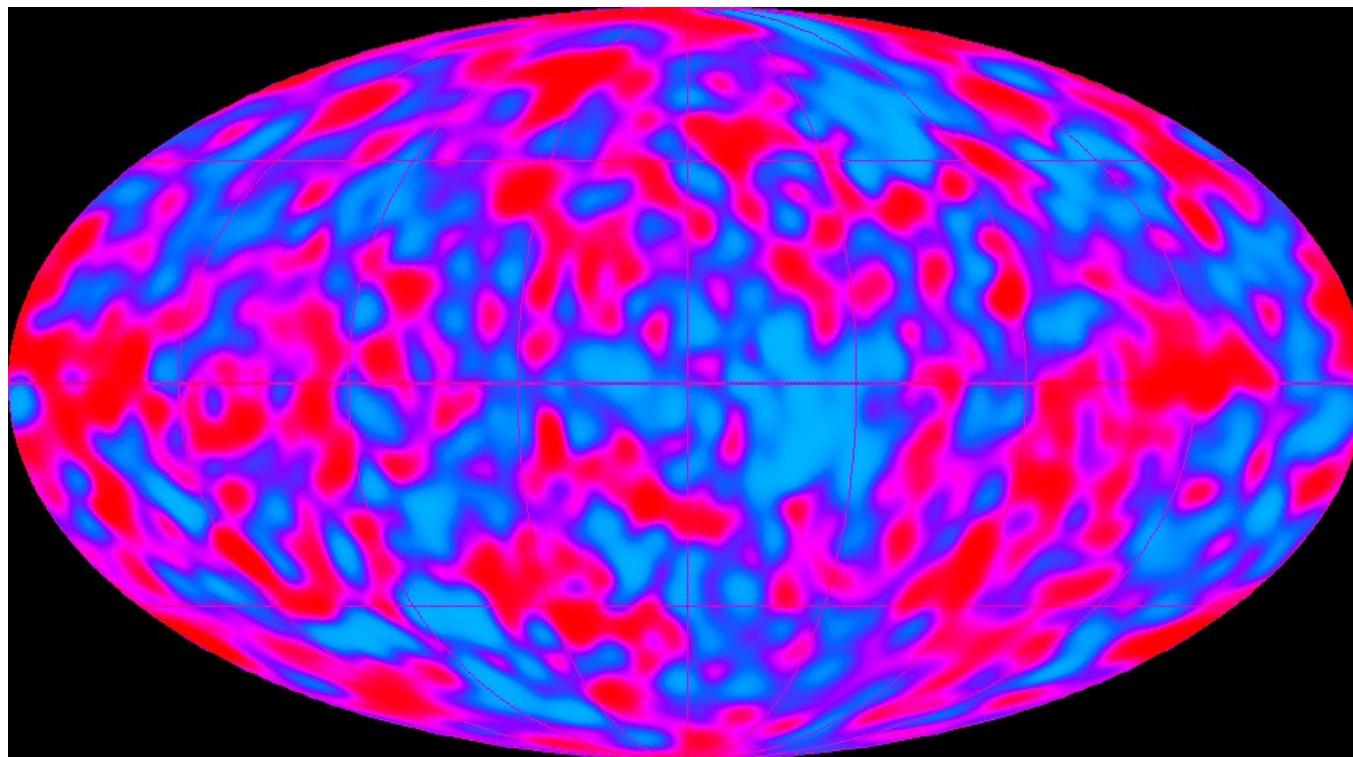
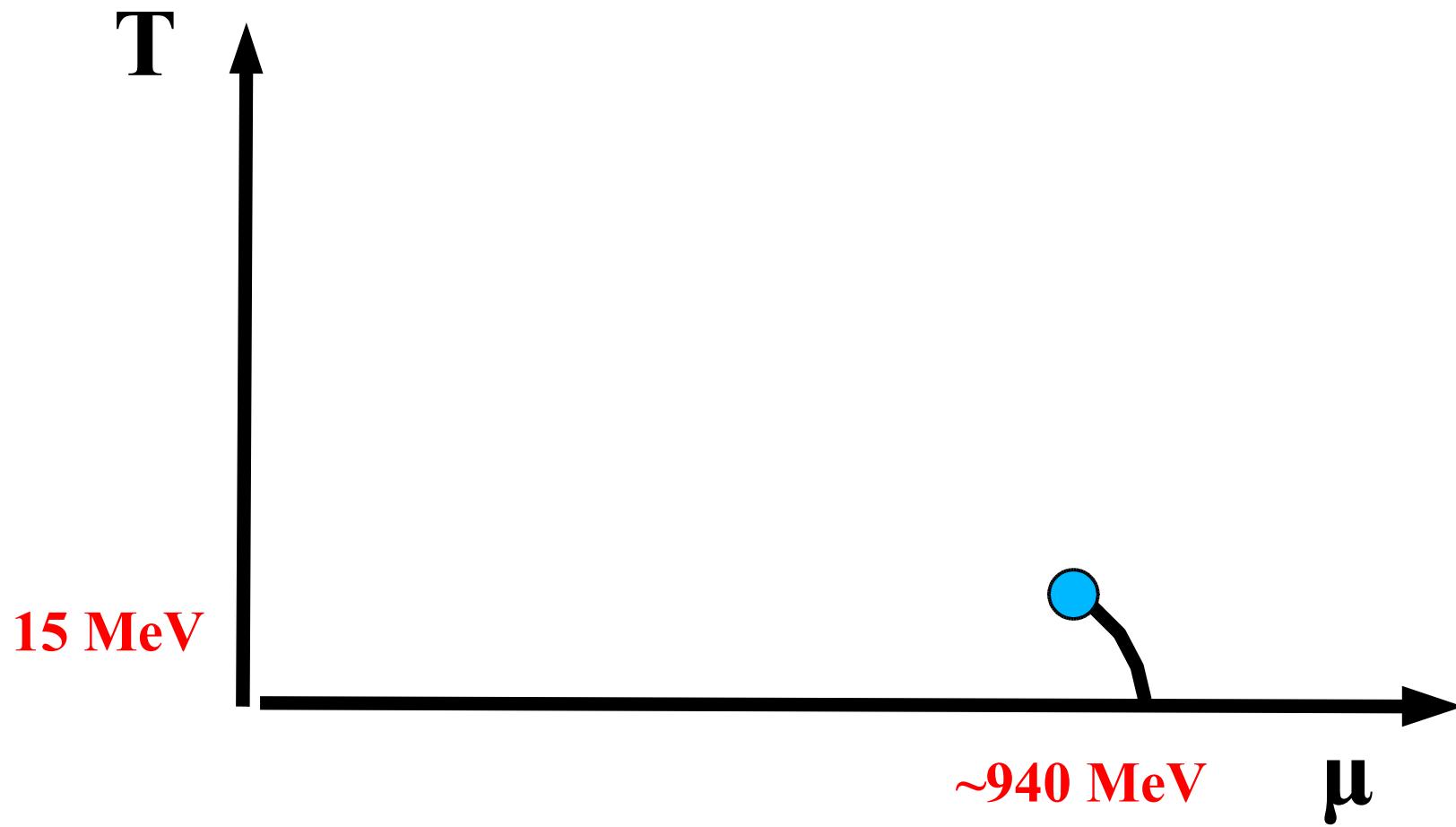


Fluctuations and Correlations

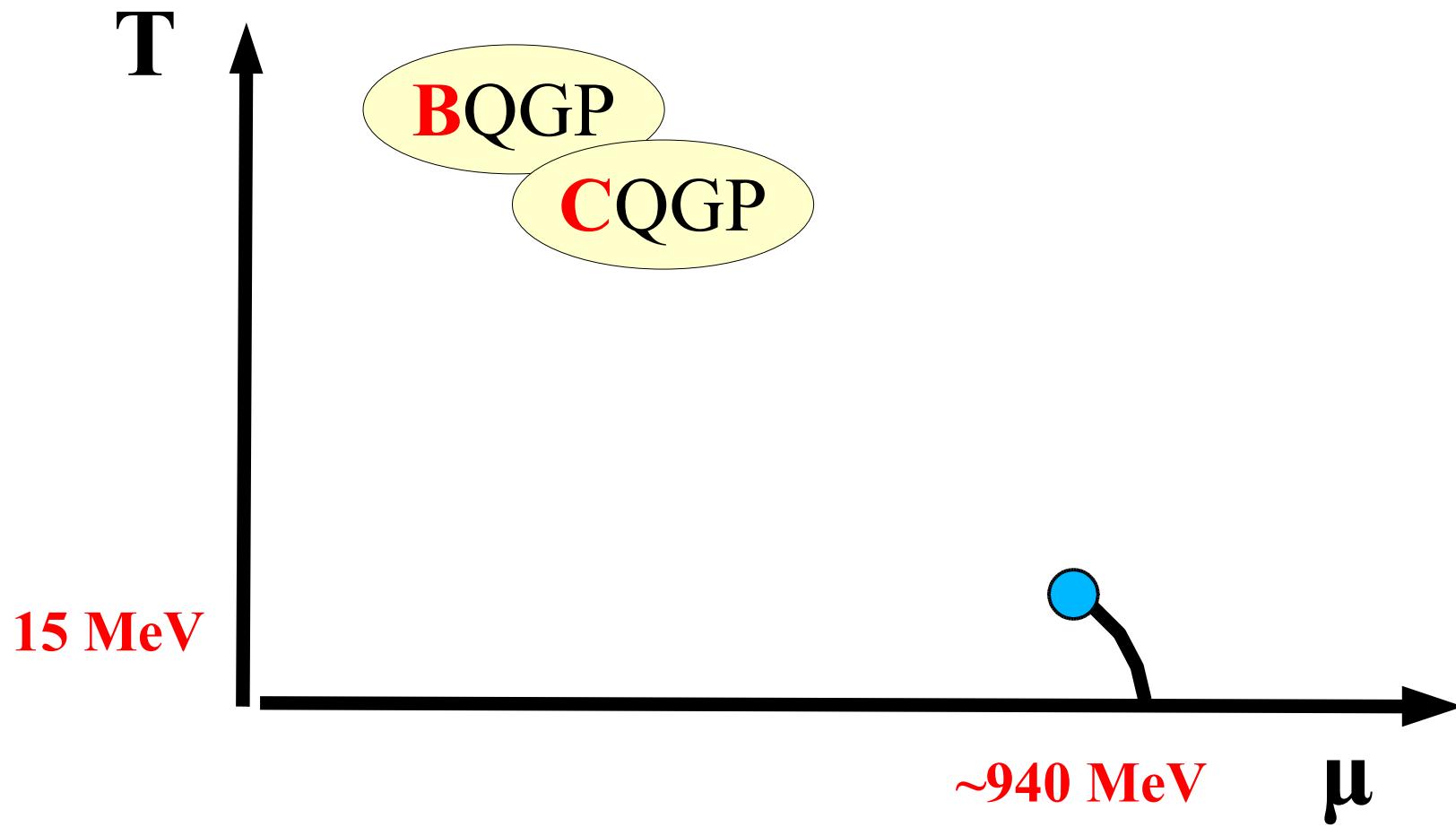
Status and future perspectives



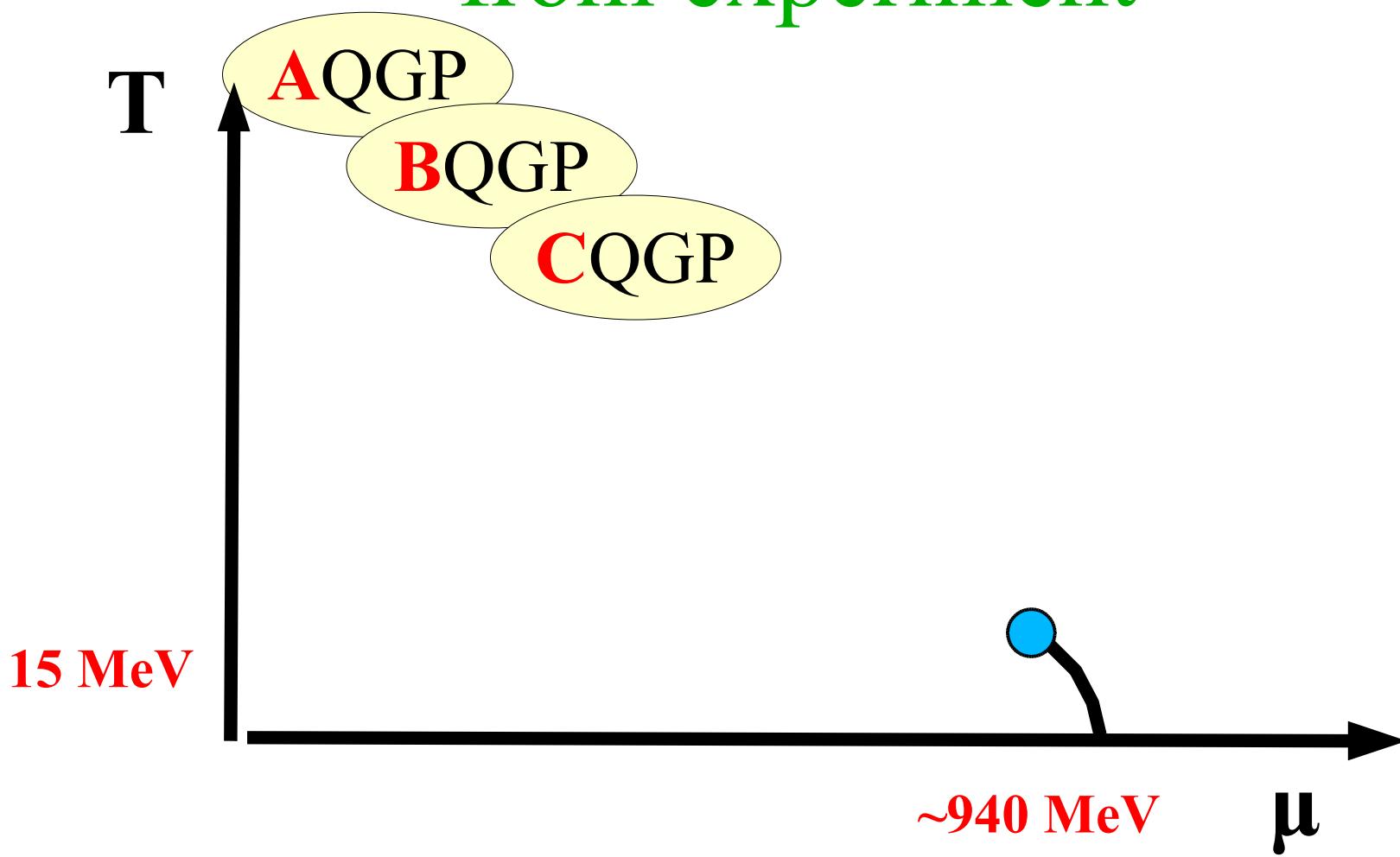
The QCD Phase Diagram from experiment



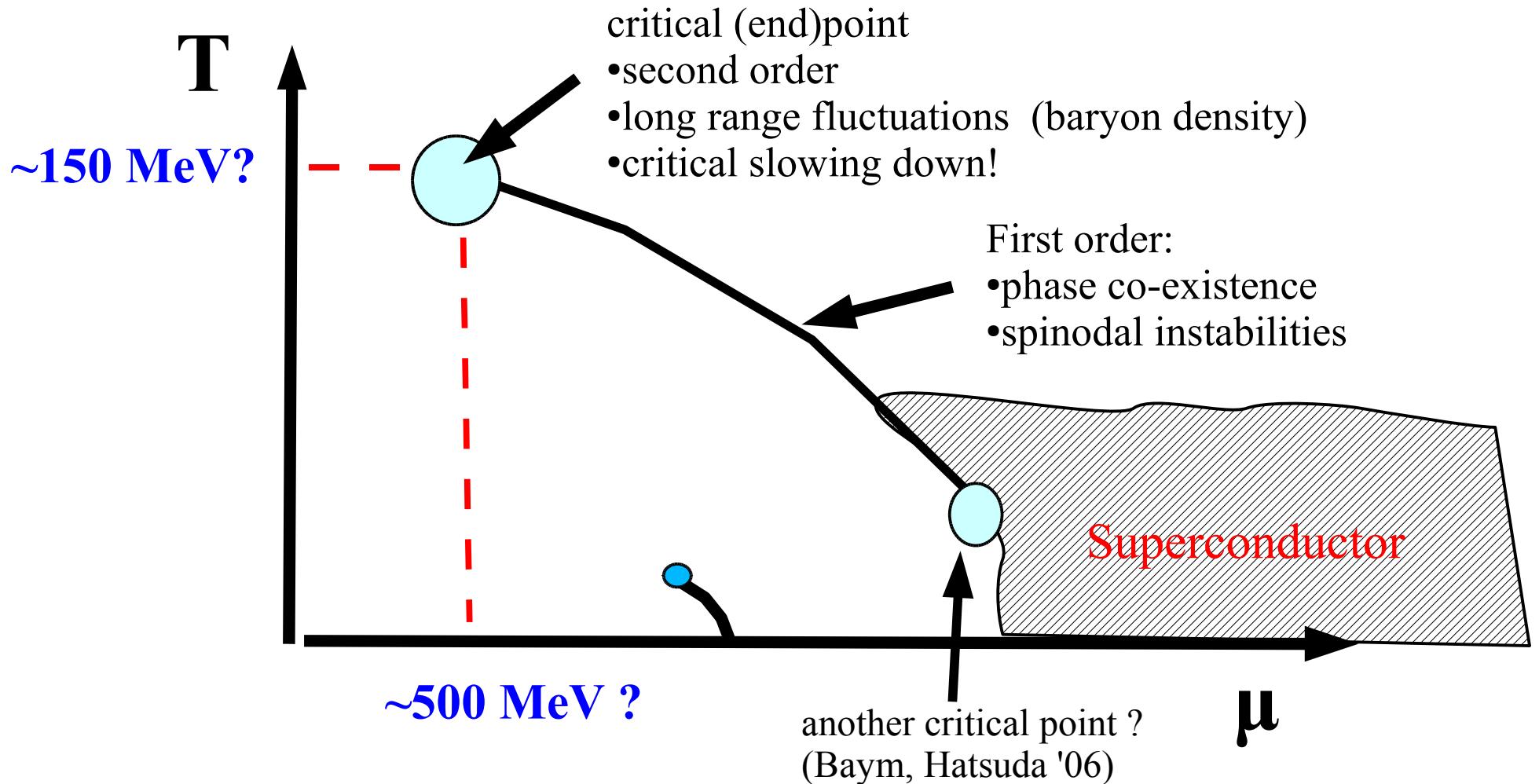
The QCD Phase Diagram from experiment



The QCD Phase Diagram from experiment



The QCD Phase Diagram (from a theorist's perspective)



N.B.: Critical point of water: $T_c = 647.096$ K, $p_c = 22.064$ MPa, $\rho_c = 322$ kg/m³

Fluctuations and Correlations in thermal system

e.g. Lattice QCD

$$Z = \text{Tr} [\exp(-\beta(H - \mu_Q Q - \mu_B B - \mu_S S))]$$

Mean : $\langle \alpha \rangle = T \frac{\partial}{\partial \mu_\alpha} \log(Z) = -\frac{\partial}{\partial \mu_\alpha} F$

Variance: $\langle (\delta \alpha)^2 \rangle = T^2 \frac{\partial^2}{\partial \mu_\alpha^2} \log(Z) = -T \frac{\partial^2}{\partial \mu_\alpha^2} F \quad \alpha, \beta = Q, B, S$

Co-Variance: $\langle (\delta \alpha)(\delta \beta) \rangle = T^2 \frac{\partial^2}{\partial \mu_\alpha \partial \mu_\beta} \log(Z) = -T \frac{\partial^2}{\partial \mu_\alpha \partial \mu_\beta} F$

Susceptibility: $\chi_{\alpha\beta} = -\frac{1}{V} \frac{\partial^2}{\partial \mu_\alpha \partial \mu_\beta} F = -\frac{1}{V} \frac{\partial}{\partial \mu_\alpha} \langle \beta \rangle$

Fluctuations and Correlations in thermal system

e.g. Lattice QCD

$$Z = \text{Tr} [\exp(-\beta(H - \mu_Q Q - \mu_B B - \mu_S S))]$$

Susceptibility: $\chi_{\alpha\beta} = -\frac{1}{V} \frac{\partial^2}{\partial \mu_\alpha \partial \mu_\beta} F = -\frac{1}{V} \frac{\partial}{\partial \mu_\alpha} \langle \beta \rangle$ $\alpha, \beta = Q, B, S$

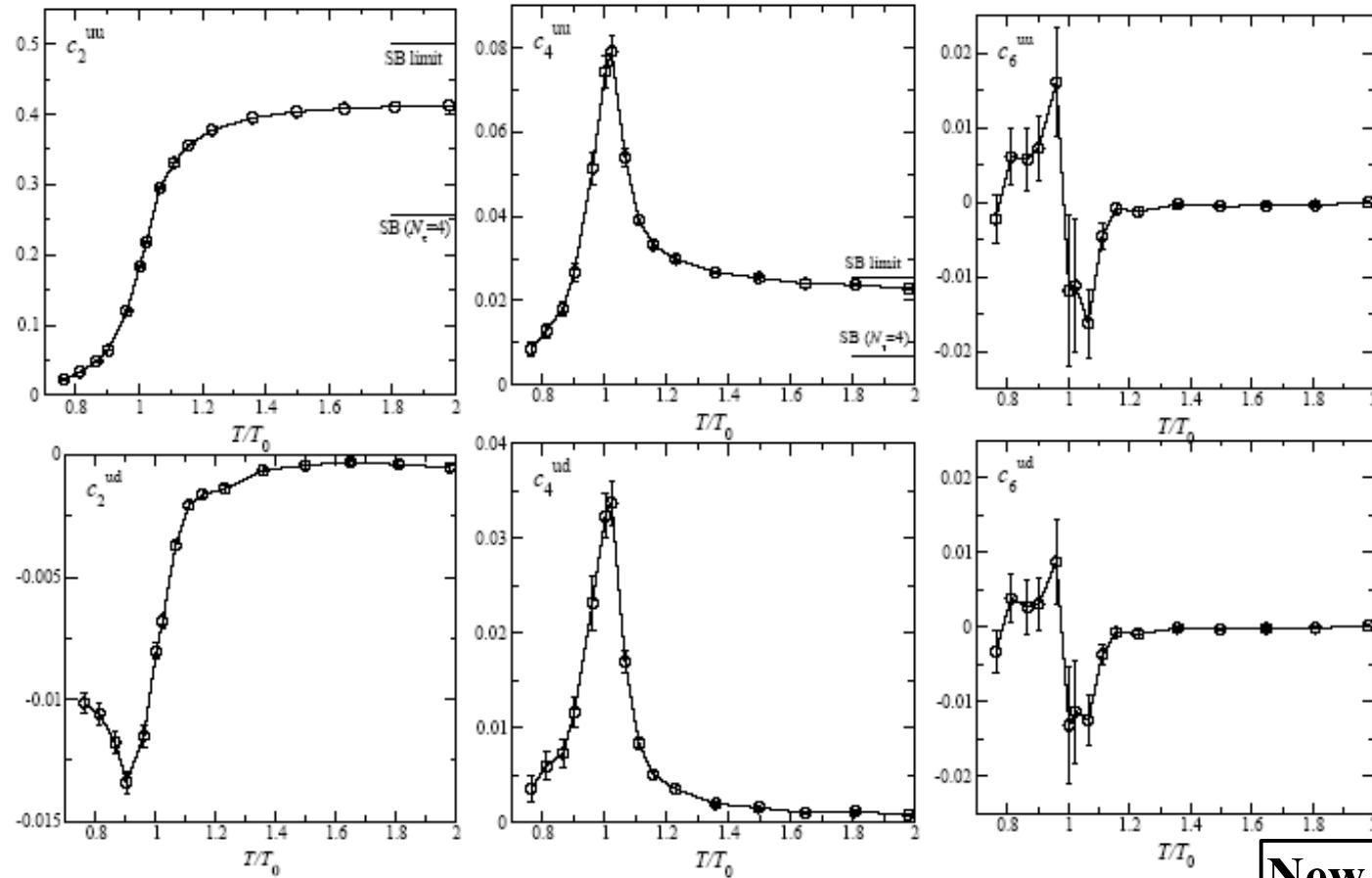
More general: $\chi_{\alpha, \beta} \sim \frac{1}{V} \int d^3x d^3y \langle \rho_\alpha(x) \rho_\beta(y) \rangle_{connected} = \int d^3r C_{\alpha, \beta}(r) \sim \xi_{\alpha, \beta}^2$

Correlation Function:

$$C_{\alpha, \beta}(r) = \langle \rho_\alpha(r) \rho_\beta(0) \rangle_{connected} \sim \frac{e^{(-r/\xi_{\alpha, \beta})}}{r} \quad \xi_{\alpha, \beta} = \text{correlation length}$$

Lattice-QCD susceptibilities

$$\frac{\chi(T, \mu_q)}{T^2} = 2c_2 + 12c_4\left(\frac{\mu_q}{T}\right)^2 + 30c_6\left(\frac{\mu_q}{T}\right)^4 + \dots$$



Alton et al, PRD 66 074507 (2002)

Rule of thumb:

$$c_n \sim \langle X^n \rangle \\ X = B, Q, S, \dots$$

**New results: C. Schmidt,
this afternoon, session X**

Susceptibilities and Phasetransitions

$$Z = Tr[\exp(-\beta(H - \mu N))]$$

Susceptibility: $\chi \sim \frac{1}{V} \frac{\partial^2}{\partial \mu^2} \log(Z) = \frac{1}{V} (\langle N^2 \rangle - \langle N \rangle^2)$

Poisson: $\chi \sim \frac{\langle N \rangle}{V}$ independent of volume $\rightarrow \langle (\delta N)^2 \rangle = N \sim V$

In general: $\chi \sim \frac{1}{V} \int d^3x d^3y \langle \rho(x) \rho(y) \rangle_{connected} = \int d^3r \langle \rho(r) \rho(0) \rangle_{connected} \sim \xi^2$

$$\langle \rho(r) \rho(0) \rangle_{connected} \sim \frac{e^{(-r/\xi)}}{r} \quad \xi = correlation\ length$$

Cross-over: $\xi = const \rightarrow \chi = const \rightarrow \langle (\delta N)^2 \rangle \sim V$

Second Order: $\xi \sim V^{(1/3)} \rightarrow \chi \sim V^{(2/3)} \rightarrow \langle (\delta N)^2 \rangle \sim V^{(5/3)}$

First Order: $\langle \rho(r) \rho(0) \rangle = const \rightarrow \chi \sim V \rightarrow \langle (\delta N)^2 \rangle \sim V^2$

Susceptibilities and Observables

Susceptibility: $\chi \sim \frac{1}{V} \frac{\partial^2}{\partial \mu^2} \log(Z) = \frac{1}{V} (\langle N^2 \rangle - \langle N \rangle^2)$

Fluctuations of some sort!

Cross-over: $\xi = \text{const} \rightarrow \chi = \text{const} \rightarrow \langle (\delta N)^2 \rangle \sim V$

Second Order: $\xi \sim V^{(1/3)} \rightarrow \chi \sim V^{(2/3)} \rightarrow \langle (\delta N)^2 \rangle \sim V^{(5/3)}$

First Order: $\langle \rho(r) \rho(0) \rangle = \text{const} \rightarrow \chi \sim V \rightarrow \langle (\delta N)^2 \rangle \sim V^2$

Since fluctuations diverge at phase transition **any** sort will do!

System size dependence!

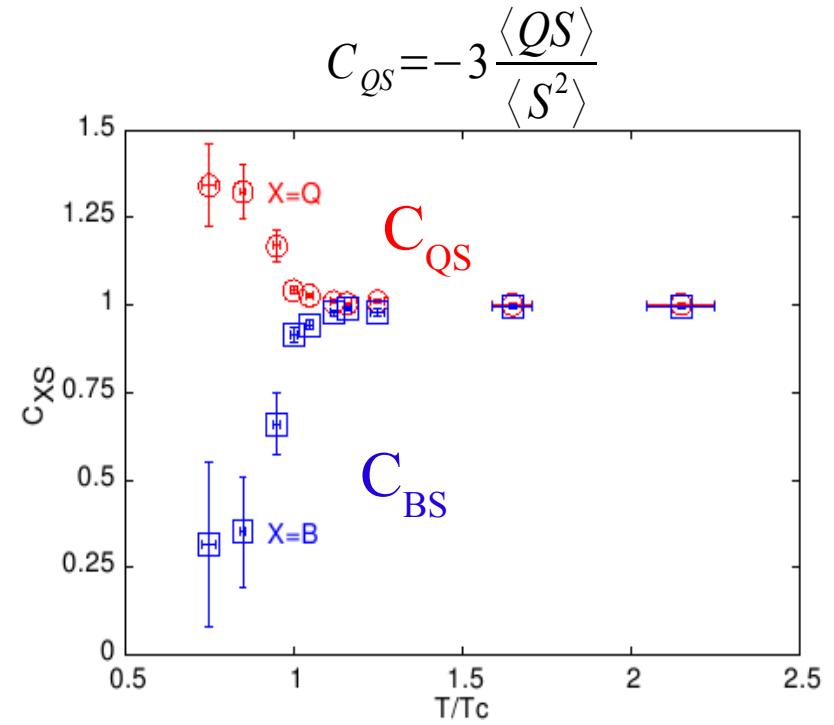
Note of caution: Co-variances also diverge; trigger?

“Charge” (co)-variances

$\langle Q^2 \rangle$, $\langle BS \rangle$, $\langle QS \rangle$

$$C_{BS} = -3 \frac{\langle BS \rangle}{\langle S^2 \rangle}$$

- Net Charge fluctuations (QGP)
- Baryon-Strangeness correlations (QGP)
- Baryon number fluctuations (CP)
- Balance Functions
-

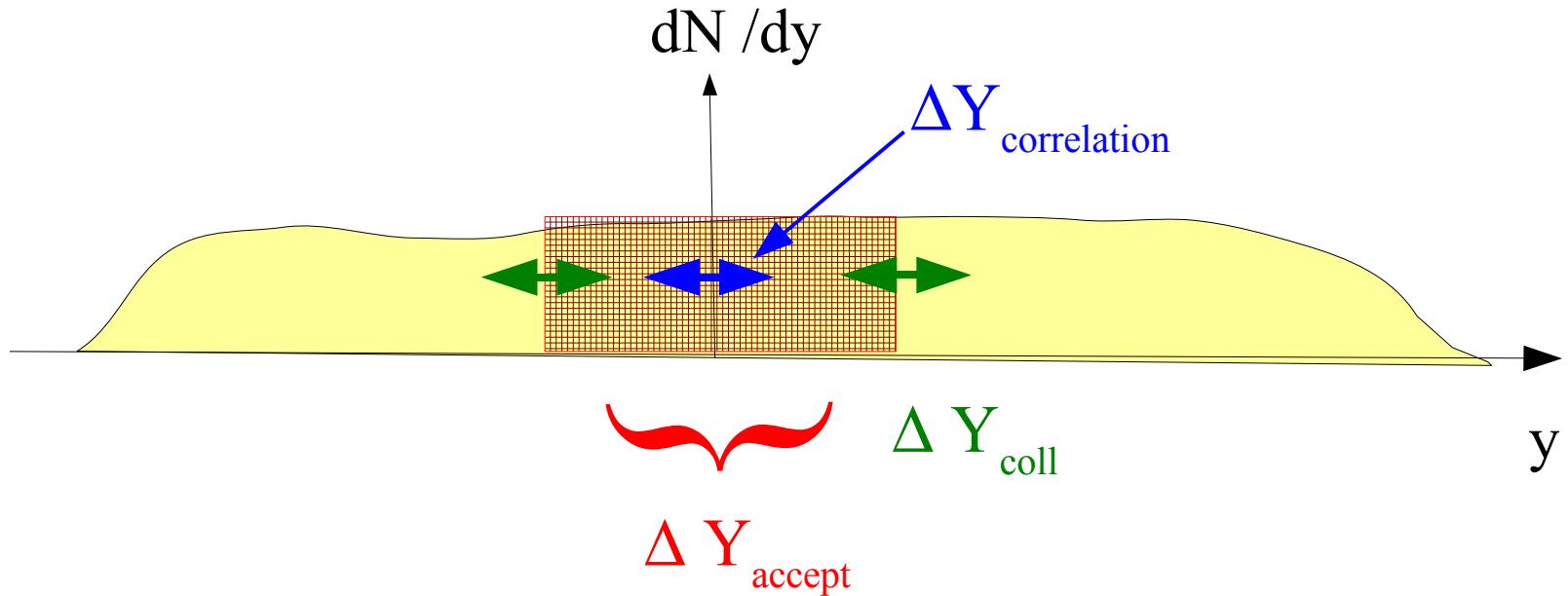


Fluctuations of conserved charges can survive hadronization!

S. Jeon, V.K; Stephanov, Shuryak (2000)

Gavai,Gupta, hep-lat/0510044
V.K, Majumder, Randrup PRL95:182301,2005
Majumder, Mueller Phys.Rev.C74:054901,2006

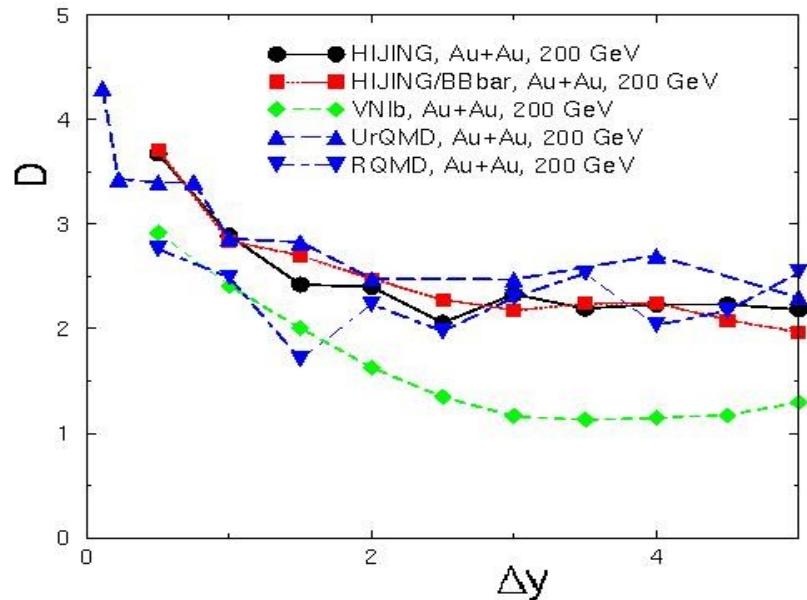
“Charge” fluctuations



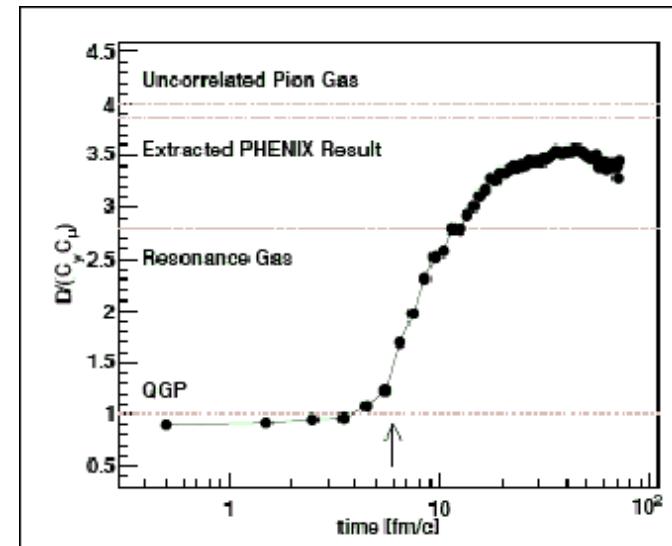
Condition for “charge” fluctuations:

- 1) $\Delta Y_{\text{correlation}} \ll \Delta Y_{\text{accept}}$ **(catch the physics)**
- 3) $\Delta Y_{\text{total}} \gg \Delta Y_{\text{accept}} \gg \Delta Y_{\text{coll}}$ **(keep the physics)**

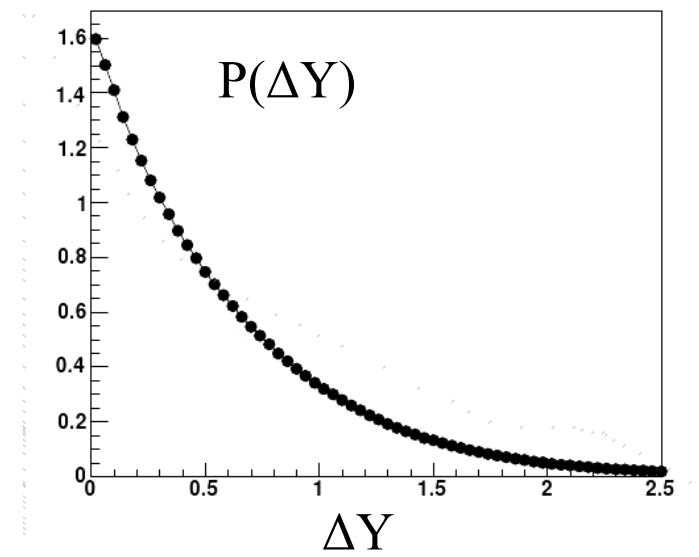
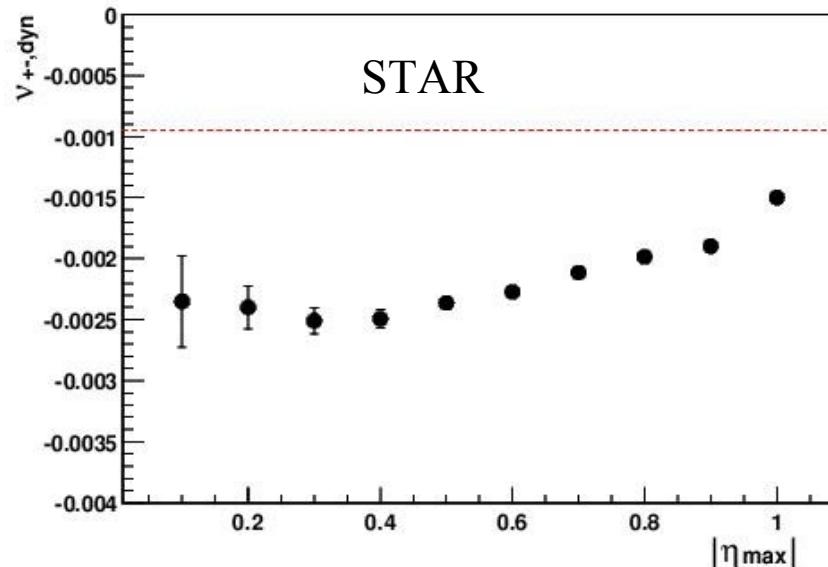
Net-Charge fluctuations



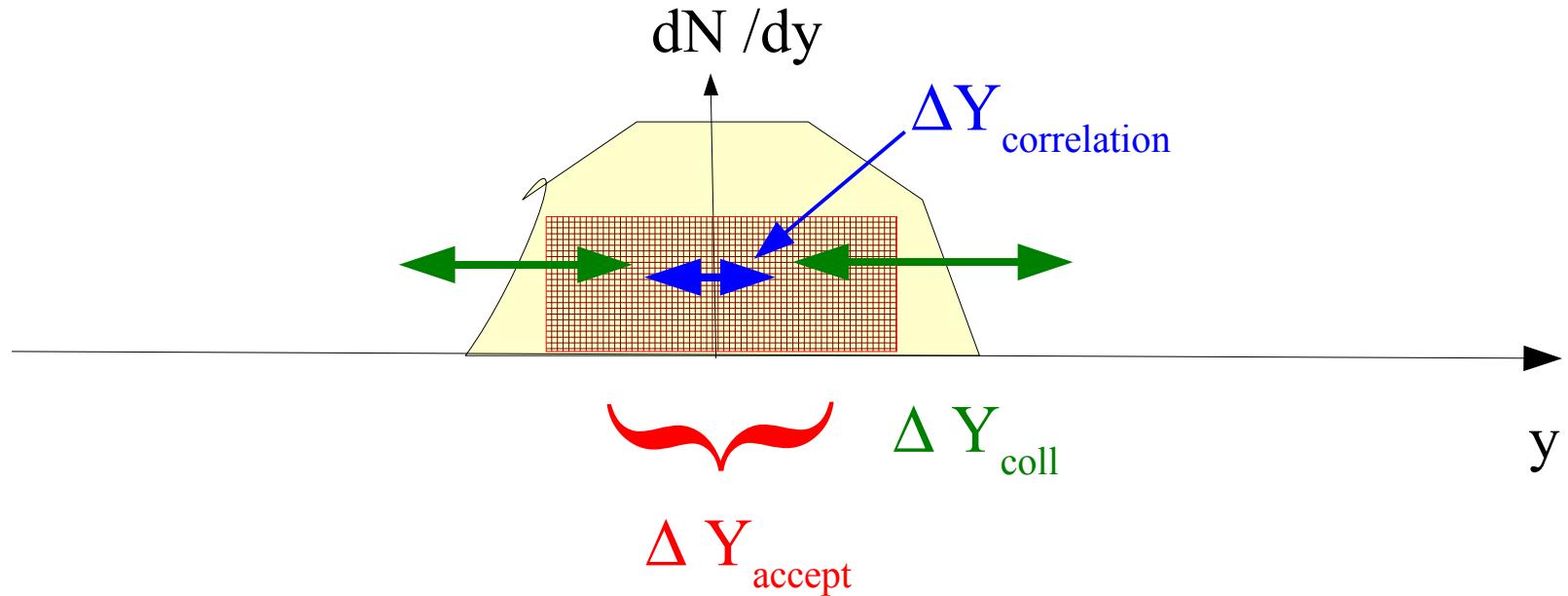
Gale et al.



Bleicher et al



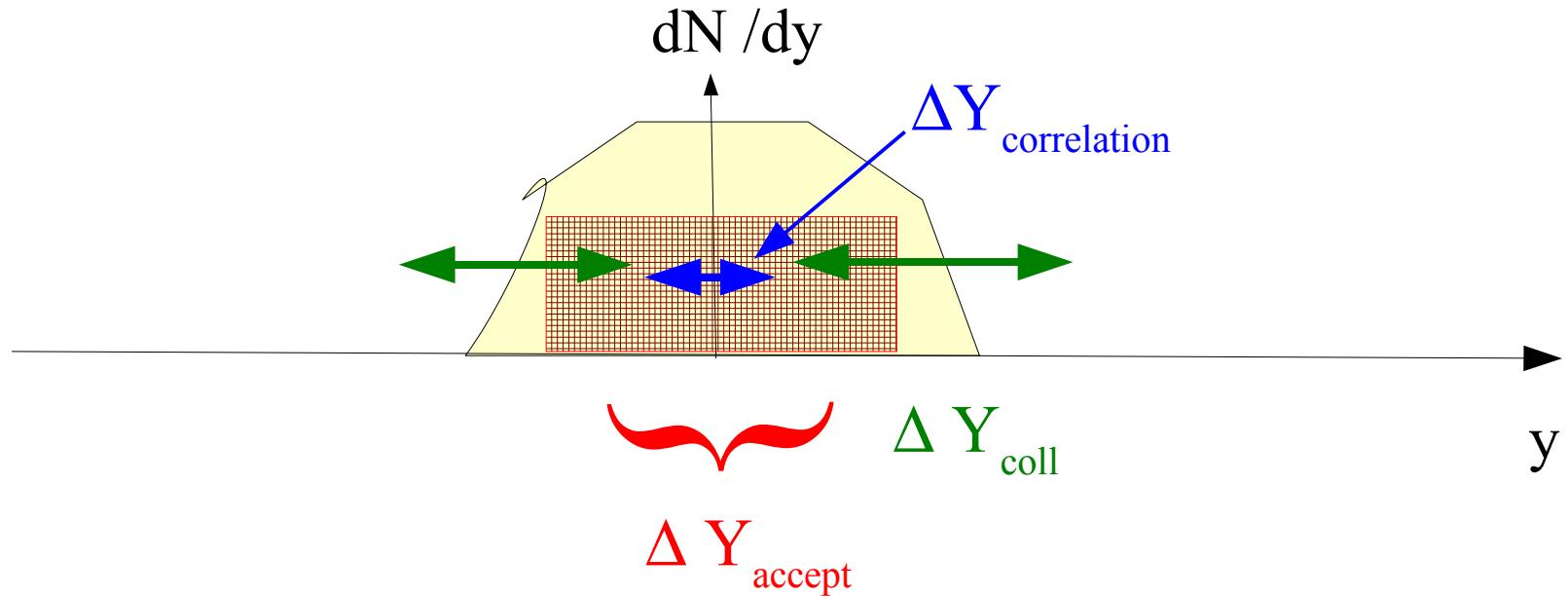
“Charge” fluctuations at SPS and below



Condition for “charge” fluctuations:

- 1) $\Delta Y_{\text{correlation}} \ll \Delta Y_{\text{accept}}$ **(catch the physics)**
- 3) $\Delta Y_{\text{total}} \gg \Delta Y_{\text{accept}} \gg \Delta Y_{\text{coll}}$ **(keep the physics)**

“Charge” fluctuations at SPS and below



Condition for “charge” fluctuations:

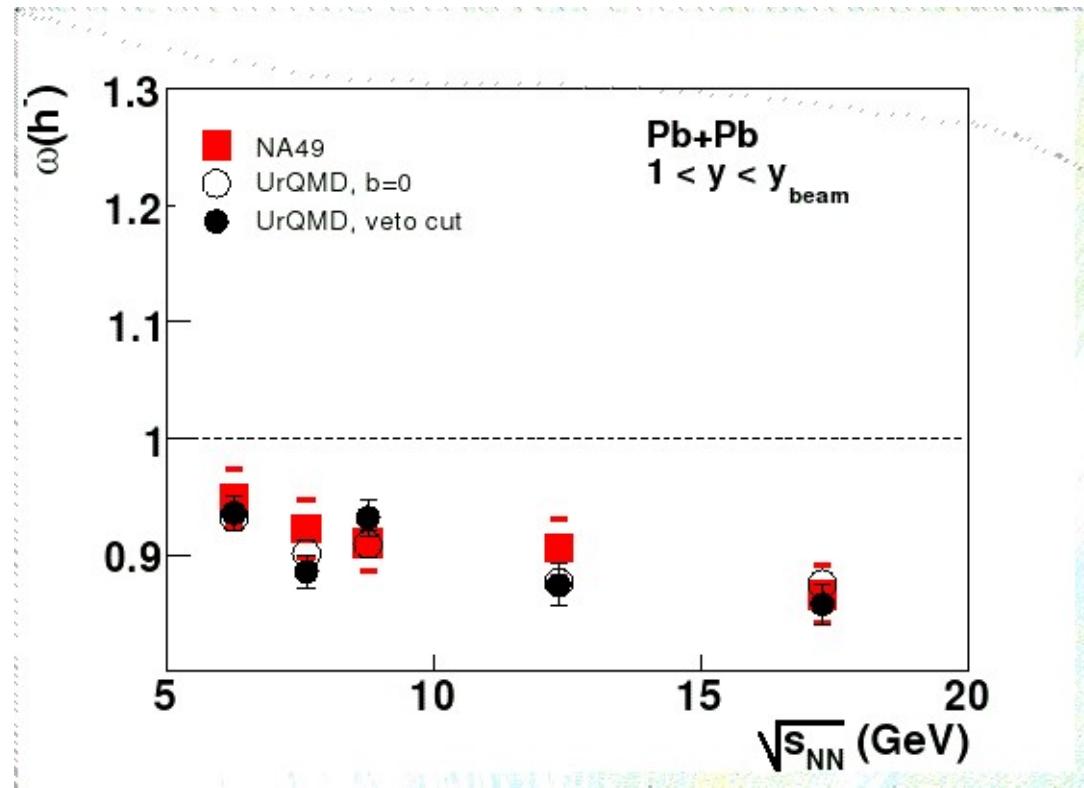
1) $\Delta Y_{\text{correlation}} \ll \Delta Y_{\text{accept}}$ **(catch the physics)**

3) $\Delta Y_{\text{total}} \gg \Delta Y_{\text{accept}} \gg \Delta Y_{\text{coll}}$ **(keep the physics)**

E-by-E observables

- Net-Charge fluctuations
 - Resonance gas at RHIC
 - no sensitivity at SPS and below
- Multiplicity fluctuations
 - interesting centrality dependence at top SPS energies
 - **No** energy dependence

Multiplicity Fluctuations

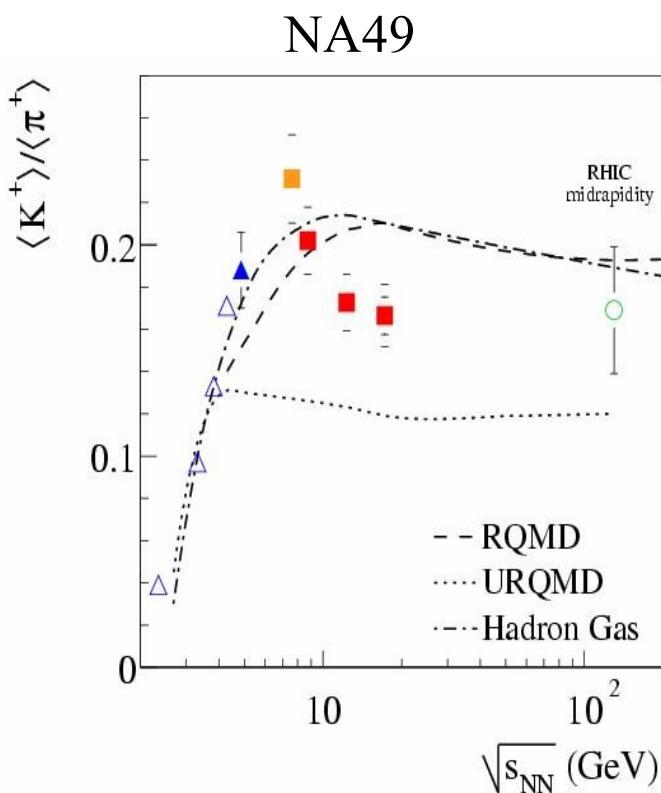
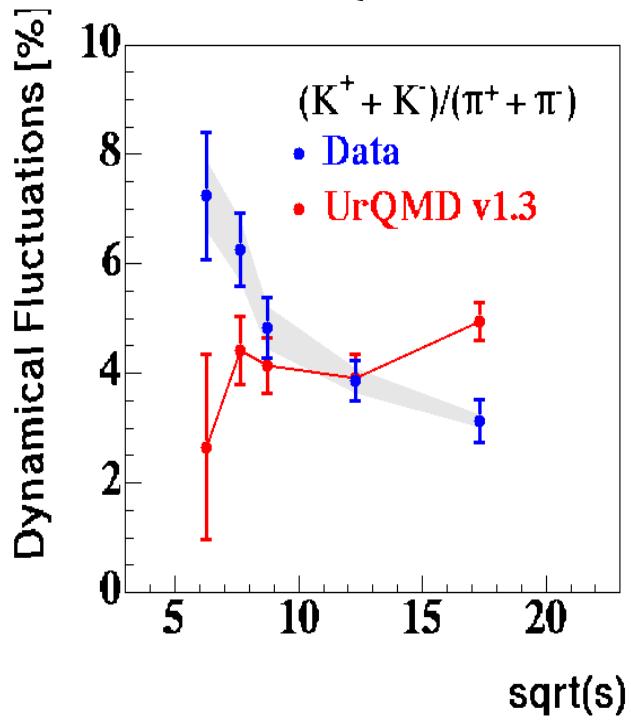


E-by-E observables

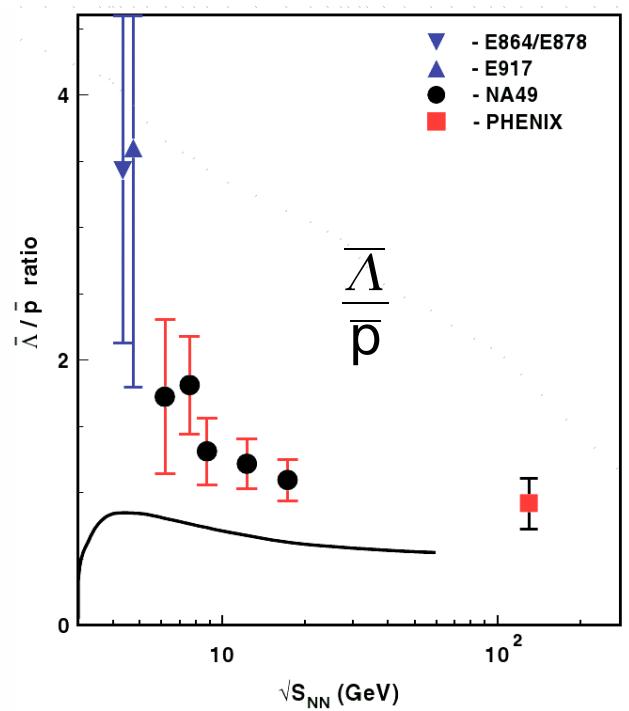
- Net-Charge fluctuations
 - Resonance gas at RHIC
 - no sensitivity at SPS
- Multiplicity fluctuations
 - interesting centrality dependence at top SPS energies
 - **No** energy dependence
- Transverse momentum fluctuations
 - some signal at SPS & RHIC (mostly “jets”)
 - **No** energy dependence
- Ratio (K/π) fluctuations
 - statistical at top SPS, possible signal at low SPS

Strange things...

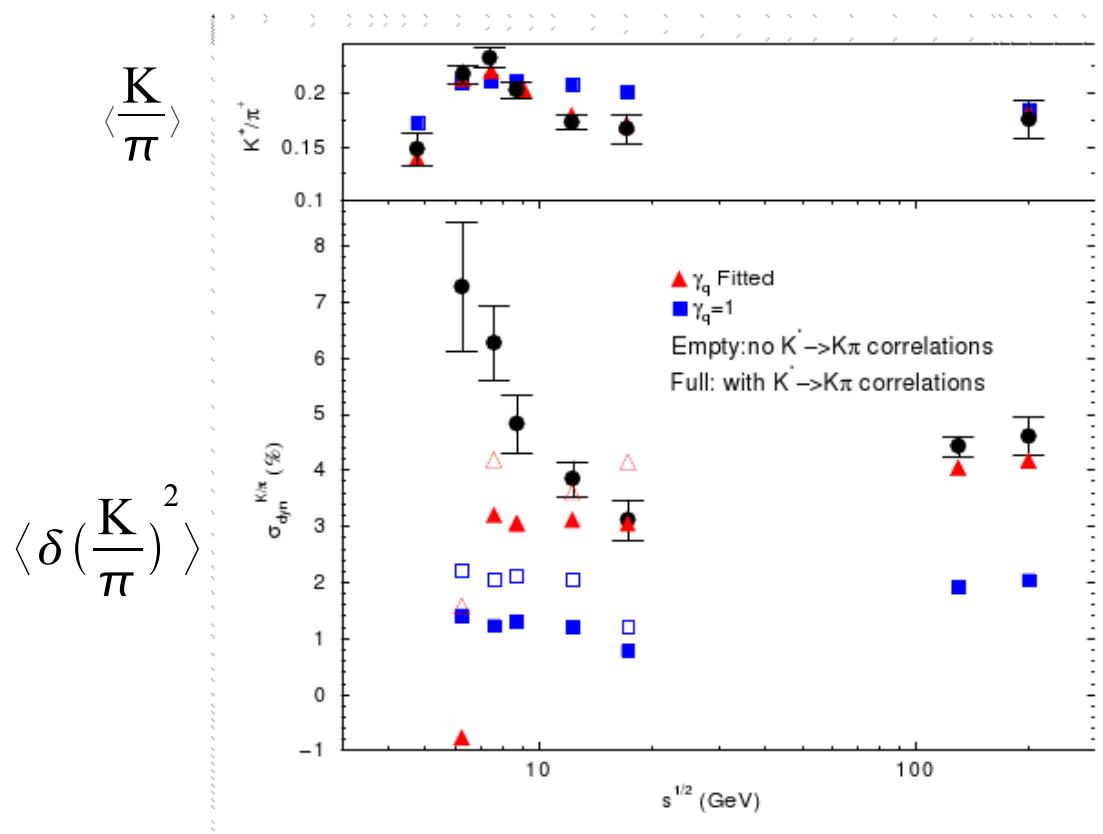
C. Roland, QM04, NA49



NA49, PRC73, 044910 (2006)

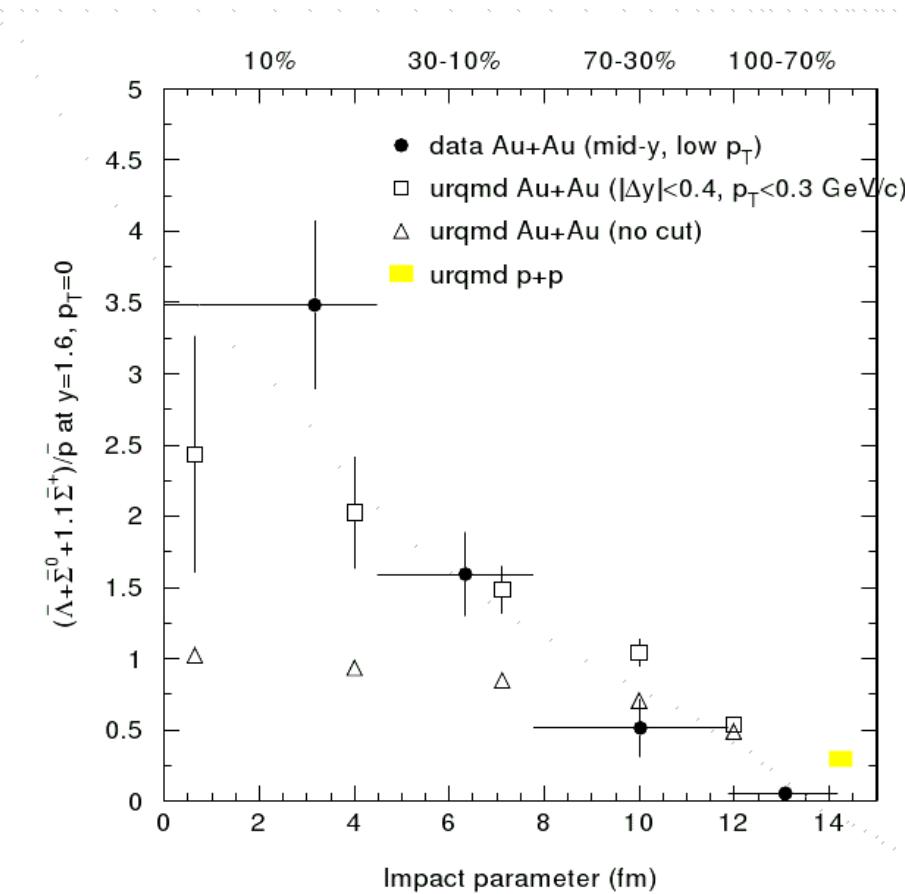


Hadron gas predictions



G. Torrieri, QM2006

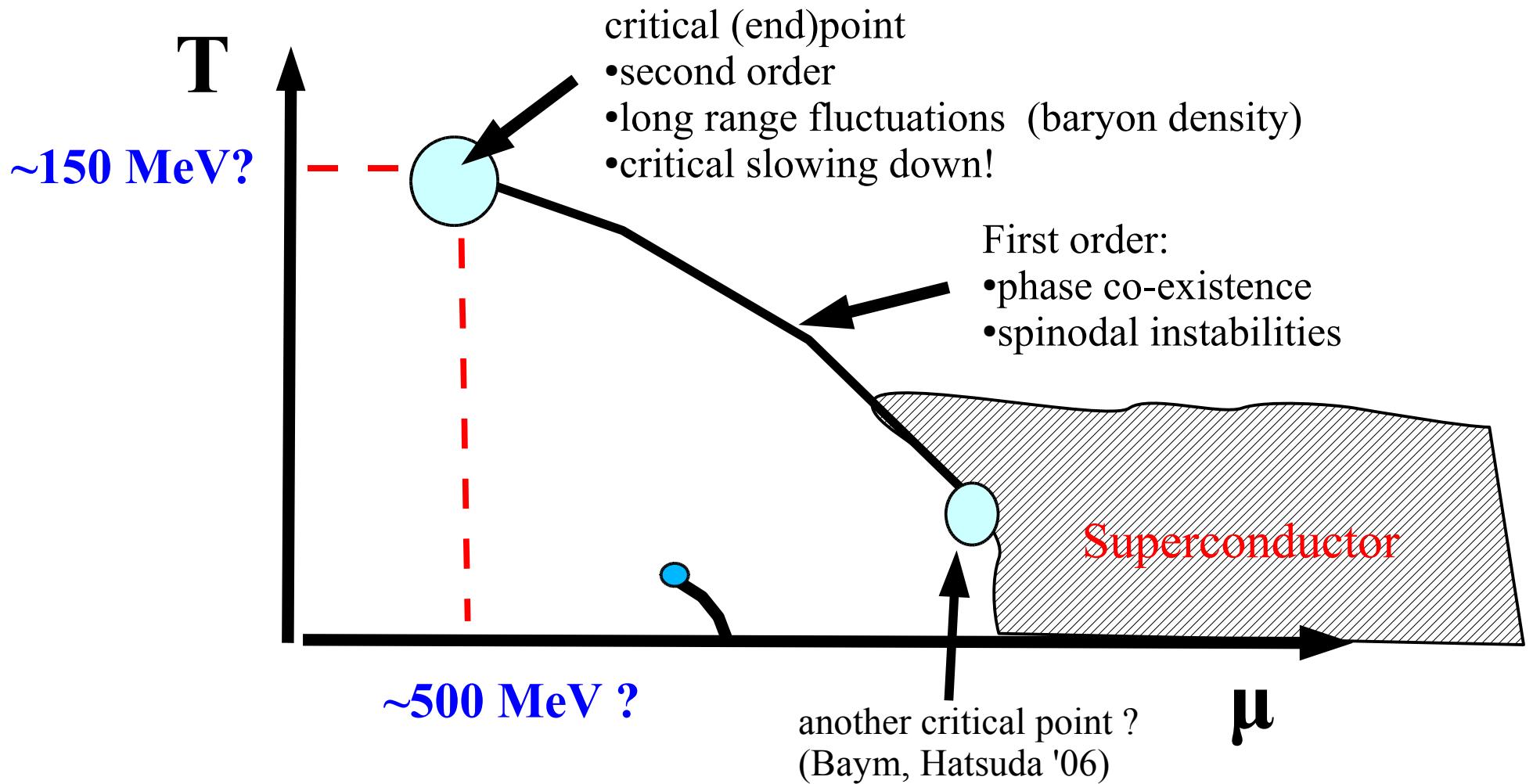
UrQMD and Lambda-bar / p-bar



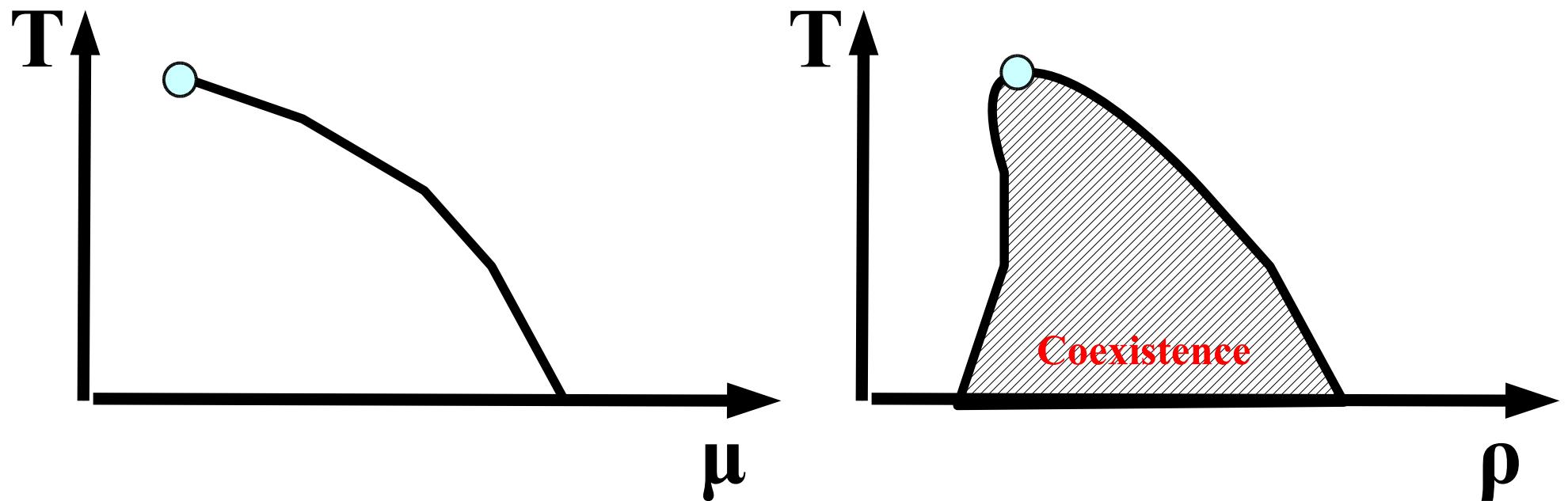
F. Wang
nucl-ex/0010002

Strong enhancement
mostly an effect of
acceptance cut !?

The QCD Phase Diagram (from a theorist's perspective)



Order Parameter



→ Baryon density is a good order parameter
density fluctuations are a good observable (theoretically...)
Baryon Number fluctuations also good in principle, but
same problem as Charge fluctuations at low Energy

Nambu model

(Fuji et al, hep-ph/0401028,0403039)

Sigma remains massive at CP; CP driven by **spacelike** p-h excitations

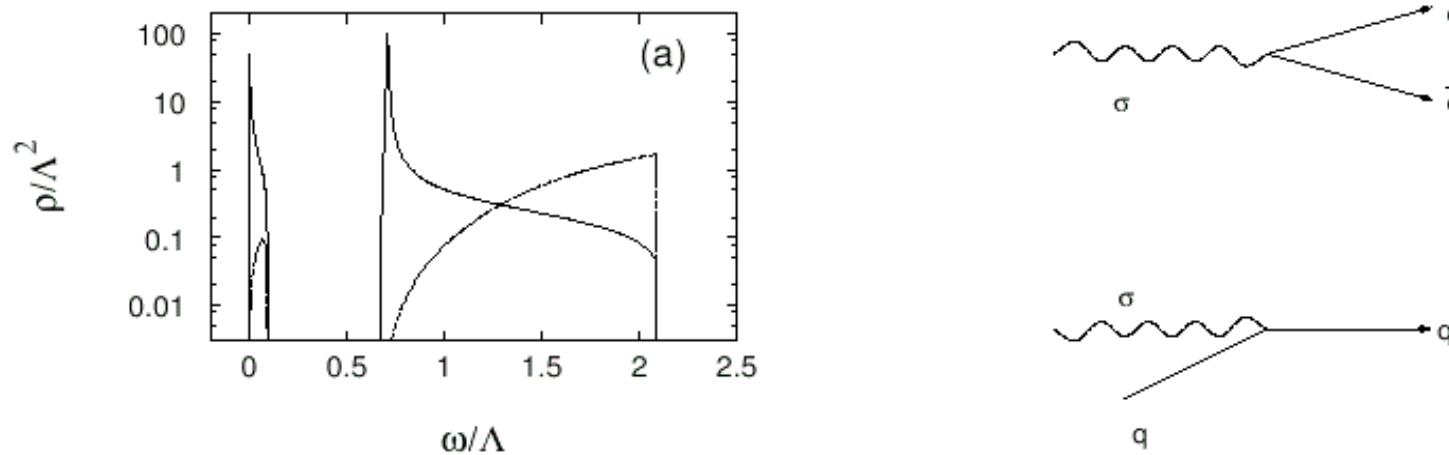


Fig. 2. (a) Spectral function in the scalar channel (solid) with $|q|/\Lambda = 0.1$ at a CEP with $m/\Lambda = 0.01$. The free gas spectrum (dashed) is also shown for reference. (b) Typical processes contributing to the spectrum.

Nambu model p-h excitations

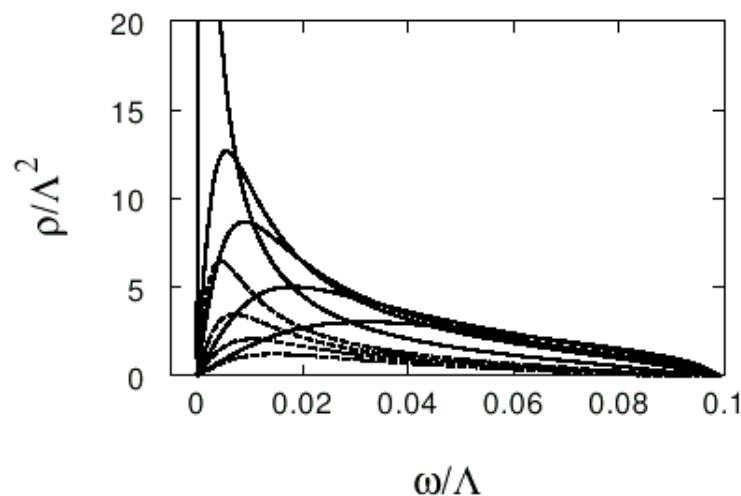
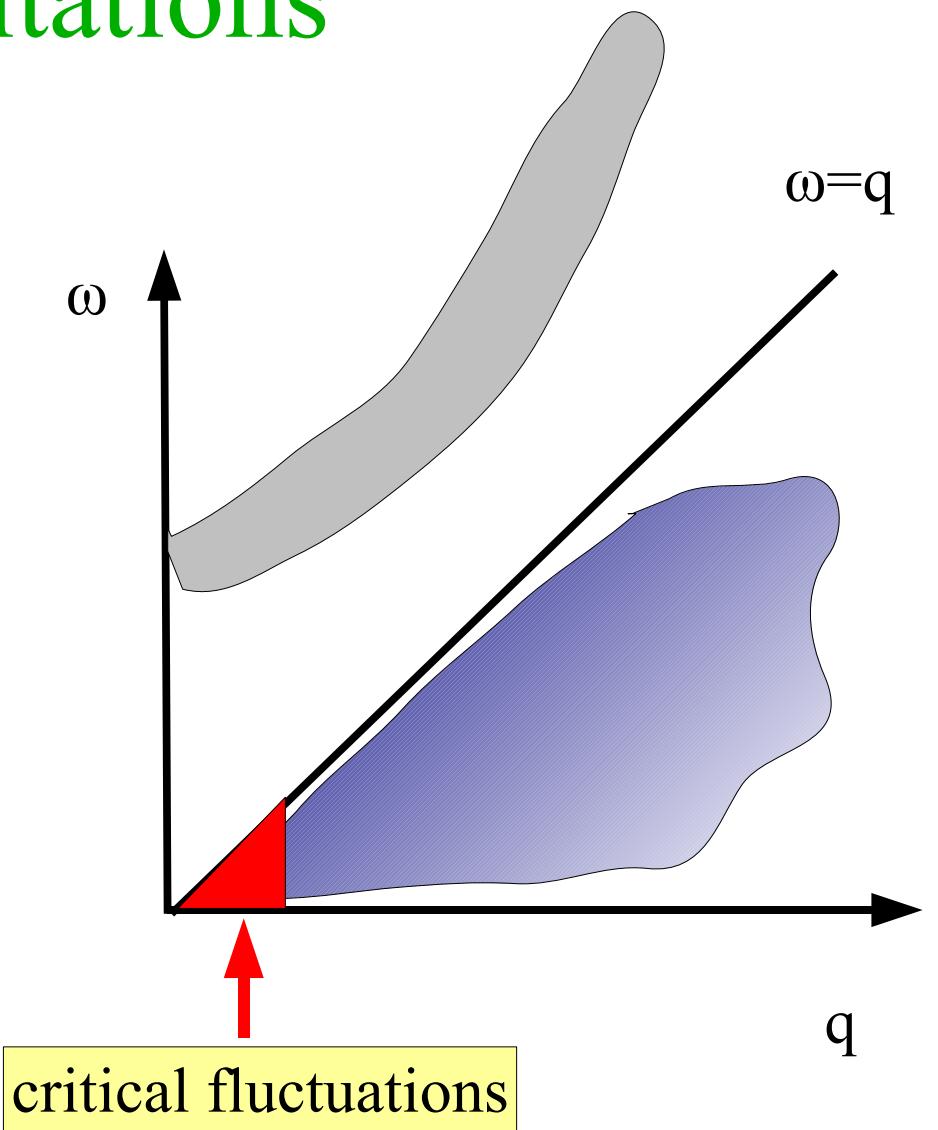


Fig. 3. Spectral function in the spacelike momentum region with $|q|/\Lambda = 0.1$, $T = T_c$ and $m/\Lambda = 0.01$ for several μ (see text).



QCD critical point

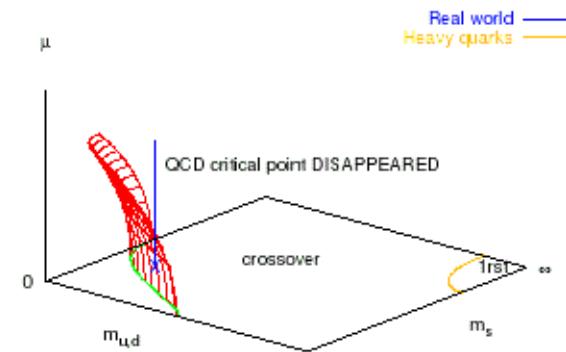
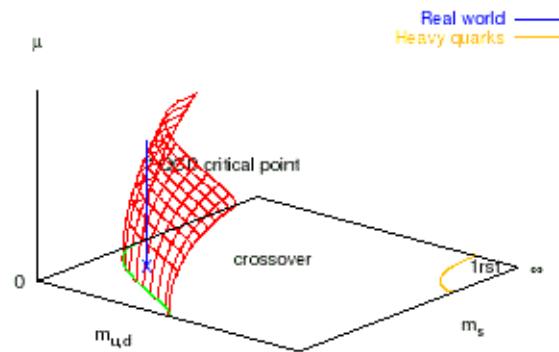
- Order parameter: baryon density or scalar density
 - Actually it is a superposition
- Both scalar (chiral) an quark number susceptibilities diverge
- **Screening** (“space like”) masses vanish (“omega”, “sigma”)
 - not accessible by (time-like) dileptons
- Is it related to chiral transition at $m_q=0$?

Lattice and the critical point

Forcrand, Philipsen

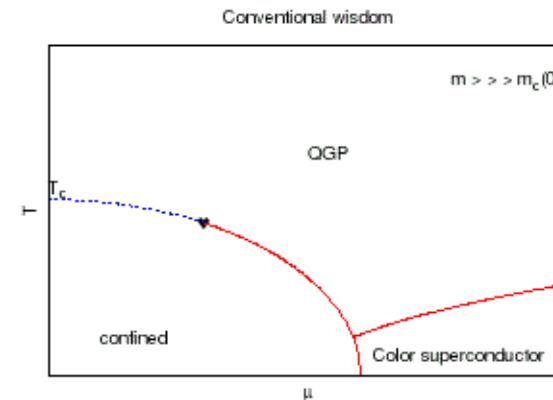
A non-standard scenario: no critical point?

$$\text{sign of } c_1 = \frac{dm_c(\mu)}{d\mu^2}|_{\mu=0}$$



If Phase transition is not chiral but
"liquid gas" and "anchored at low T high μ
the "conventional scenario is still ok.

Question/Challenge: Can we find a ROBUST
phasetransition at high μ e.g. such as in the
nuclear liquid gas case?

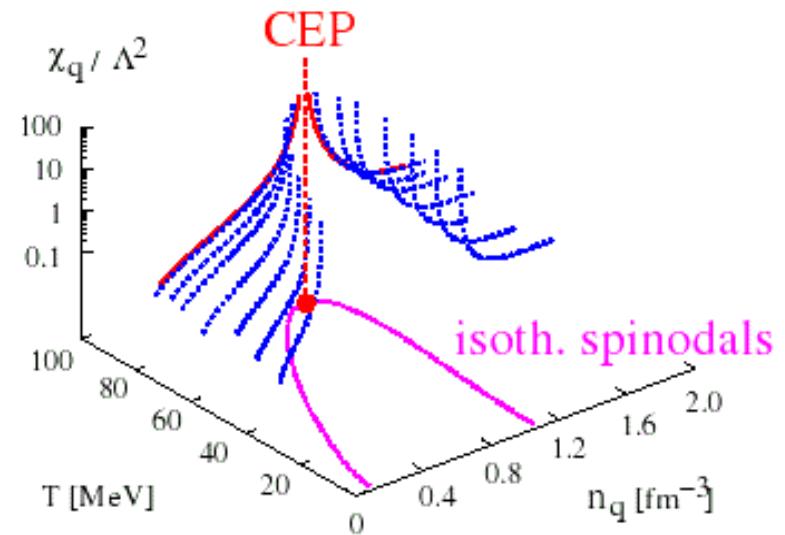
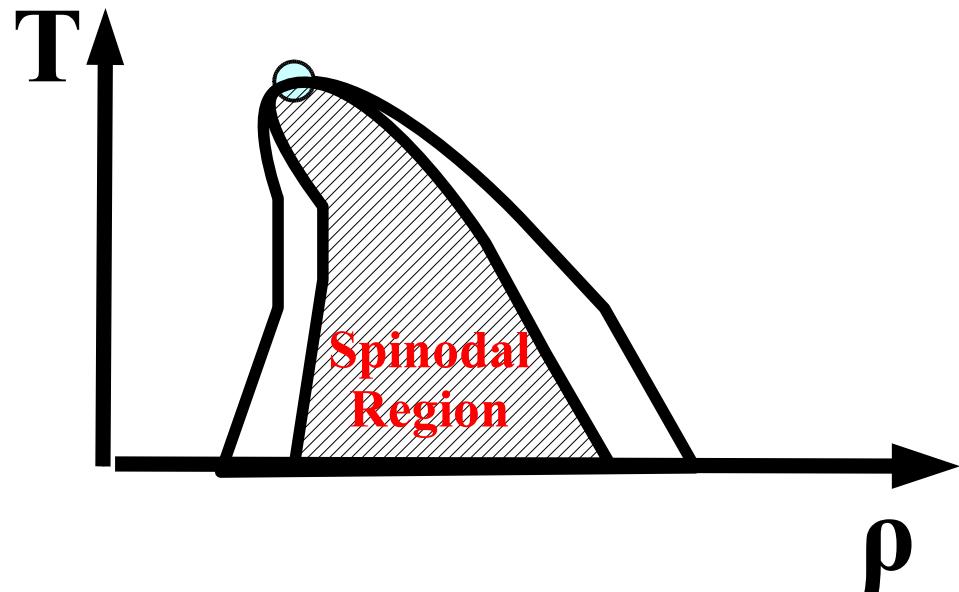


See P. de Forcrand, Parallel Session XI

QCD critical point

- Order parameter: baryon density or scalar density
 - Actually it is a superposition
- Both scalar (chiral) an quark number susceptibilities diverge
- **Screening** (“space like”) masses vanish (“omega”, “sigma”)
 - not accessible by (time-like) dileptons
- Is it related to chiral transition at $m_q = 0$?
 - The transition is in same universality class as liquid gas! (Son, Stephanov)
 - Fluctuations are driven by density fluctuations; chiral field is just tagging (H. Fuji, parallel Session IV)
 - CP “just” the end of of 1st Order transiition
 - Spinodal instabilities

Spinodal region



→ Inside spinodal region: mechanical instability
formation of “blobs” of typical size;
(may explain K/pi flcuts, Majumder et al)

Baryon susceptibility diverges at spinodal boundary
(Sasaki et al, see also Sasaki, parallel session X)

Things to do!

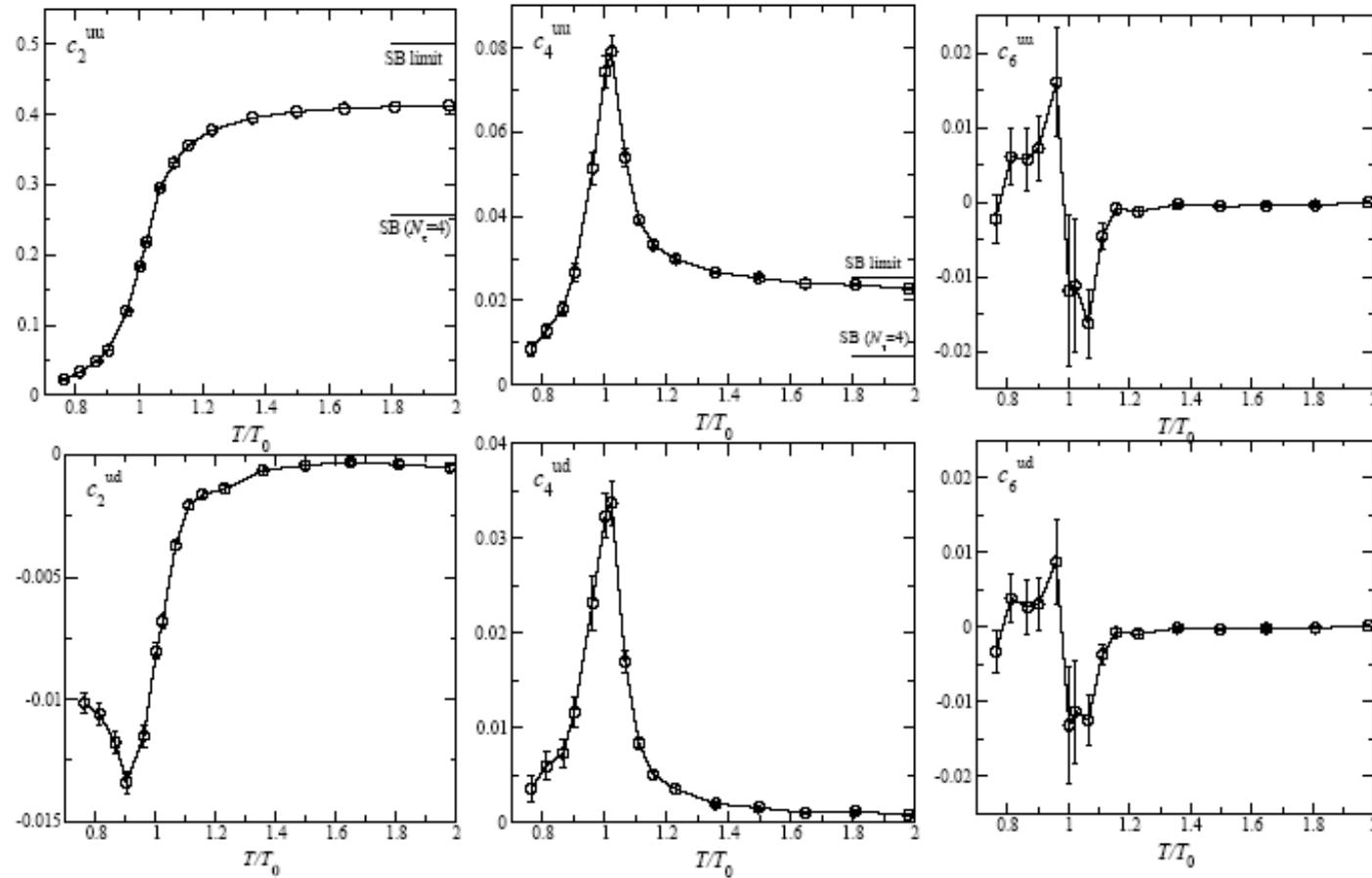
- Characterize the Phases
 - what are useful order parameters
 - what observables couple well to order parameter but survive hadron phase
- Test observables using static and dynamical models
 - Effects are small, comparable with “trivial ones” such as quantum statistics, dynamics etc.
 - Only a well chosen observable / set of observables will prevent us from seeing Poisson
 - e.g. can we live without neutrons?
 - CONSERVATION LAWS
- **Don't forget the first order phase co-existence!**

Observables for CP and co-existence

- Fluctuations (probably not of conserved charges)
- Correlations (spiondal blobs)
- Energy scan
- System size dependence (finite volume scaling)
 - centrality may not do
- Good, HOMOGENIOUS acceptance
- Be prepared to measure everything
 - not clear (yet?) which observable couples strongest to baryon density
 - want to see finite volume scaling in more than one observable
- Get you hands dirty now!

Lattice-QCD susceptibilities

This should be in your present (RHIC) dataset !



Rule of thumb:

$$c_n \sim \langle X^n \rangle$$
$$X = B, Q, S, \dots$$

Alton et al, PRD 66 074507 (2002)

Critical Point and viscosities

CP is in universality class of liquid gas (Son, Stephanov)

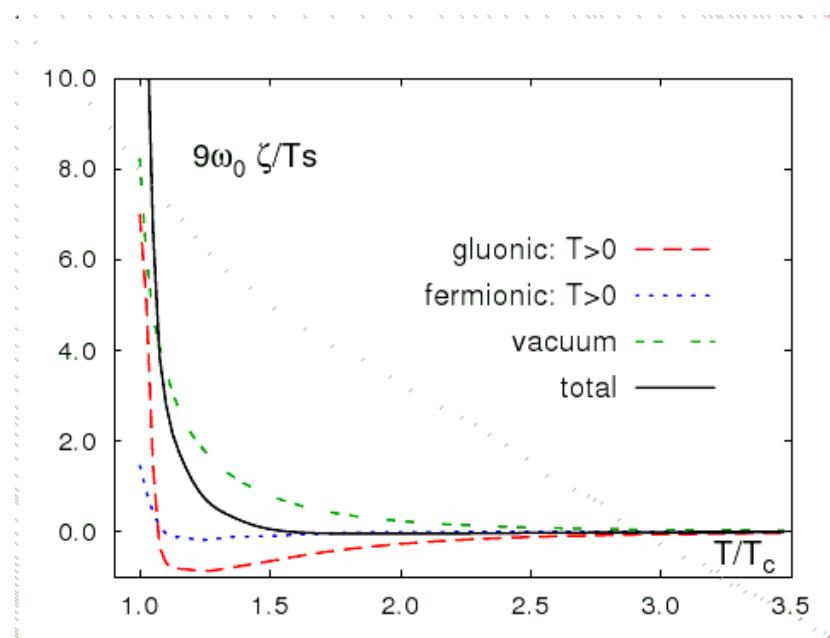
Hohenberg - Halperin Model H (Rev. Mod. Phys 49 (1977)):

$$\eta \sim \xi^{0.065}, \quad \xi = \text{Correlation Length}$$

Shear viscosity **diverges** at CP

Bulk viscosity also **diverges**:
(Kharzeev, Turchin, Karsch arXiv:0711.0914)

Note: even large increase without PT
due to vacuum contribution





INT Summer Program
The QCD Critical Point
July 28 - August 22, 2008

Organizers: G. Roland, M. Stephanov and V.K.

The End

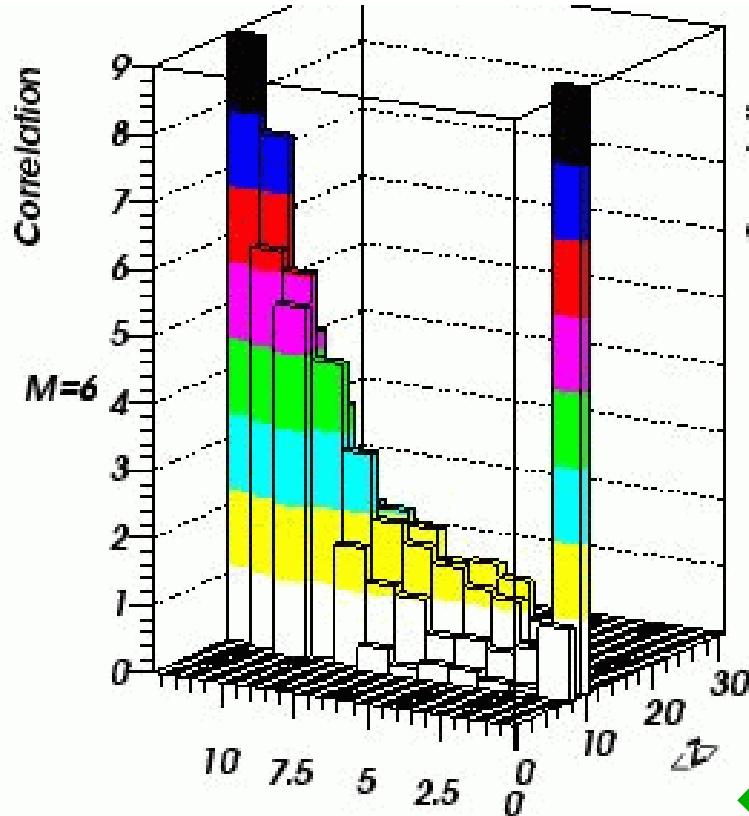
Charge Fluctuations and Balance Functions

$$\frac{\langle \delta Q^2 \rangle_{\Delta\eta}}{\langle N_{\text{ch}} \rangle_{\Delta\eta}} \approx 1 - \int d\eta B(\eta|\Delta\eta)$$

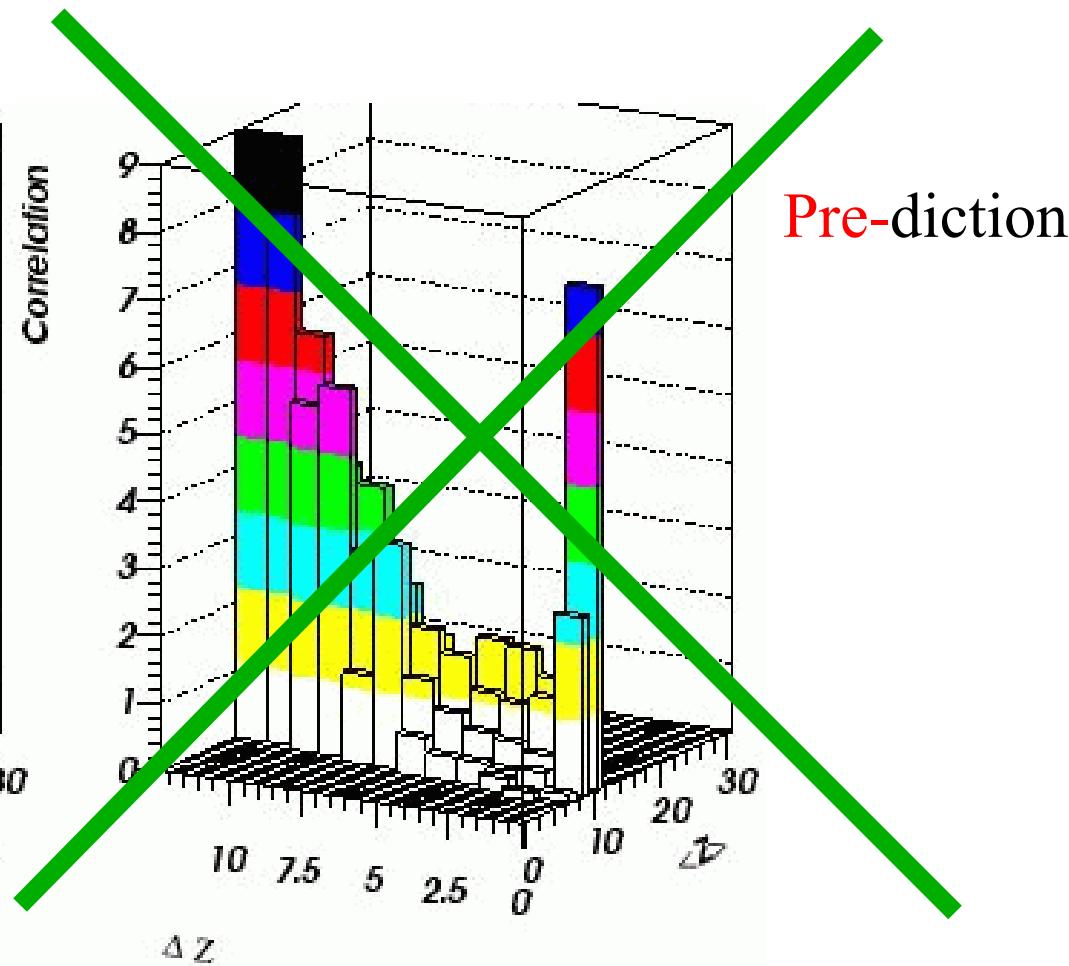
Spinodal decomposition in nuclear multifragmentation

occurs!

Data speak for themselves!



Experiment (INDRA @ GANIL)
Borderie et al, PRL 86 (2001) 3252



Theory (Boltzmann-Langevin)
Chomaz, Colonna, Randrup, ...

Correlation length and correlation functions

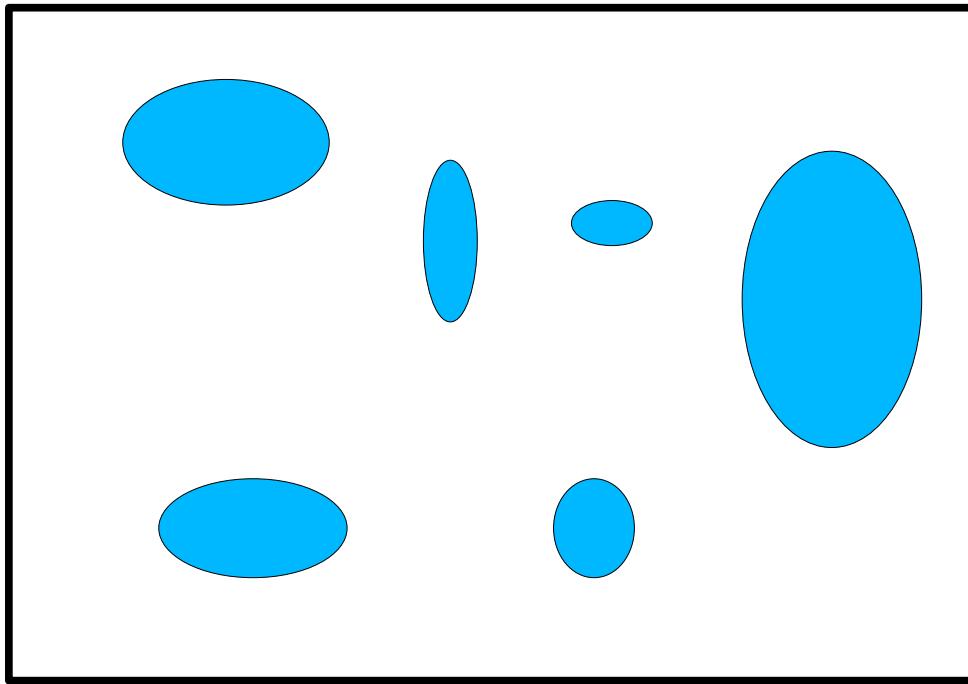
Theory: $\langle \rho(r) \rho(0) \rangle$ Correlation in **coordinate** space

Experiment: $\frac{dN}{dp_1 dp_2} \sim \langle \tilde{\rho}(p_1) \tilde{\rho}(p_2) \rangle$ Correlation in **momentum** space

Flow helps maybe....

Better: Calculation of momentum space correlations

Baryon number fluctuations (the obvious choice?)



Strong spatial fluctuations

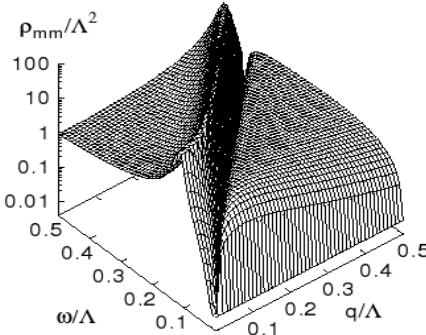
If $V_{\text{domain}} \ll V$, small effect
on integrated Baryon Number
fluctuations

$$\frac{\langle (\delta N)^2 \rangle}{\langle N \rangle} \approx \left(1 + \frac{(\Delta \rho)^2}{4 \bar{\rho}^2} \right)$$

CP and the chiral transition

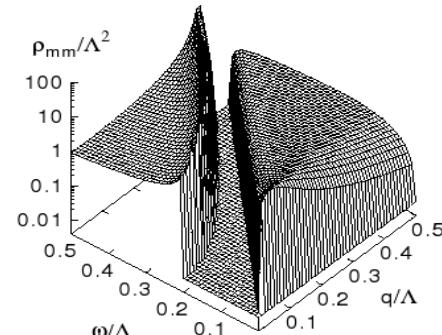
Chiral transition $m_q = 0$

$T > T_c$



(a)

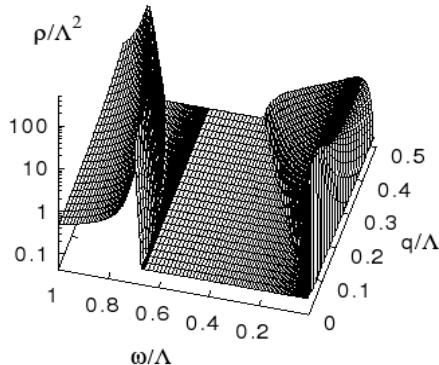
$T < T_c$



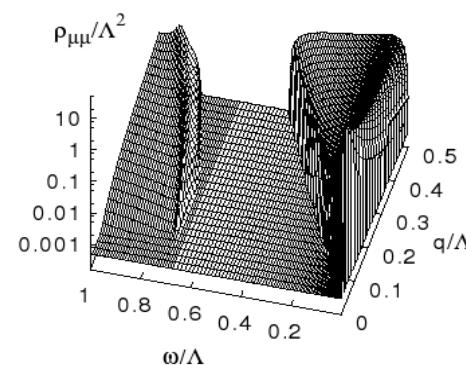
(b)

Critical point $m_q > 0$

Massive Scalar

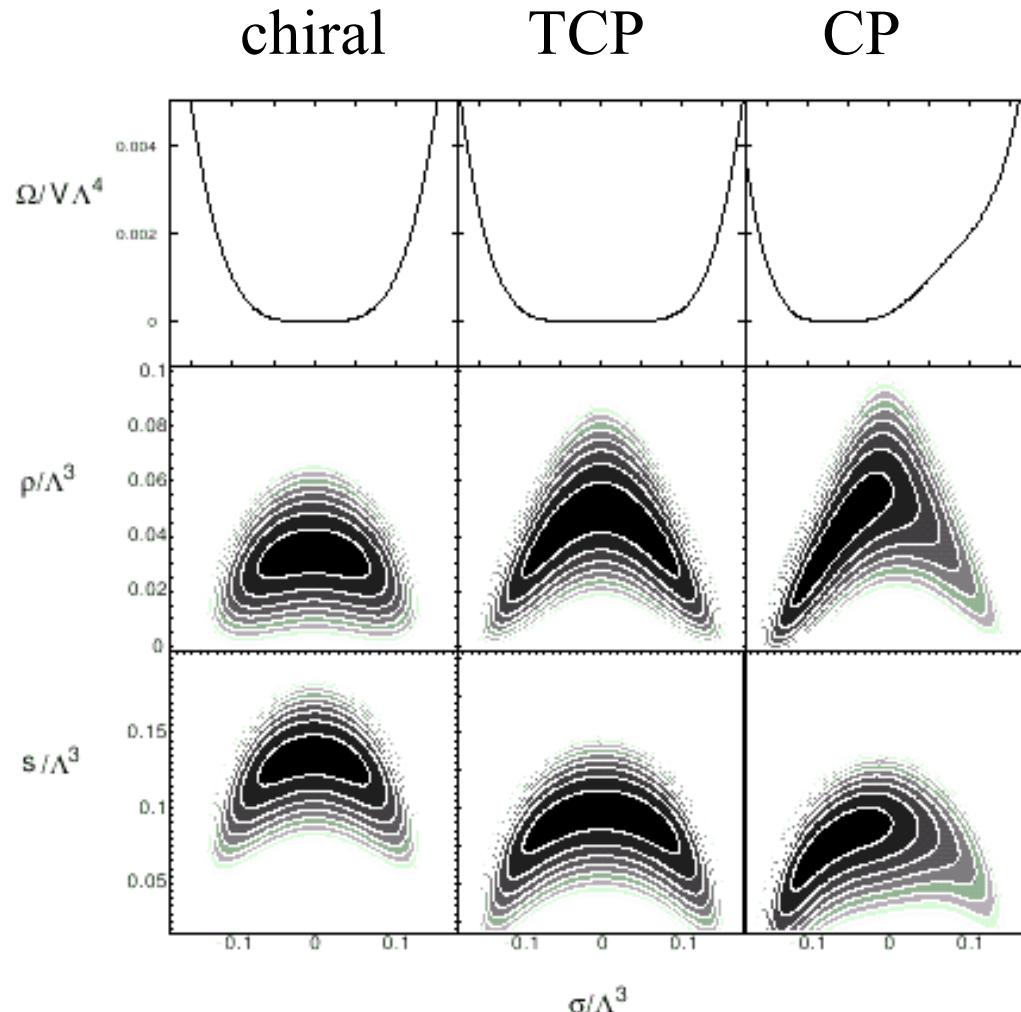


Scalar



Vector (Baryon-density)

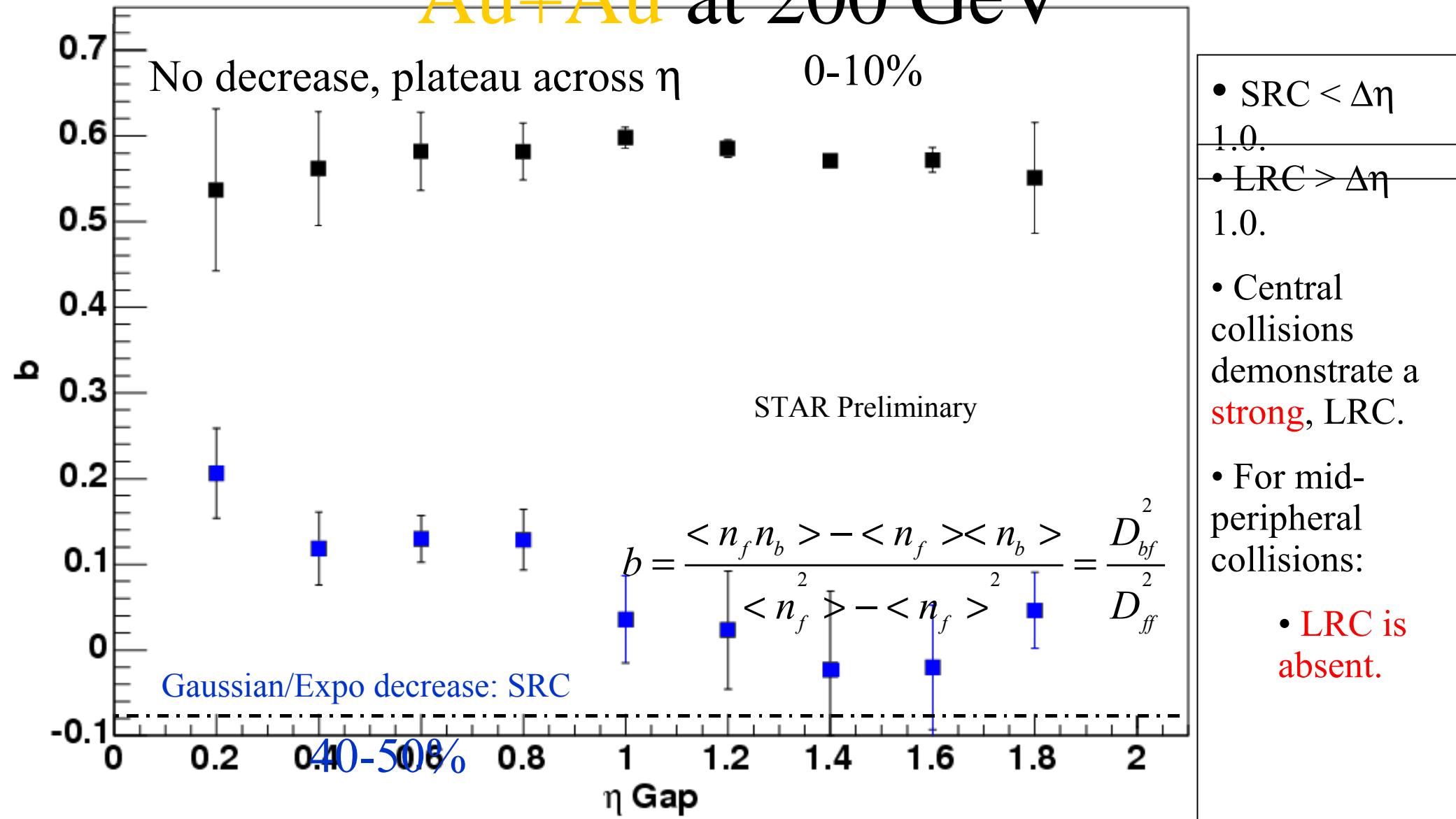
Mixing of scalar and vector



Note: mixing of scalar
vector (p-h) only for
spacelike excitations
close to $q=0$
(Friman, Soyeur et al)

Note: entropy also mixes
in!
Multiplicity fluctuations
Trigger?

Centrality Dependence: Au+Au at 200 GeV



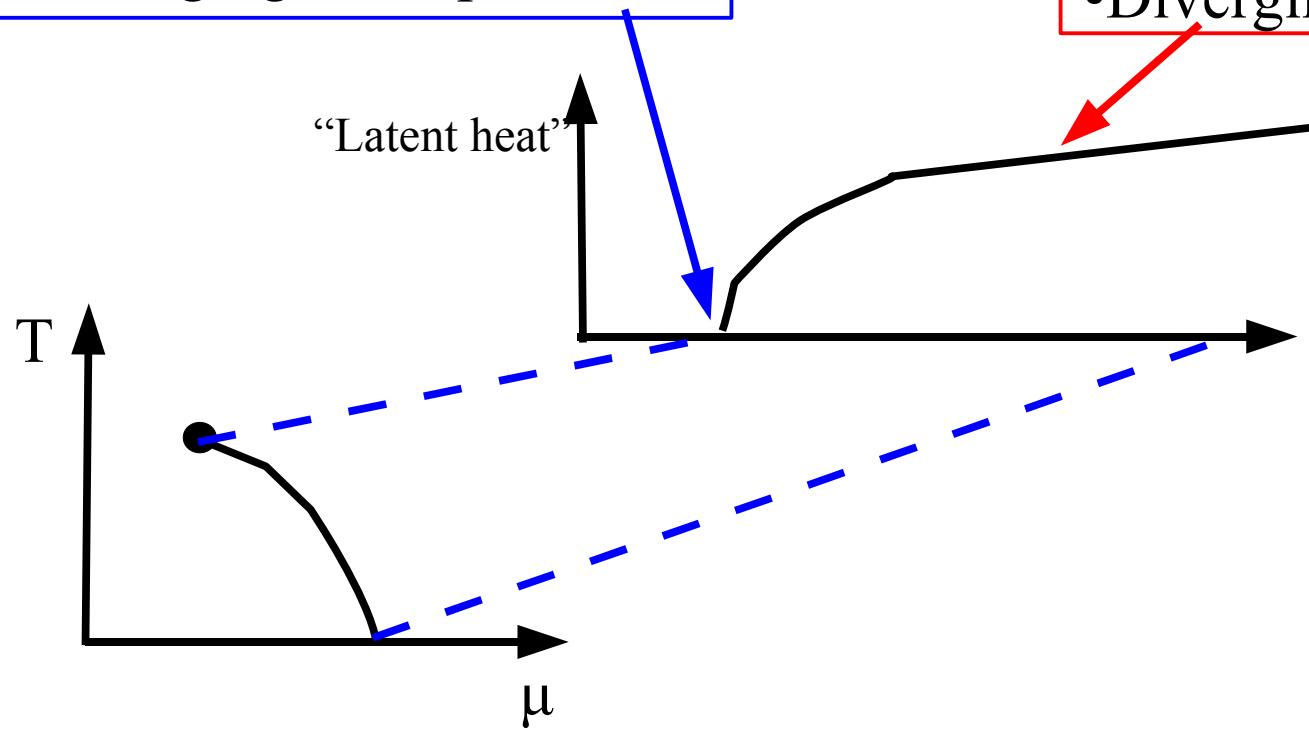
First order or second order?

Second order:

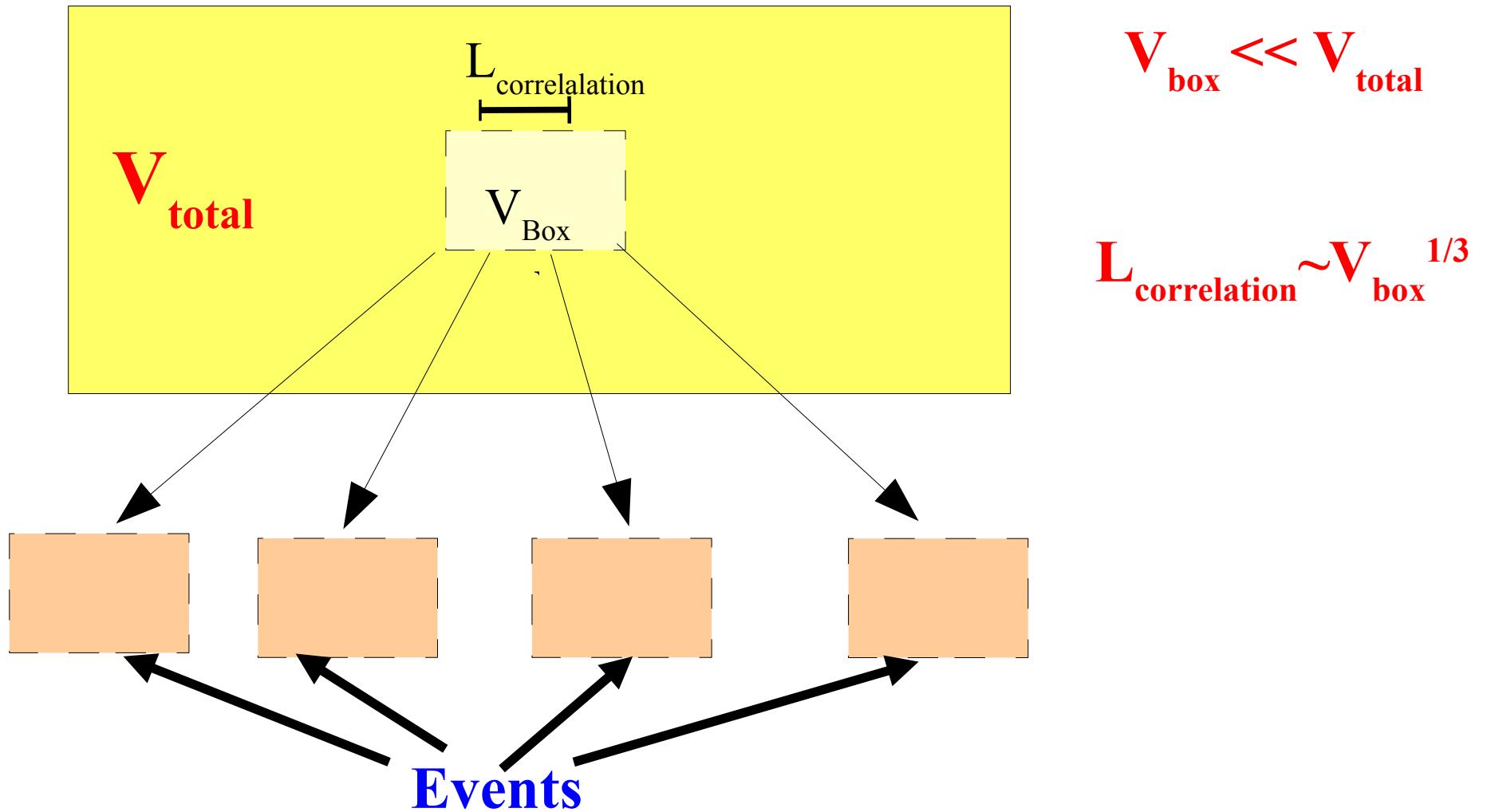
- Critical fluctuations
- Diverging Susceptibilities

First order:

- Phase coexistence, bubbles
- Spinodal instabilities
- Diverging Susceptibilities



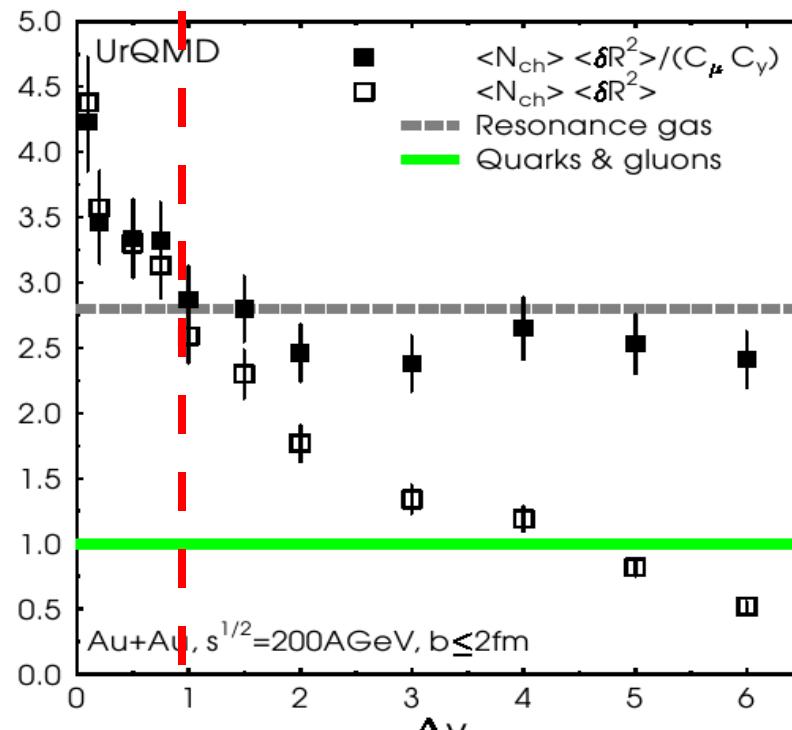
“Charge” Fluctuations



Charge conservation

$$C_\mu = \tilde{R}_{\Delta y}^2 = \frac{\langle N_+ \rangle_{\Delta y}^2}{\langle N_- \rangle_{\Delta y}^2}$$

$$C_y = 1 - P = 1 - \frac{\langle N_{ch} \rangle_{\Delta y}}{\langle N_{ch} \rangle_{total}}$$



Correlations (meson decay)

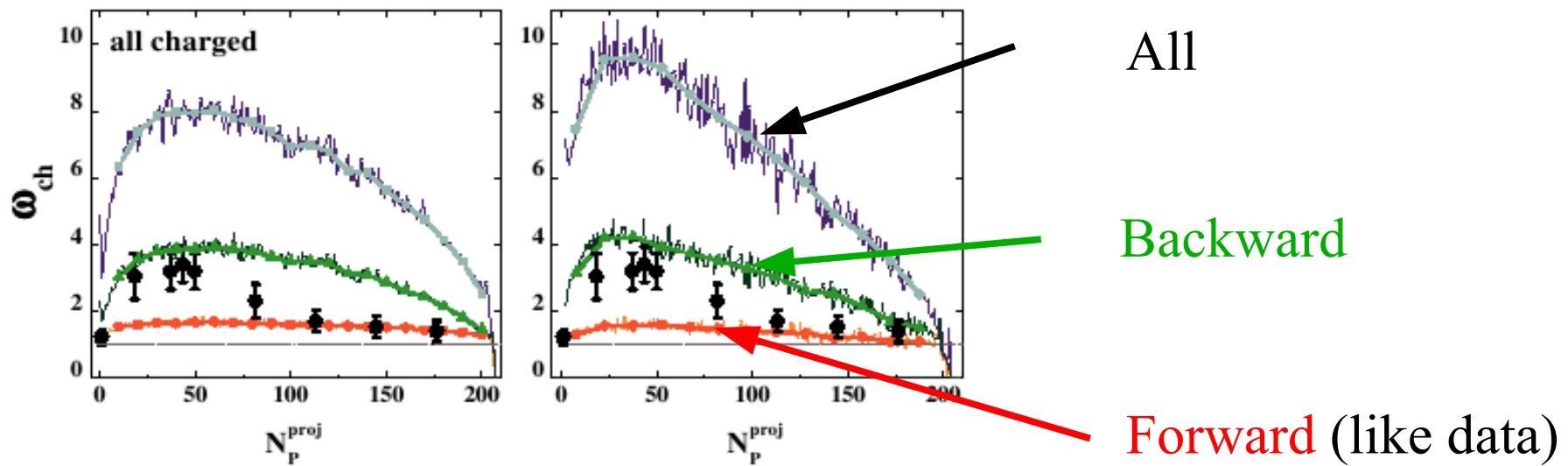
Bleicher et al.
Phys.Rev.C62:061902,2000

QCD first order transition

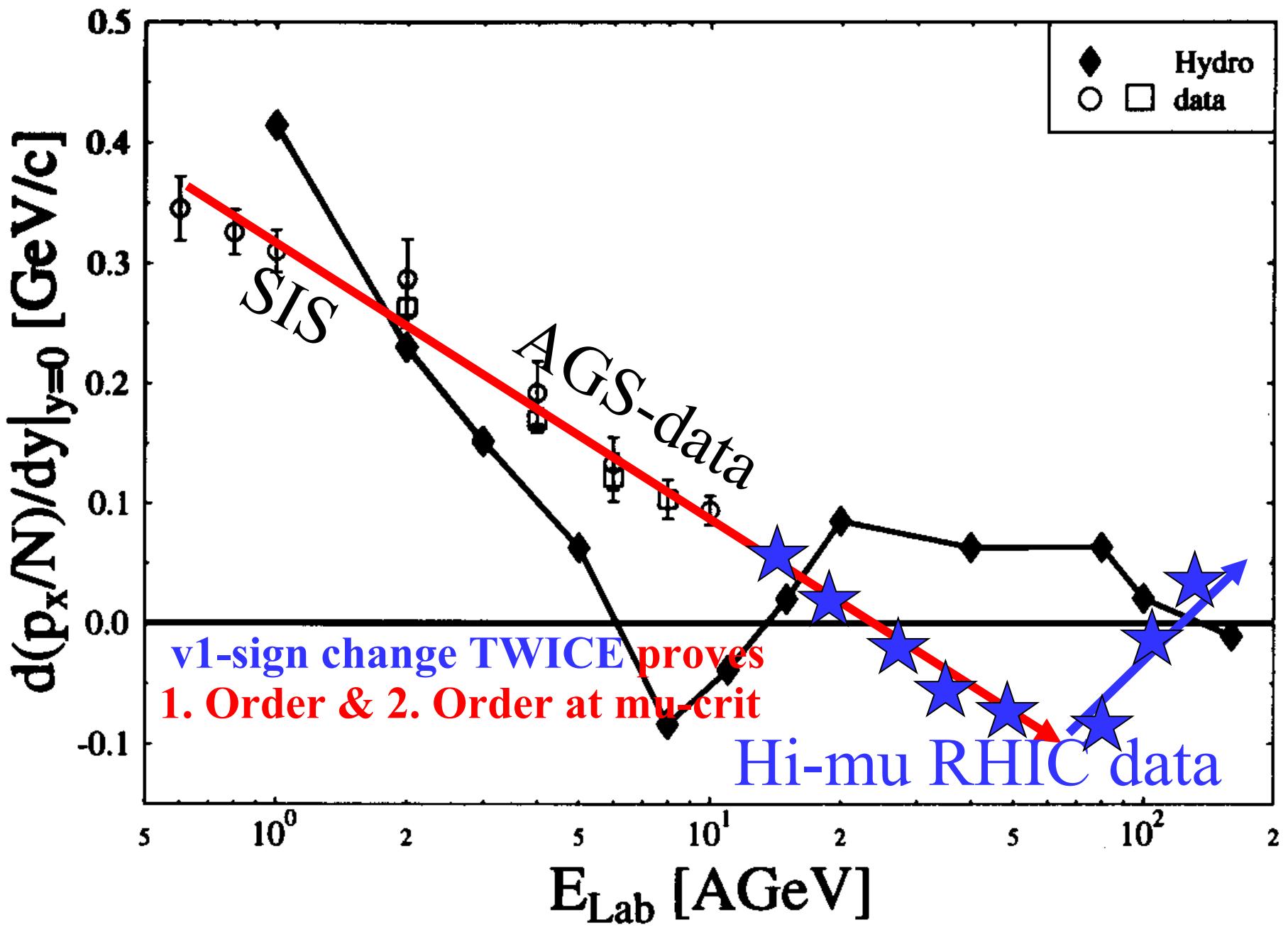
- Order parameter: baryon density or scalar density
- Both scalar (chiral) and quark number density are discontinuous
- Density fluctuations / blobs due to spinodal instability.
- If latent heat is sizable: Long lived mixed phase
 - two peak structure of omega a la C.M. Ko for phi.
 - assuming the TAPS results hold up in this case you want to measure **low p_t** omegas

Dynamics, event selection ... (or why a symmetric detectors are good)

Konchakovski et al, nucl-th/0511083



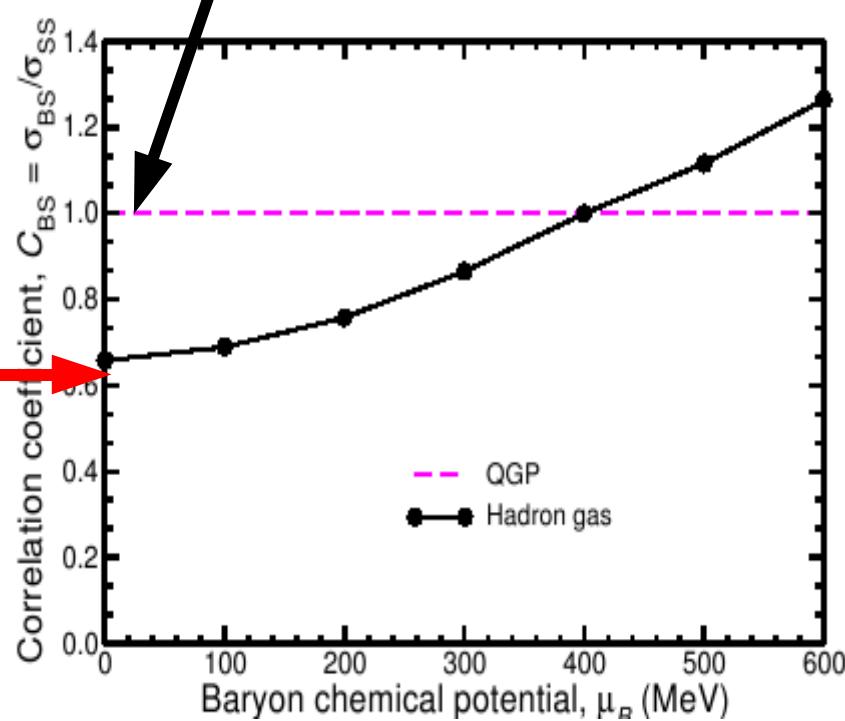
- Fluctuations are sensitive to dynamics (mixing of projectile and target material?)
- Event selection/trigger affects fluctuations → large Acceptance!



Paech, Dumitru

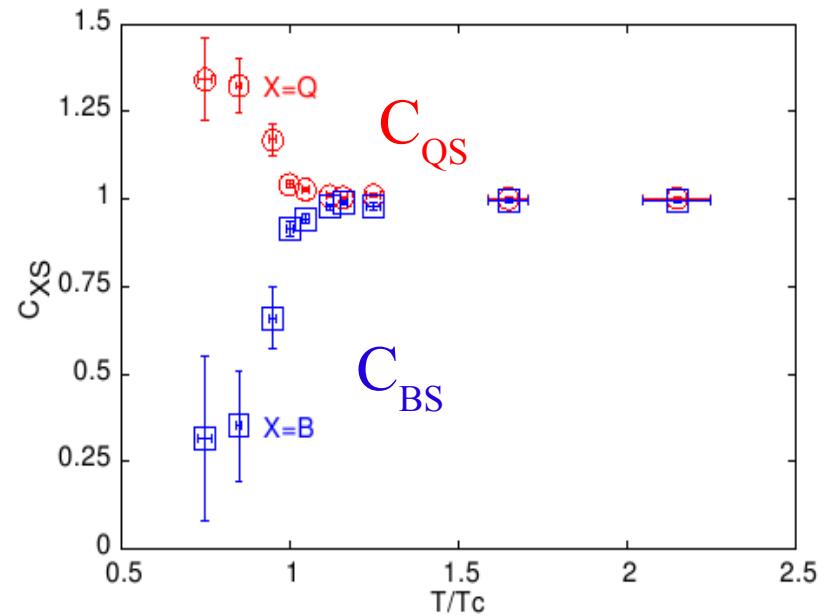
Something else: $\langle BS \rangle$, $\langle QS \rangle$

Independet quarks and
LATTICE QCD for $T > 1.1 T_c$



$$C_{BS} = -3 \frac{\langle BS \rangle}{\langle S^2 \rangle}$$

$$C_{QS} = -3 \frac{\langle QS \rangle}{\langle S^2 \rangle}$$

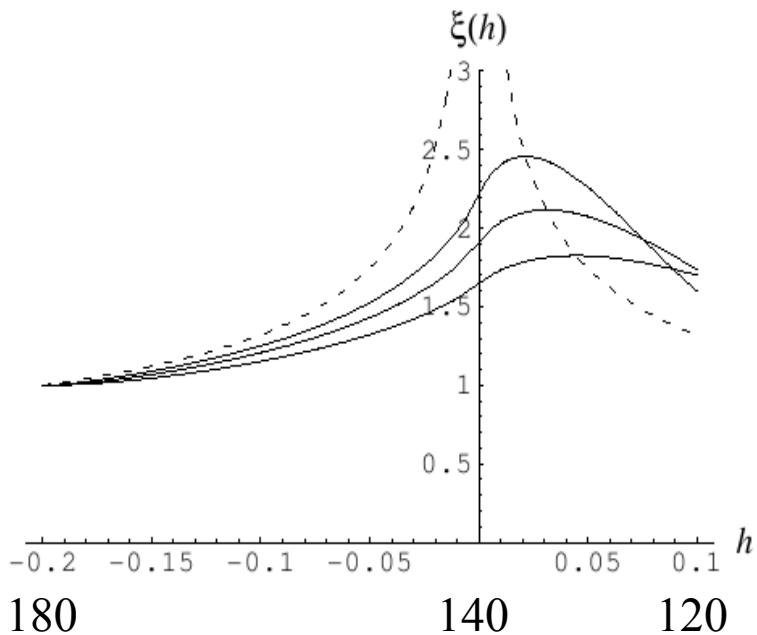


V.K. Majumder, Randrup PRL95:182301,2005

Gavai,Gupta, hep-lat/0510044

Second order

correlation length $\sim 1/m_\sigma$



- Critical slowing down
- limited sensitivity on model parameters
- Max. correlation length 2-3 fm
- Translates in **3-5%** effect in p_t -fluctuations

Bernikov, Rajagopal, hep-ph/9912274