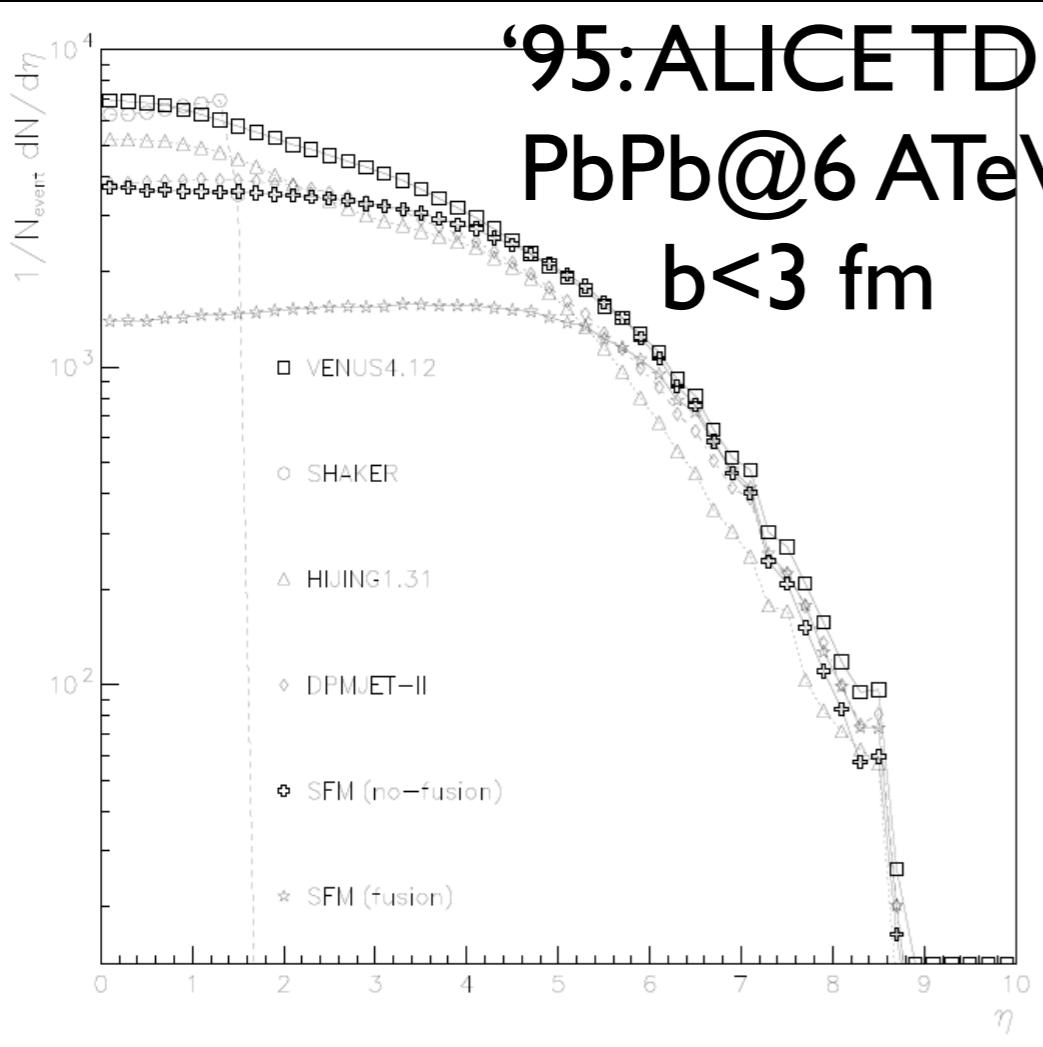


# RHIC-tested models face the LHC era

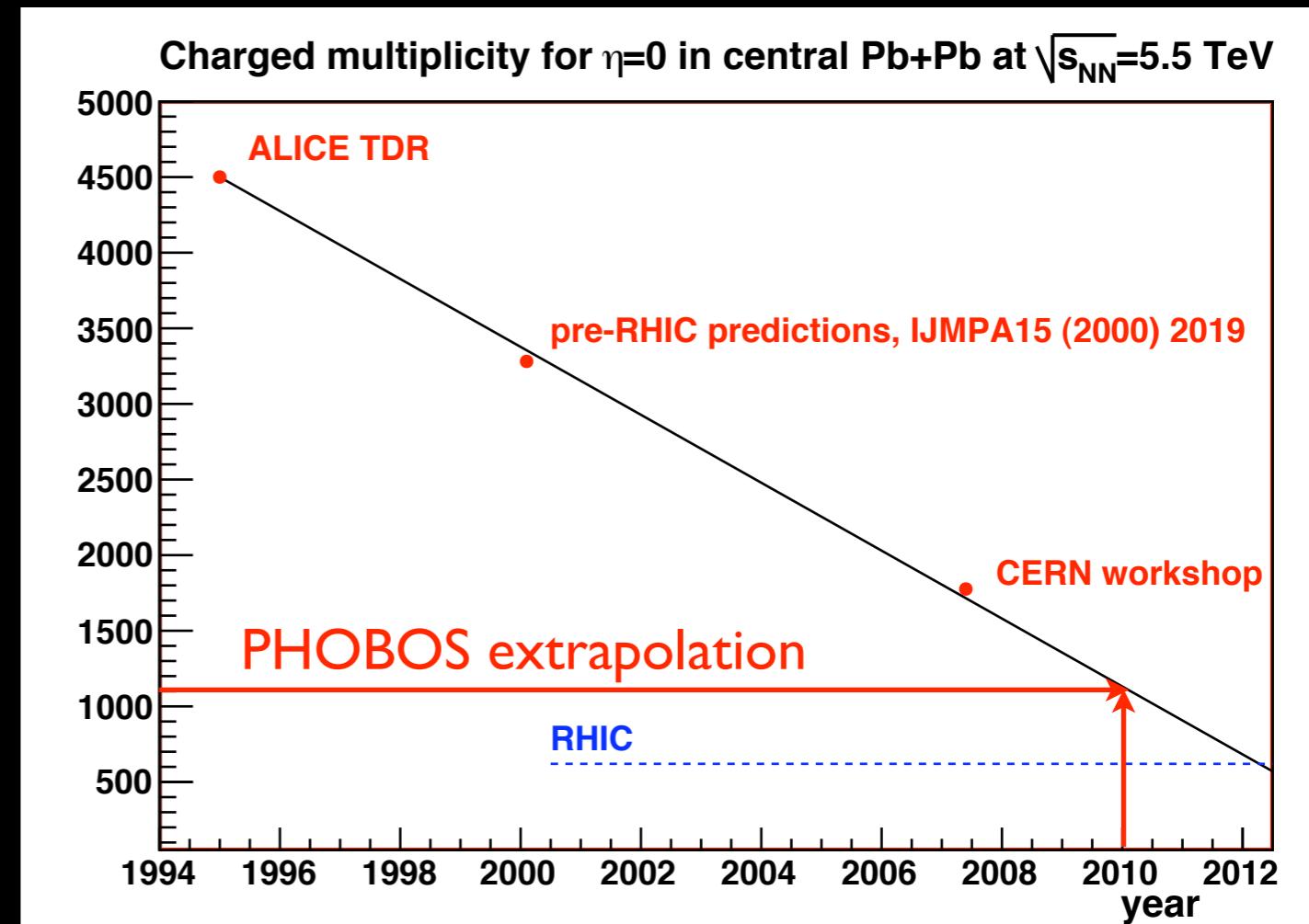
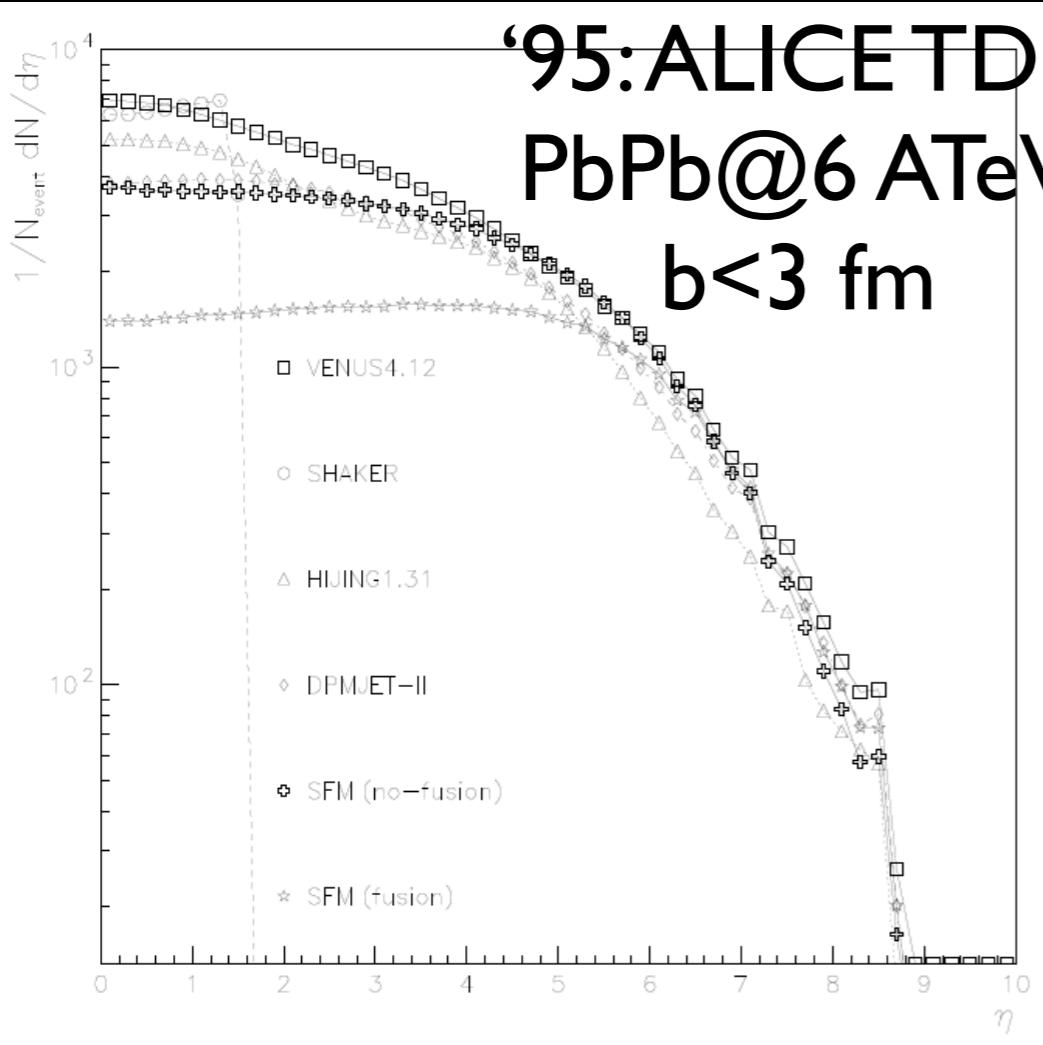
Observable at RHIC	Standard interpretation	Prediction for the LHC
Low multiplicity	Strong coherence in particle production	$dN_{ch}/d\eta _{\eta=0} < 2000$ for central collisions
$v_2$ in agreement with ideal hydro	Almost ideal fluid	Similar or smaller $v_2(p_T)$
Strong jet quenching	Opaque medium	$R_{AA}(20 \text{ GeV}) \sim 0.1-0.2$ for $\pi^0$

- \* This picture has motivated **new theoretical developments**: application of AdS/CFT, early thermalization, viscous hydro, CGC,...
- \* **Major deviations from expectations** will enlarge our understanding of Ultra-Relativistic Heavy-Ion Collisions: naive extrapolations tend to disagree with those from successful models at RHIC.
- \* **New experimental observables at LHC**: higher  $p_T$ , jets, higher quarkonium states,... demand **new theoretical tools**.

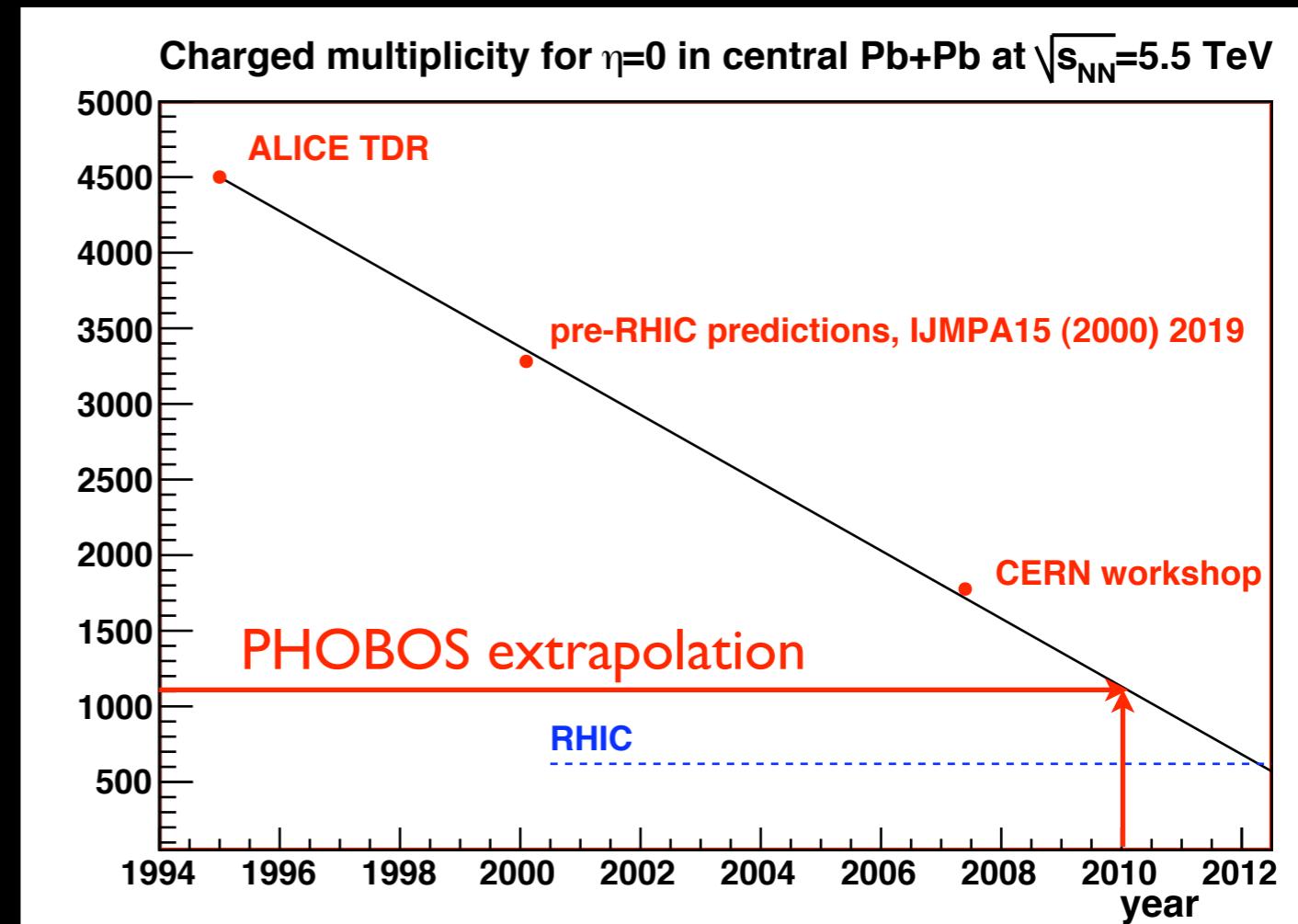
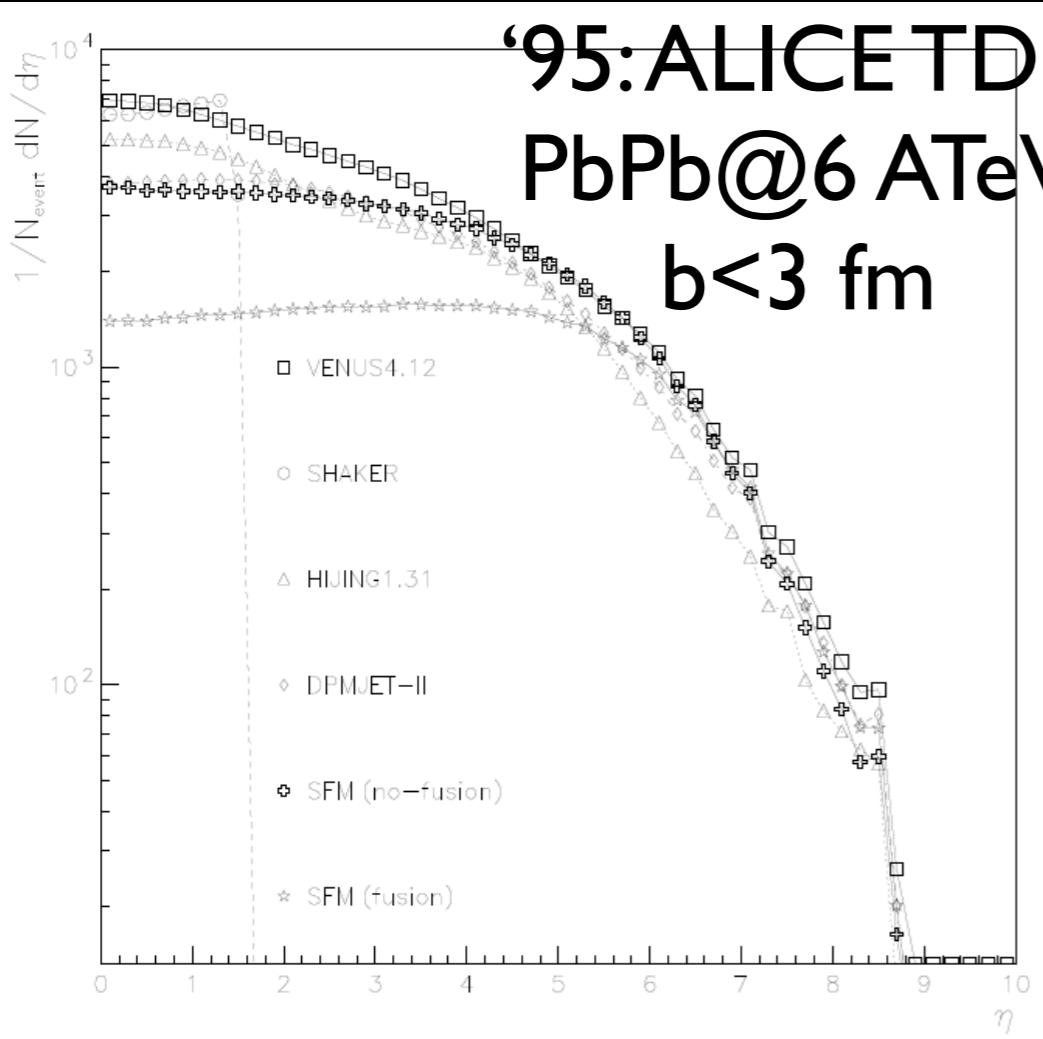
# To conclude: a little bit of history



# To conclude: a little bit of history



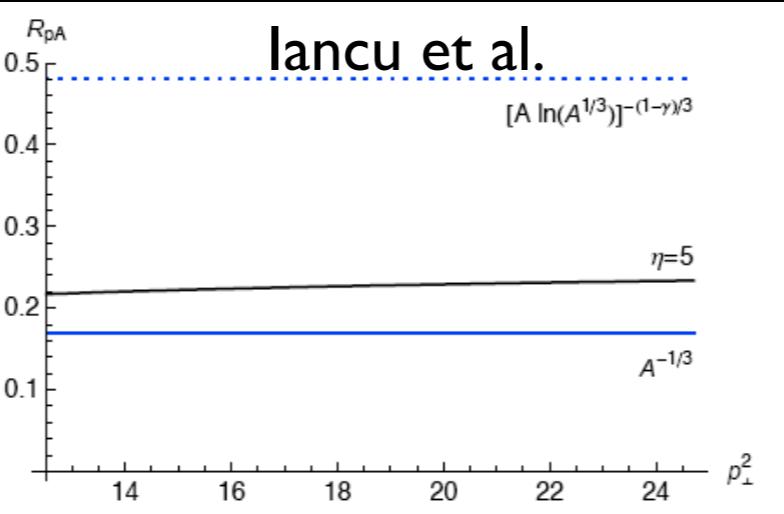
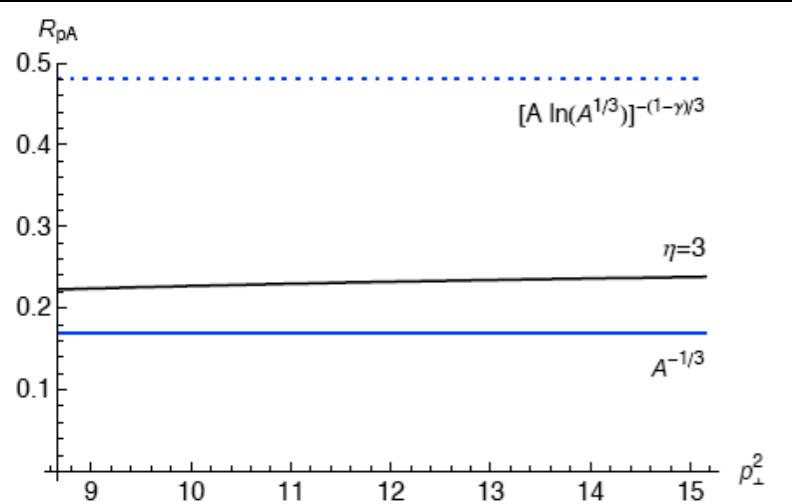
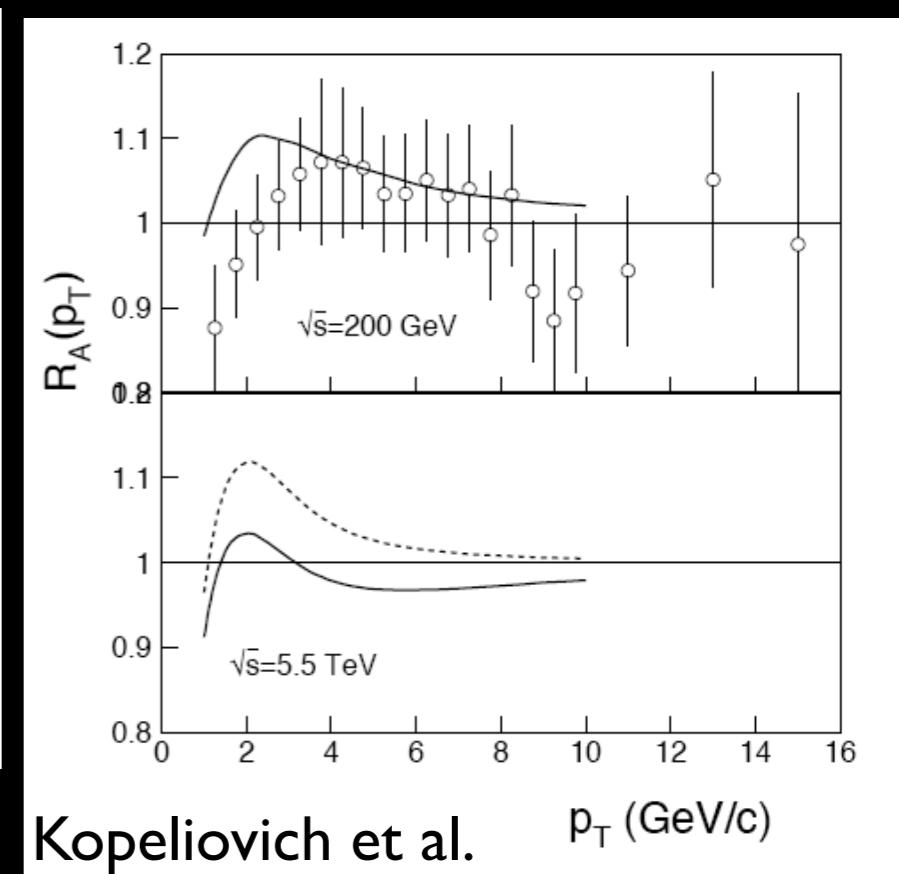
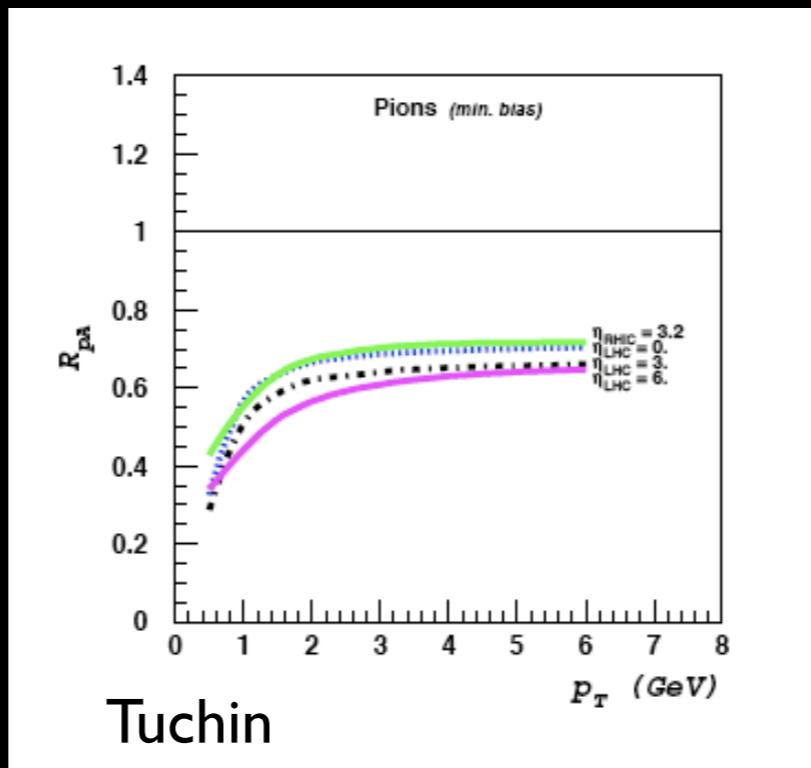
# To conclude: a little bit of history



Thanks to J. Albacete, F. Bopp, W. Busza, L. Cunqueiro, A. Dainese, A. El, K. Eskola, U. Heinz, C.-M. Ko, I. Lokhtin, G. Milhano, C. Pajares, V. Pantuev, T. Renk, V. Topor Pop, R. Venugopalan, I. Vitev, X. N. Wang, K. Werner and G. Wolschin for feedback on their predictions, and to J. Albacete, J. Casalderrey-Solana, K. Eskola, E. Ferreiro, U. Heinz, P. Jacobs, C. Salgado, X. N. Wang and U. Wiedemann for discussions on the talk.

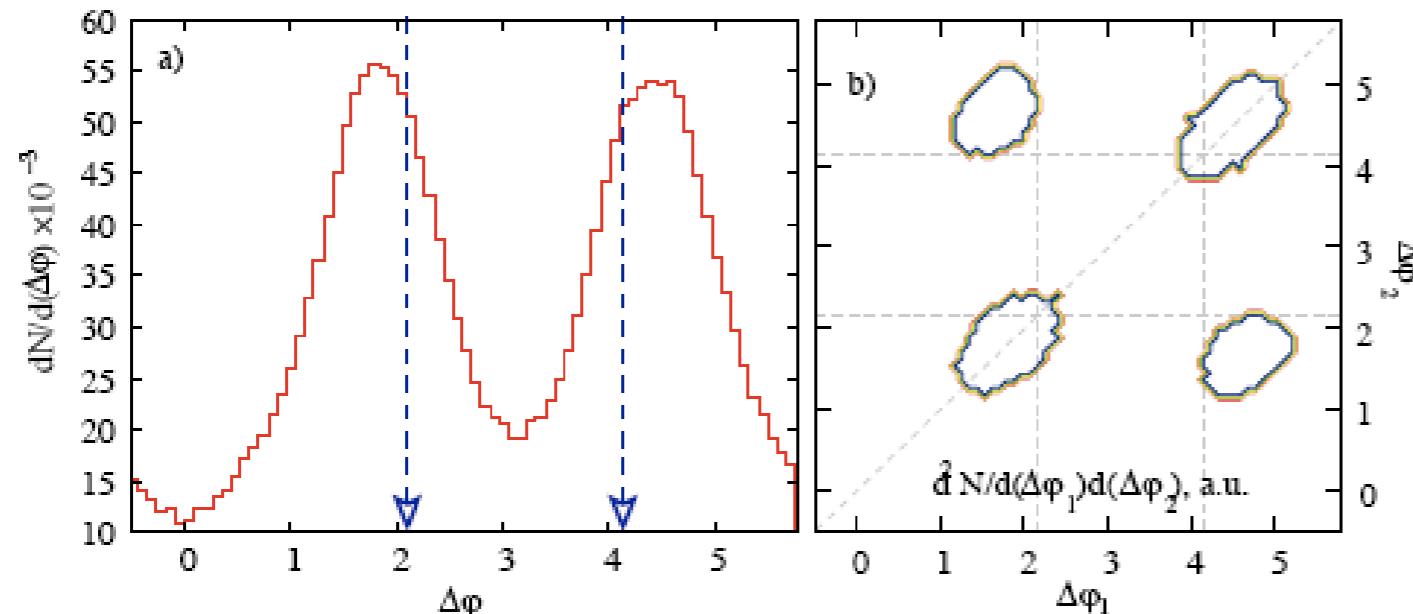
# Backup I: pA

**Forward  $R_{pA}(p_T)$  as a probe of high density dynamics:** saturation for light (Armesto et al., Boer et al.) and heavy (Fujii et al., Tuchin), or absence of saturation (Arsene et al.);  $\Lambda$  polarization (Boer et al.).

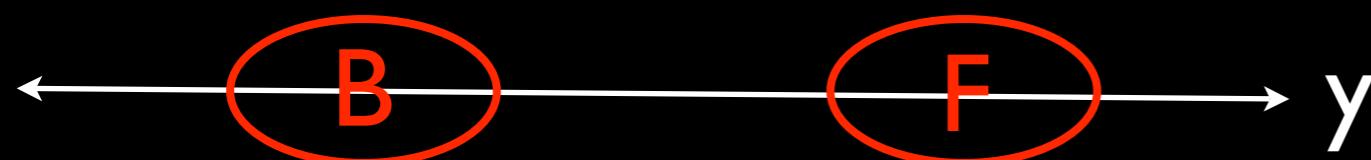


Within CGC, total shadowing due to fluctuations (Kozlov et al.) or running coupling effects.

## A.4. Correlations at low $p_T$ (II):



Bauchle et al., **Mach cones in hydro**, peak displaced from naive expectation; Betz et al., gradual energy deposition required. Other mechanism possible: **Cherenkov** (Dremin), and **instabilities** (Mannarelli et al.) - unstable modes for jets with speed larger than  $c_s$ , which peak at larger angles for larger speed.



$$\langle n_B \rangle_F = a + b n_F, \quad b \equiv D_{FB}^2 / D_{FF}^2,$$

$$b = \frac{\langle n_B \rangle / \langle n_F \rangle}{1 + K / \langle n_F \rangle}$$

$$1/K = D_{FB}^2 / \langle n_F \rangle \langle n_B \rangle$$

Dias de Deus et al.: **FB correlations** may help to establish the dynamics of particle production.

# Backup II: fluctuations

$$\Psi_{N_1/N_2}^{N_1} = \frac{dN_1}{dy} v_{N_1/N_2}^{dyn}$$

$$v_{N_1/N_2}^{dyn} = (\sigma_{N_1/N_2}^{dyn})^2 = \sigma_{N_1/N_2}^2 - (\sigma_{N_1/N_2}^{Poisson})^2 = \\ = \frac{\langle N_1(N_1-1) \rangle}{\langle N_1 \rangle^2} + \frac{\langle N_2(N_2-1) \rangle}{\langle N_2 \rangle^2} - 2 \frac{\langle N_1 N_2 \rangle}{\langle N_1 \rangle \langle N_2 \rangle}$$

Torrieri: fluctuations in particle ratios as a tool to verify the statistical model and decide which ensemble or non-eq. situation holds.

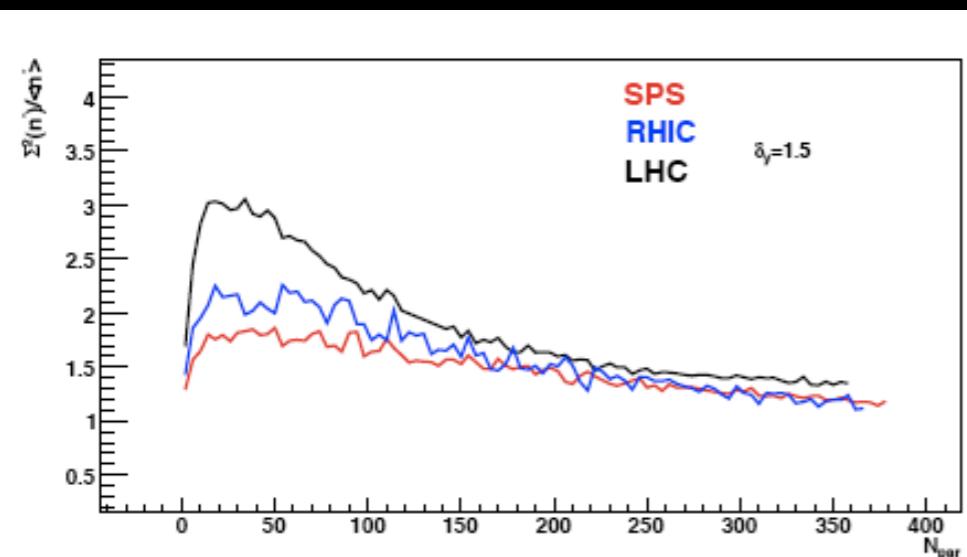
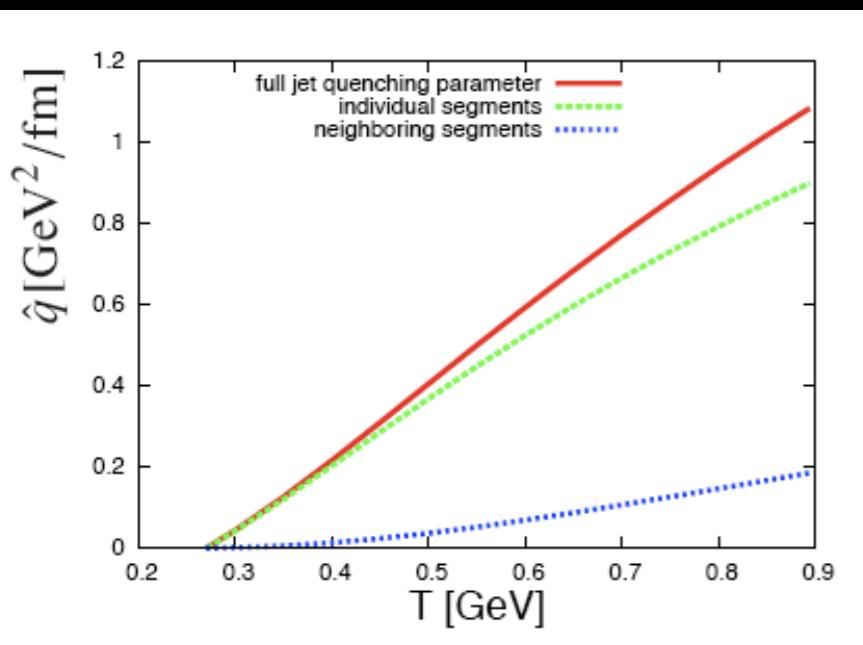


Figure 71: Scaled variance on negatively charged particles at, from up to down, LHC, RHIC and SPS.

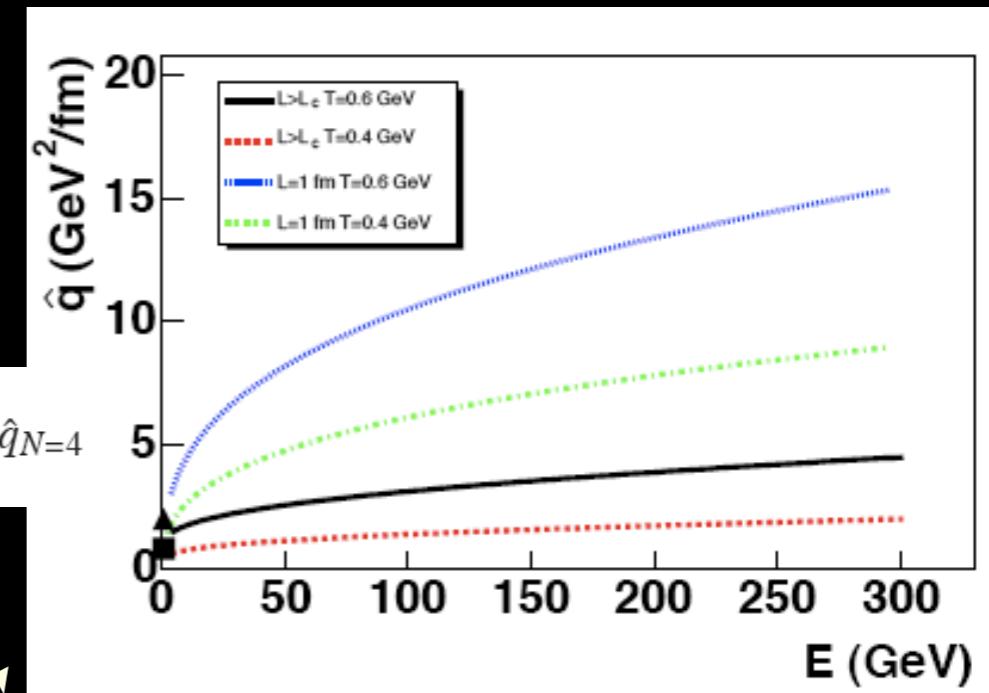
Cunqueiro et al.: multiplicity fluctuations determined by the number of coherent particle sources, possibility to verify phase transition scenarios.

# B.I. High- $\text{p}_T$ observ.: theory



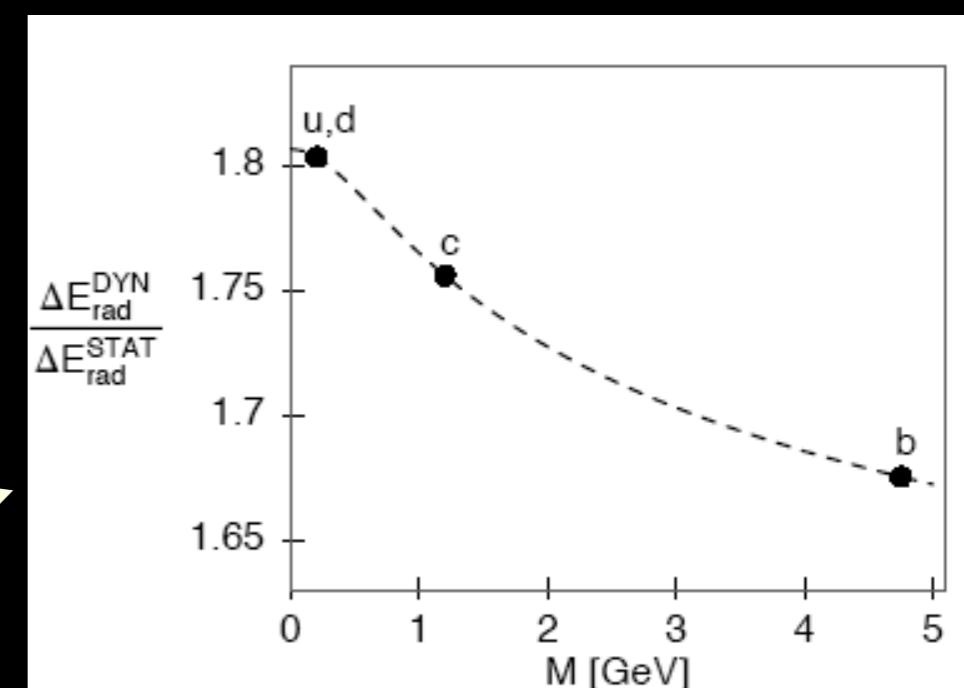
$$\hat{q}_{SYM} = \frac{\pi^{3/2} \Gamma(3/4)}{\Gamma(5/4)} \sqrt{\lambda} T^3$$

$$\hat{q}_{QCD} = \sqrt{\frac{s_{QCD}}{s_{N=4}}} \hat{q}_{N=4} = \sqrt{\frac{47.5}{120}} \hat{q}_{N=4} \approx 0.63 \hat{q}_{N=4}$$

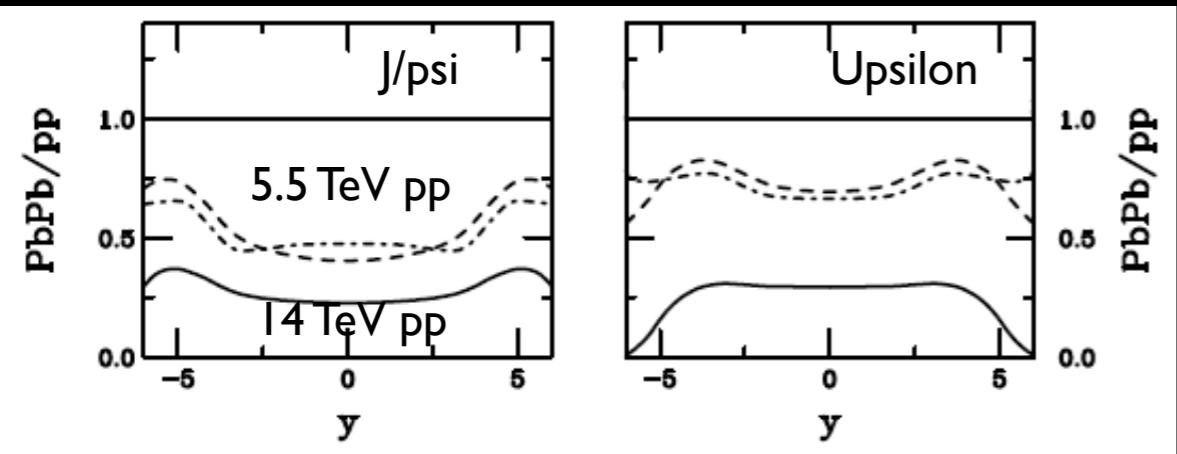


First-principle calculations of  $\hat{q}$  in the vacuum correlator model (Antonov et al.), in  $N=4$  SYM (H. Liu et al.) and in ThFT with coherence effects (Casalderrey et al.).

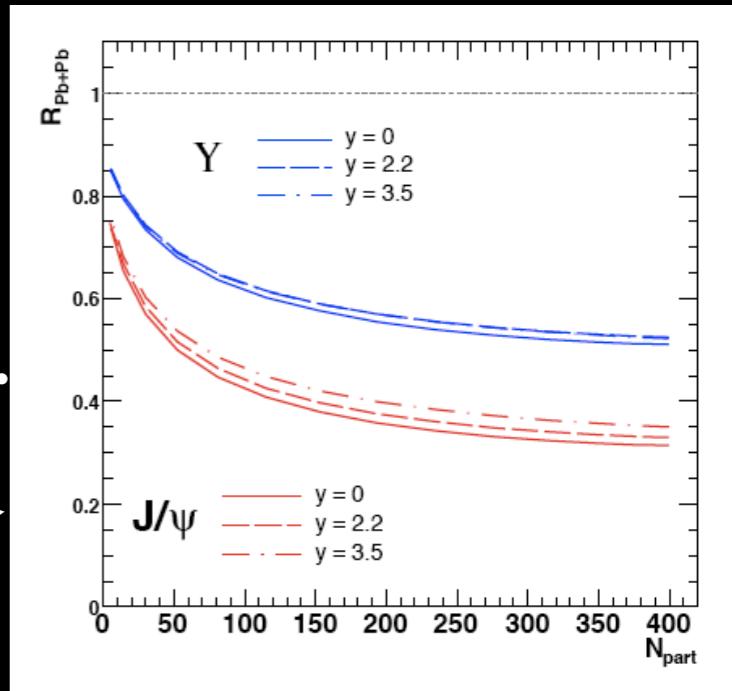
Inclusion of elastic eloss in DGLAP-like evolution (Pirner et al.) and consideration of a **dynamical medium** in GLV eloss (Djordjevic et al.).



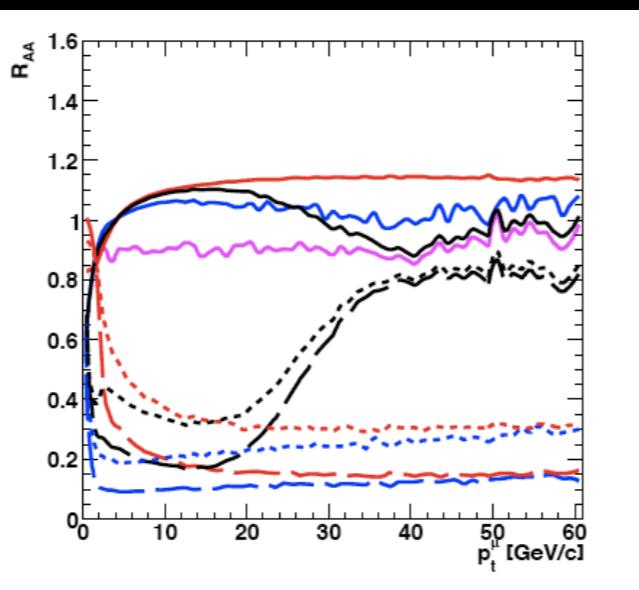
## B.2. HQ and quarkonium: reference



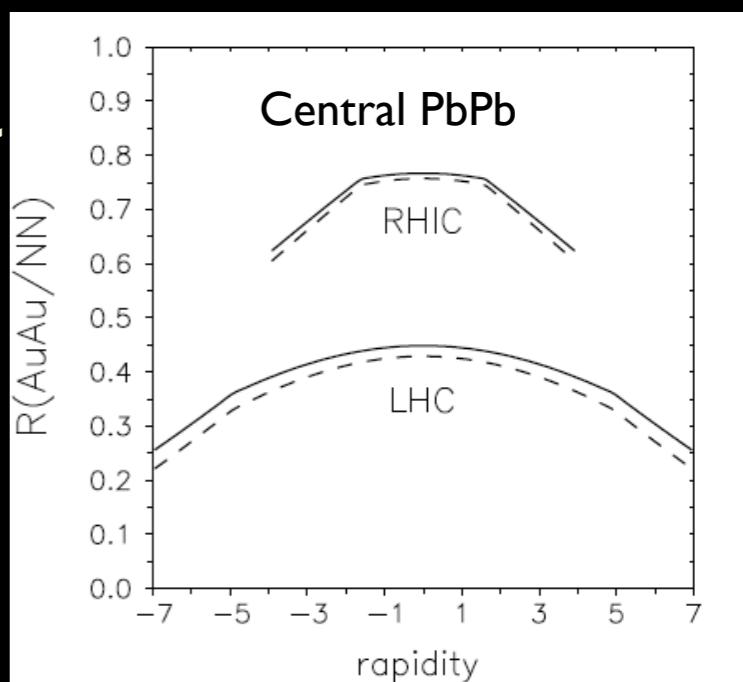
Arsene et al.,  
Gribov  
shadowing (HT).



Vogt, effect of npdf's (LT: EKS and nDSg): results very similar to other approaches.

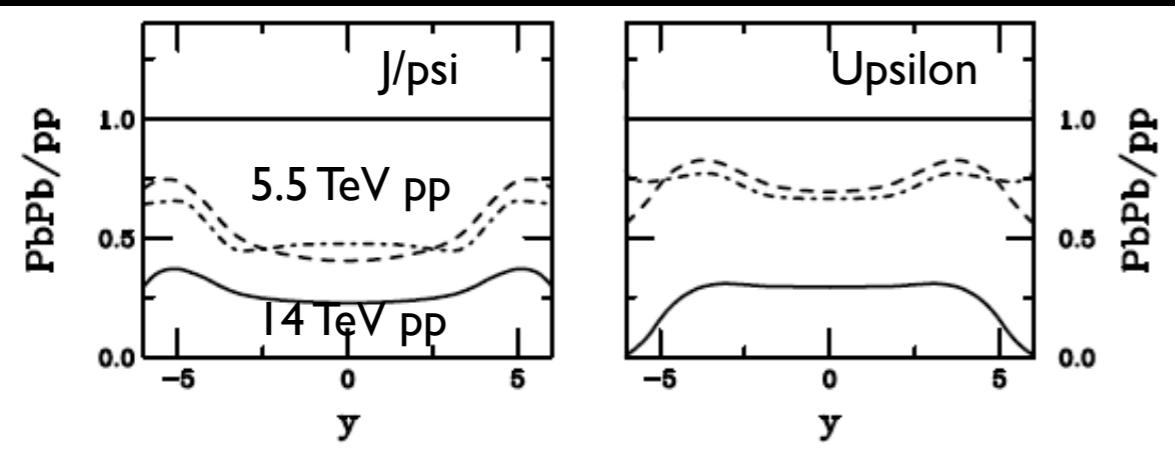


Kopeliovich et al., LT+HT process-dependent shadowing.

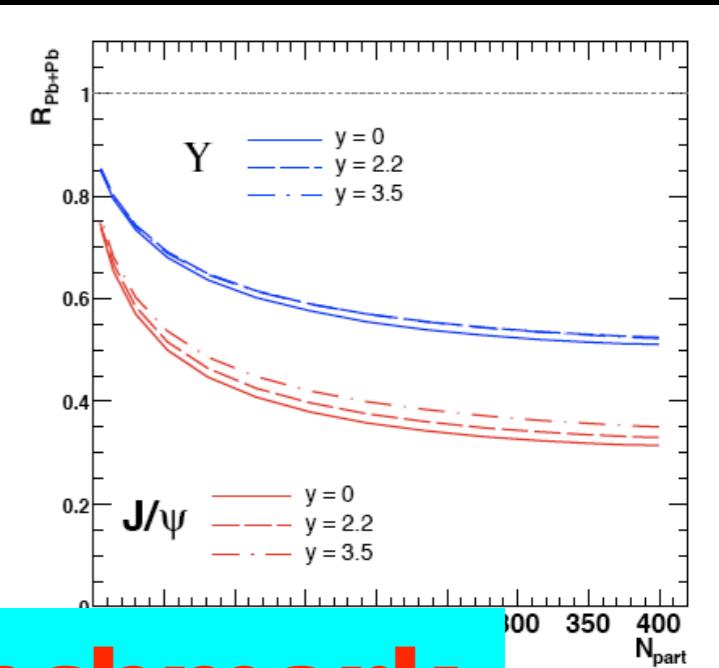


Conesa del Valle et al.: muons from  $W, Z$  as a reference for  $R_{AA}$ .

## B.2. HQ and quarkonium: reference

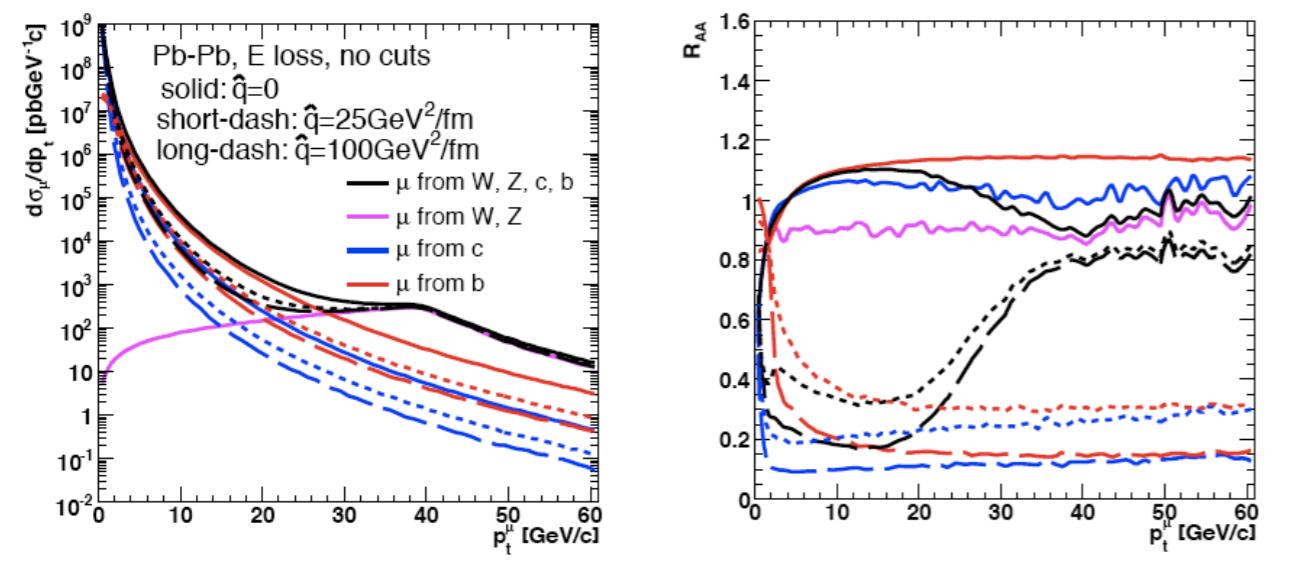


Arsene et al.,  
Gribov  
shadowing (HT).



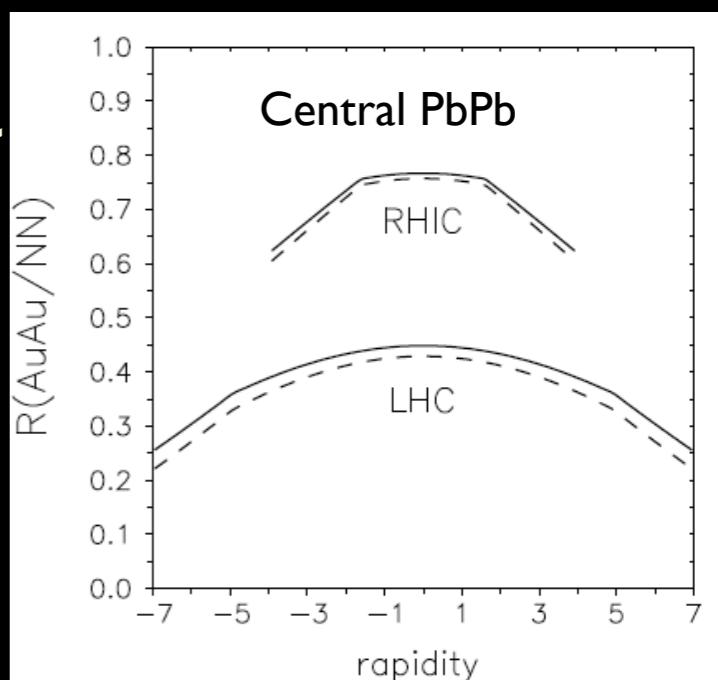
Vogt, effect of npdf's (LT: EKS and nDS) and Lutz-Krebs (HT)

**pPb most welcome for the benchmark.**



Conesa del Valle et al.: muons  
from  $W, Z$  as a reference for  $R_{AA}$ .

Kopeliovich et  
al., LT+HT  
process-  
dependent  
shadowing.



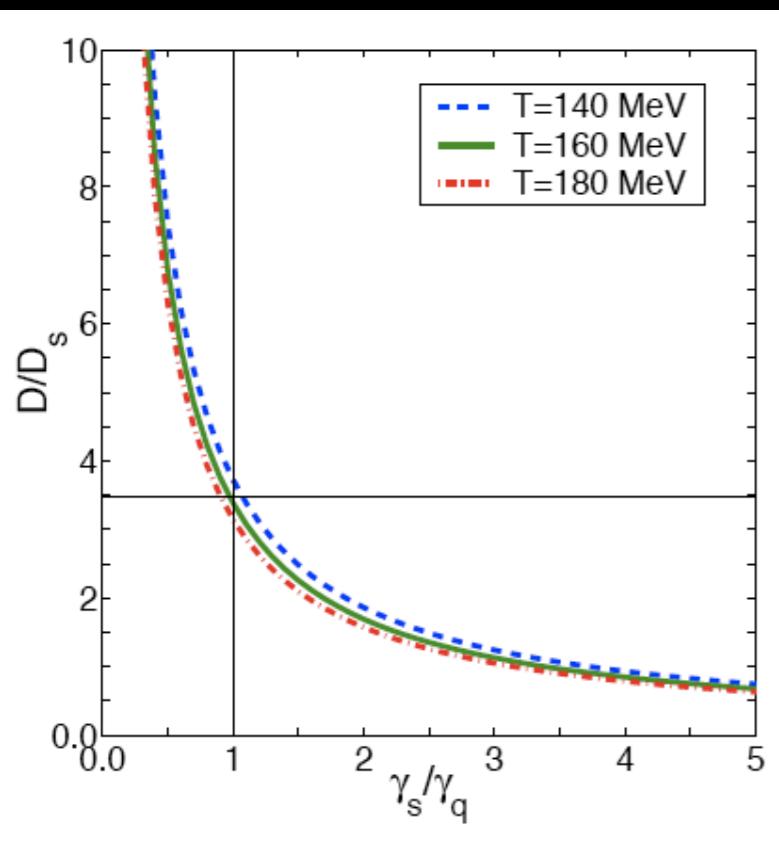
## B.2. HQ and quarkonium: ratios

$D^-/D^+$	$\bar{D}_0/D_0$	$D^{*-}/D^{*+}$	$D_s^-/D_s^+$	$\bar{\Lambda}_c/\Lambda_c$	$D^+/D_0$	$D^{*+}/D_0$
1.00(0)	1.01(0)	1.01(0)	1.00(1)	1.00(1)	0.425(18)	0.387(15)

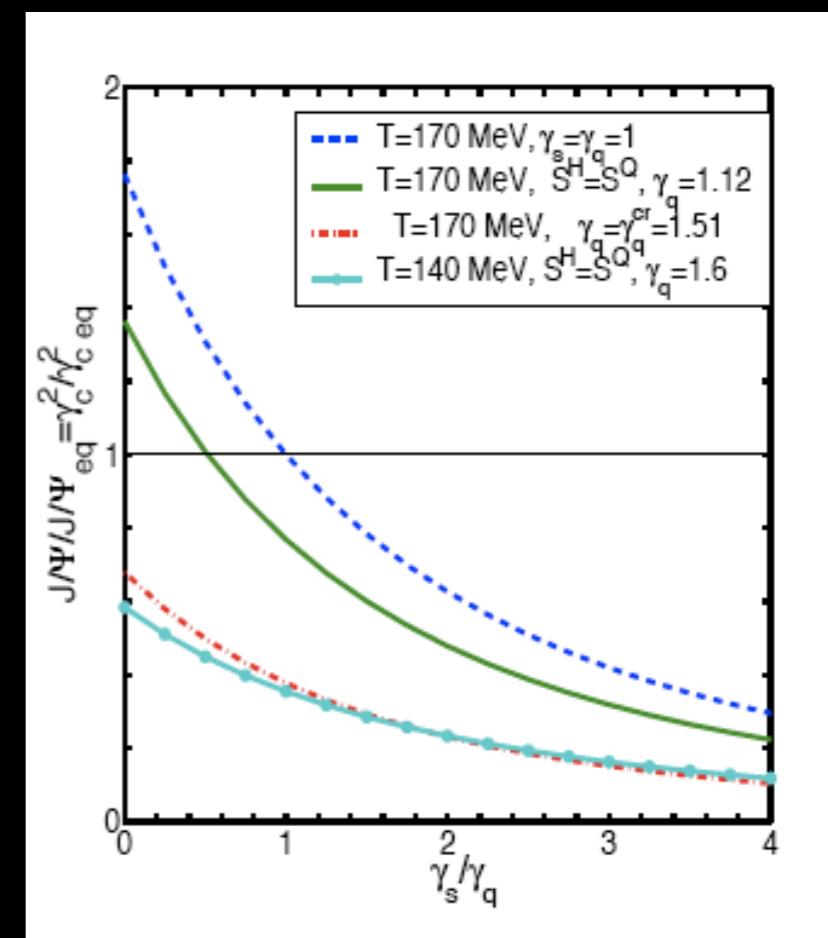
  

$D_s^+/D_0$	$\Lambda_c/D_0$	$\psi'/\psi$	$\eta_c/\psi$	$\chi_{c1}/\psi$	$\chi_{c2}/\psi$
0.349(14)	0.163(16)	0.031(3)	0.617(14)	0.086(5)	0.110(8)

Andronic et al.: ratios at equilibrium in the statistical model.



Kuznetsova et al.: strangeness oversaturation may lead to modifications in the ratios and to a suppression in the production of ccbar states.

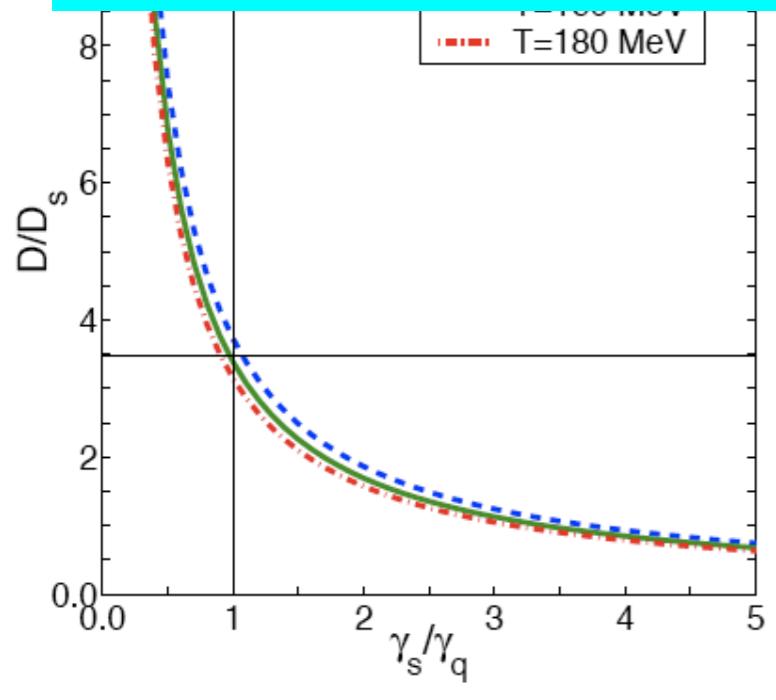


## B.2. HQ and quarkonium: ratios

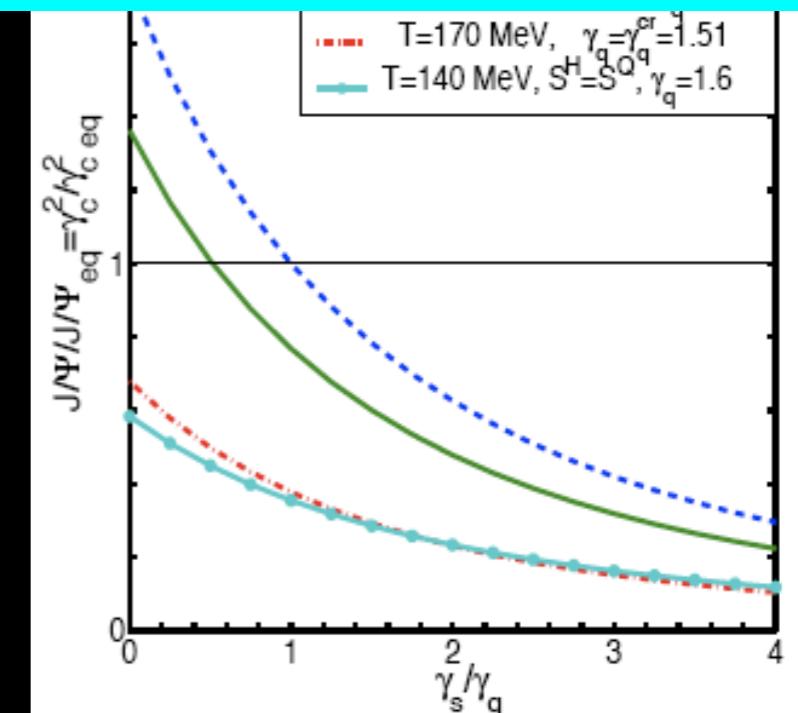
$D^-/D^+$	$\bar{D}_0/D_0$	$D^{*-}/D^{*+}$	$D_s^-/D_s^+$	$\bar{\Lambda}_c/\Lambda_c$	$D^+/D_0$	$D^{*+}/D_0$
1.00(0)	1.01(0)	1.01(0)	1.00(1)	1.00(1)	0.425(18)	0.387(15)
$D_s^+/D_0$	$\Lambda_c/D_0$	$\psi'/\psi$	$\eta_c/\psi$	$\chi_{c1}/\psi$	$\chi_{c2}/\psi$	
0.349(14)	0.163(16)	0.031(3)	0.617(14)	0.086(5)	0.110(8)	

Andronic et al.: ratios at equilibrium in the statistical model.

**Chemical composition studies will have implications on recombination models for quarkonium.**



strangeness oversaturation may lead to modifications in the ratios and to a suppression in the production of ccbar states.



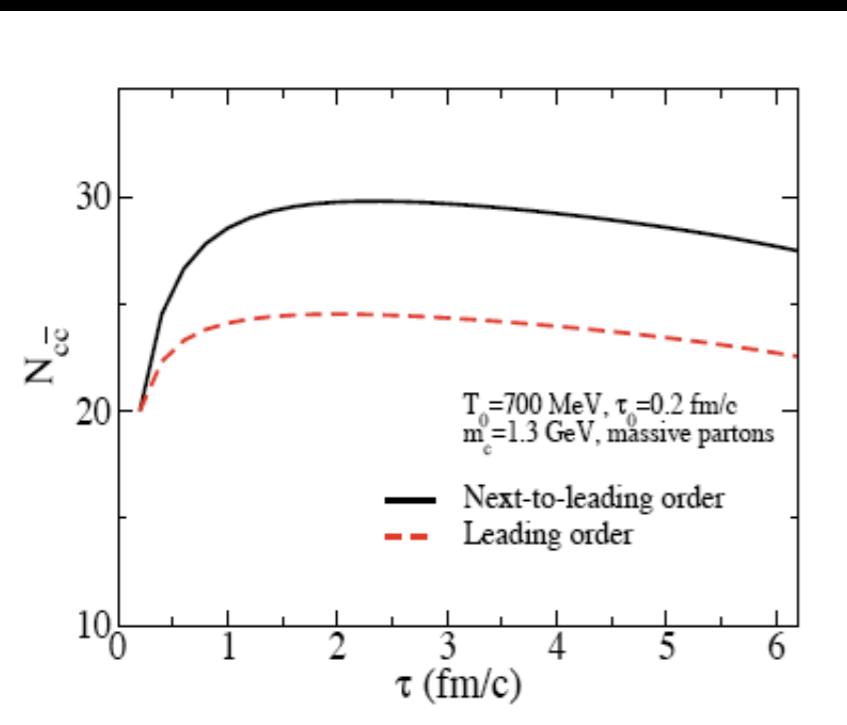
## B.2. HQ and quarkonium: others

Table 7: Upper bound on quarkonium dissociation temperatures.

state	$\chi_c$	$\psi'$	$J/\psi$	$\Upsilon'$	$\chi_b$	$\Upsilon$
$T_{\text{dis}}$	$\leq T_c$	$\leq T_c$	$1.2T_c$	$1.2T_c$	$1.3T_c$	$2T_c$

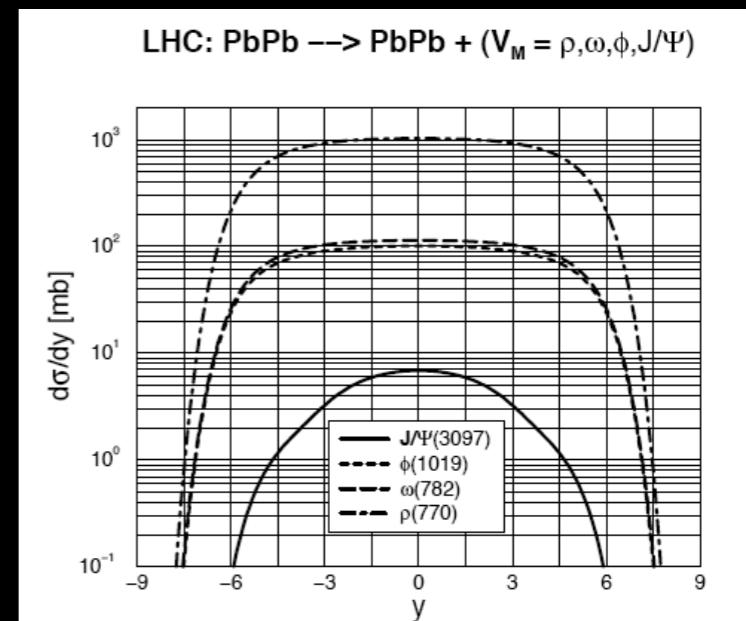
Mocsy et al.,  
dissociation T in  
potential models.

$$L_s(v, T) \simeq L_s(0, T) / \sqrt{\gamma} \rightarrow T_{\text{diss}}(v) \simeq T_{\text{diss}}(0) / \sqrt{\gamma}$$



Zhang et al., thermal  
charm at NLO.

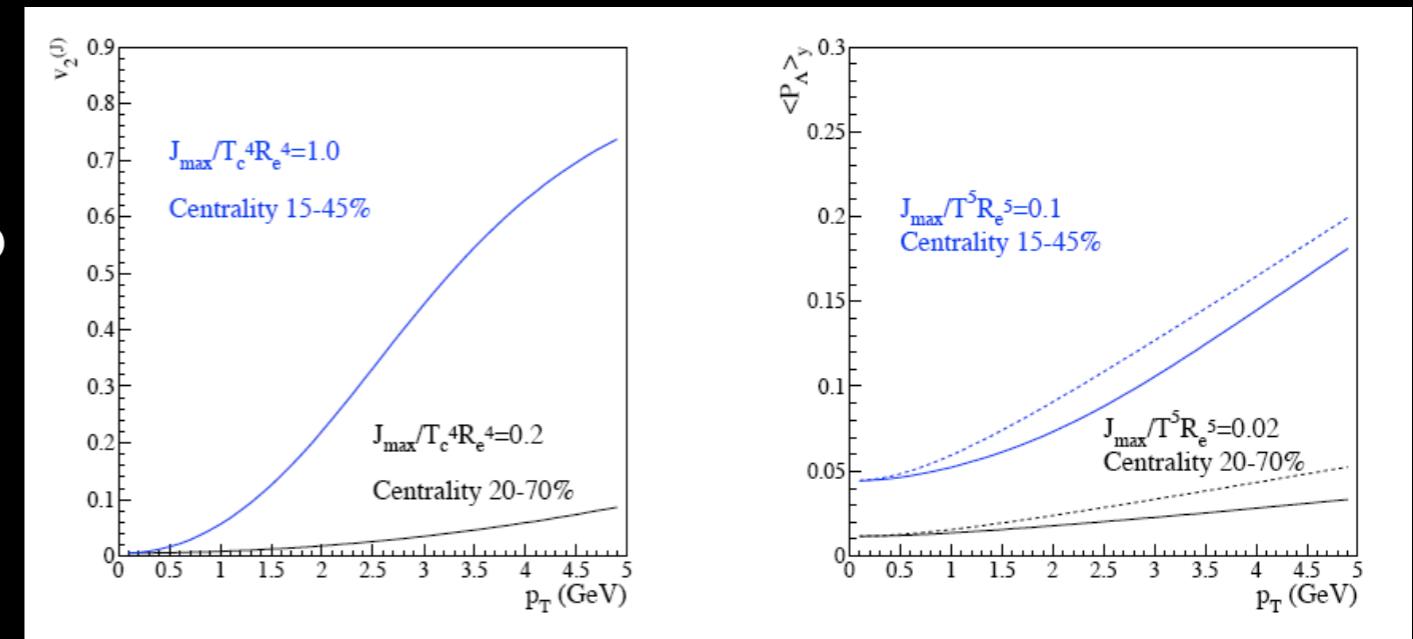
H. Liu et al., screening lengths  
through the medium, new  
suppression at larger p\_T if  
produced inside the medium



Goncalves et al., charm in UPC,  
test of production and npdf's.

# Backup III: others

Becattini et al.: in peripheral collisions, a highly spinning QGP may be formed → decreasing  $\mu_b$ , larger  $v_2$ , polarization.



Stocker et al.: creation of black holes in HIC due to low-scale extra dimensions: suppression of dijets, remnants.

Lee et al.: coalescence formation of charmed exotic, multi-quark hadrons, test of particle production mechanism.

Lokhtin et al.: exotic phenomena in HE CR, like alignment of tracks, may become visible at the LHC.