



# Quenching of light hadrons at RHIC In a collisional energy loss scenario

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# Plan

- **Motivation**
- **Present status of collisional energy loss**
- **Our formalism**
  - **Energy loss**
  - **Nuclear modification factor**
- **Results**
- **Discussion**



# Motivation

- **Non-photonic electron data shows much larger suppression than expected**
- These data reflect the energy loss of heavy quarks
- **Mainly radiative loss has been discussed in the literature : dead cone effect => less energy loss => less suppression of the decay product coming from heavy quarks**
- Radiative energy loss seems to fail to describe the non-photonic single electron data (Wicks et. al, JPG 34, 2007)
- **Activities started to re-visit the importance of collisional energy loss for heavy as well as light quarks**



# Present status of collisional energy loss

## 1. Peigne et al (JHEP 04, 2006)

Collisional energy loss is suppressed in comparison to infinite medium

## 2. Djordjevic (PRC 74, 2006)

For characteristic QCD medium scales finite size effects are negligible

## 3. Wang X.- N. (PLB650, 2007)

Interference between elastic amplitude and that of gluon radiation reduces the effective energy elastic loss

## 4. Adil et. Al (PRC75 2007)

Similar results for collisional loss as 2



## Contd.

5. Dutt-Mazumder et. al (PRD 74 2005)

Below certain energy ( $E_c$ ) collisional energy loss dominates (details will be shown)

6. Peshier et al. (EPJC 2007), Braun et al. (PRD75 2007)

running coupling constant  $\alpha_s = \alpha_s(p, T)$

$T \sim 2T_c$ ,  $dE/dx \sim 1.6$  GeV/fm for a 100 GeV quark

Collisional energy loss is of the same order as radiative loss

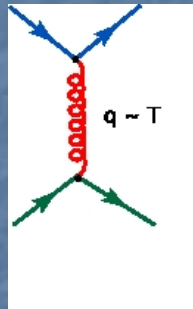


# Collisional energy loss



$$\frac{dE}{dx} = \frac{1}{v} \int dE' (E - E') \frac{d\Gamma}{dE'}$$

**Bjorken**



$$\frac{dE}{dx} = \frac{8\pi}{3} \left(1 + \frac{N_f}{6}\right) \alpha_s^2 T^2 \ln \frac{q_{\max}}{q_{\min}}$$

**Infra-red divergent**

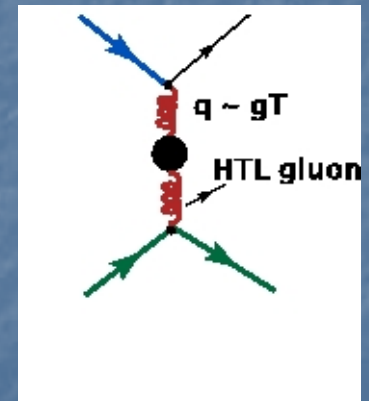
**Way out : Shielding of divergent by plasma effects**

**Sum of two diagrams separated by intermediate scale  $q^*$  ( $gT \ll q^* \ll T$ )**

**(i) Soft scale ( $gT < q < q^*$ ) : HTL propagator**

**(ii) Hard scale ( $q^* < q < T$ ) : bare propagator**

**Cancellation of arbitrary scale  $q^*$**





# Energy loss contd.



- Consider small angle qq scattering (t-channel)
- Assume energy  $E_{\text{jet}} \gg T$

$$\left(\frac{dE}{dx}\right)_{\text{hard}} \sim \alpha_s^2 T^2 c_{qq} \ln \frac{ET}{q^{*2}}$$

Bare gluon propagator

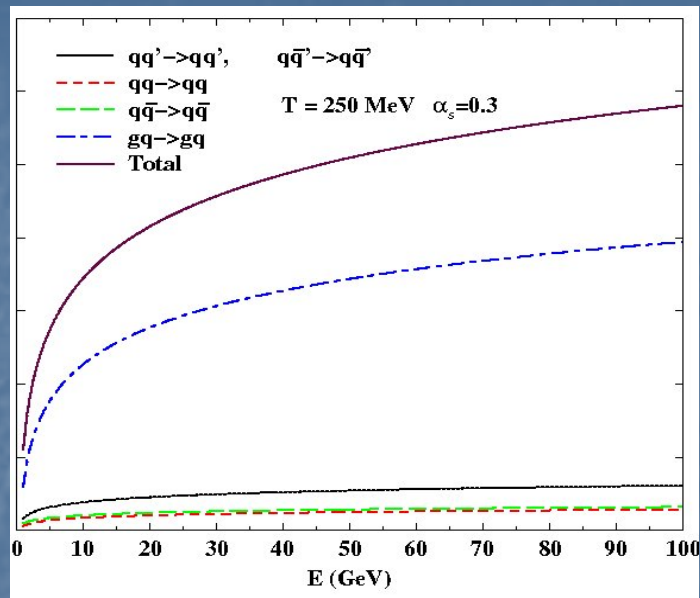
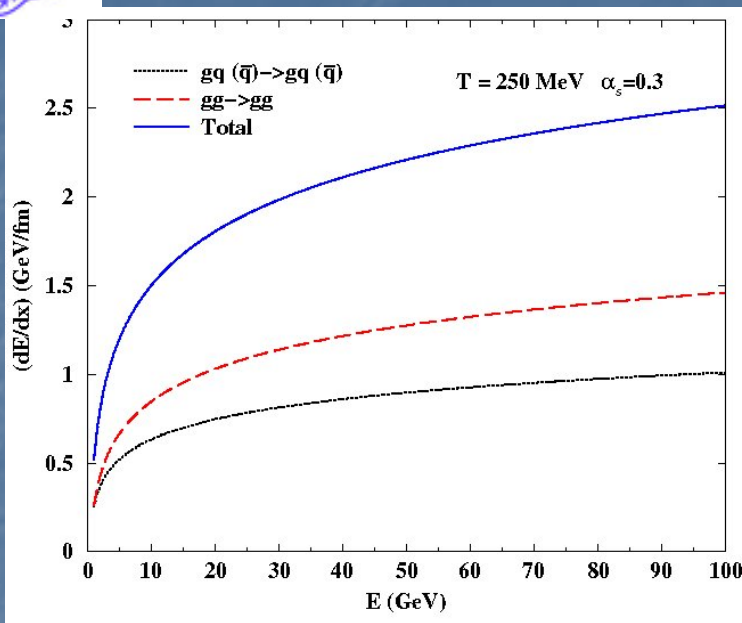
and

$$\left(\frac{dE}{dx}\right)_{\text{soft}} \sim \alpha_s^2 T^2 c_{qq} \ln \frac{q^{*2}}{g^2 T^2}$$

HTL gluon propagator

Hard + soft : independent of  $q^*$

$$\left(\frac{dE}{dx}\right)_{\text{soft}} \sim \alpha_s^2 T^2 c_{qq} \ln \frac{E}{g^2 T}$$



$\frac{dE_q}{dx}$  is of the order of 0.8 GeV/fm for a 20 GeV quark

**Factor of 2 -3 more than previously calculated**  
**Justifies not to neglect the collisional energy loss**





# Collisional vs. Radiative

(PRD 74 2005)



$$\Delta E_{GLV} = \frac{C_F \alpha_s}{N(E)} \frac{L^2 \mu^2}{\lambda} \ln\left(\frac{2E}{\mu^2 L}\right),$$

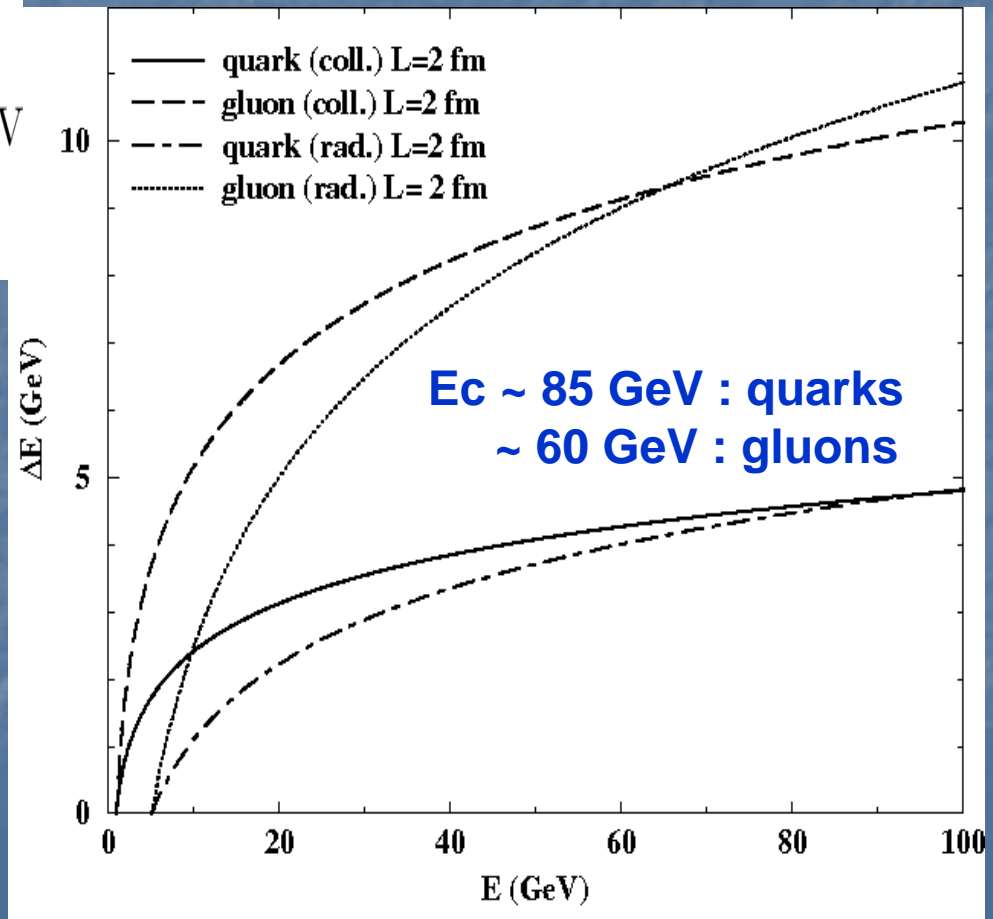
Important!!!!

$N(E) = 7.3, 10.0, 24.4$  for  $E = 500, 50, 5$  GeV

$N(E \rightarrow \infty) = 4$

there might be domains  
where collisional loss could  
be comparable to radiative  
loss

$N(E) = 10$   
 $\mu = 1$  GeV  
 $L/\lambda = 4$





# Transverse momentum distributions hadrons : Nuclear modification factor

(Roy et al. PRC 2006, Alam et al. NPA 2007)

**$p_T$  distribution :**

Standard pQCD calculation : modification of the

$$\text{FF} \quad D(z, Q^2) : z \rightarrow z^* = \frac{z}{1 - \Delta z} : \Delta z = \frac{\Delta E}{E}$$

**equal amount of energy loss for each parton**

**Essential to evaluate parton  $p_T$  spectra  
dynamically : FOKKER PLANCK  
APPROACH**

- **System of quarks, antiquarks, and gluons**
- **Inject partons with given initial distributions**
- **Study of time evolution as the system expands and cools**

**Boltzmann Equation**

- **Small angle scattering  
more frequent  $\sim O(g^2T)$**

**Collisional integral  
reduces to appropriately  
defined diffusion and  
Drag coefficients**



# Fokker-Planck Equation

$$\frac{\partial f}{\partial t} - \frac{p_z}{t} \frac{\partial f}{\partial p_z} = \frac{\partial(\eta \mathbf{p} f)}{\partial \mathbf{p}} + \frac{\partial(Df)}{\partial \mathbf{p}^2}$$

**Solution by Green's function technique:**

$$E \frac{dN}{d^3 p^f} = \int d^3 p_0^f G(p^f, t | p_0^f, t_i) E_0 \frac{dN}{d^3 p_0^f}$$

**All the jets are not produced at the same space time point**

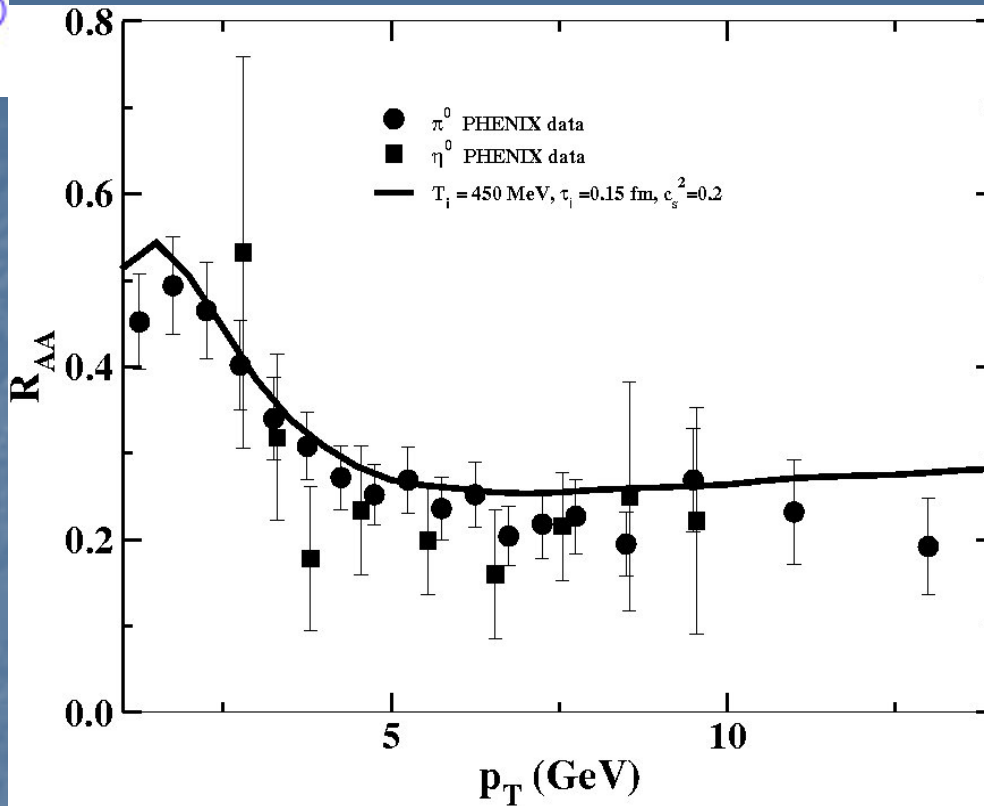
**Path length traversed by these partons before fragmentation are not same**

**Temperature also varies along the trajectory of the partons**

**Jet produced at  $(r, \phi)$  spends a time  $t_L$  or equivalently traverses a distance  $L \sim t_L$  and this is not a measurable quantity : time average**

**Finally,**

$$\frac{dN^{\pi^0(\eta)}}{d^2 p_T dy} = \sum_f \int d^2 r \mathcal{P}(r) \int_{t_i}^{t_L} \frac{dt}{t_L - t_i} \int \frac{dz}{z^2} D_{\pi^0(\eta)/f}(z, Q^2)|_{z=p_T/p_T^f} \times E \frac{dN}{d^3 p^f}$$



# Nuclear modification factor

$$R_{AA}(p_T) = \frac{\frac{dN_{AA}^{\pi^0(\eta)}}{d^2p_T dy}}{\left[ \frac{dN_{AA}^{\pi^0(\eta)}}{d^2p_T dy} \right]_0}$$

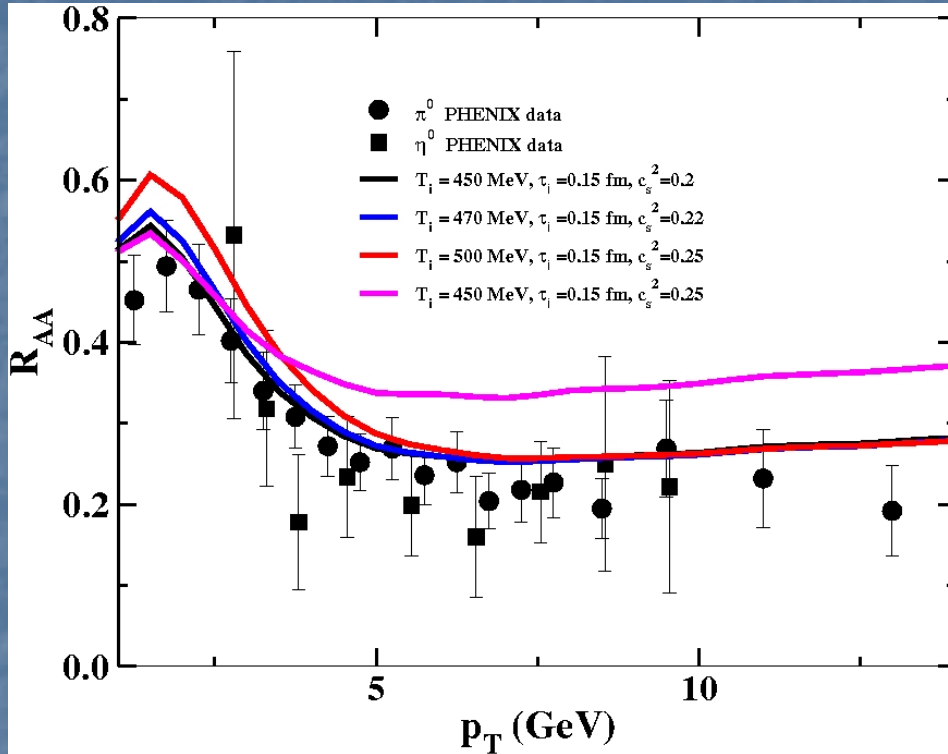
Fits the Phenix data quite well

Tendency to rise at high  $p_T$  : radiation starts to dominate

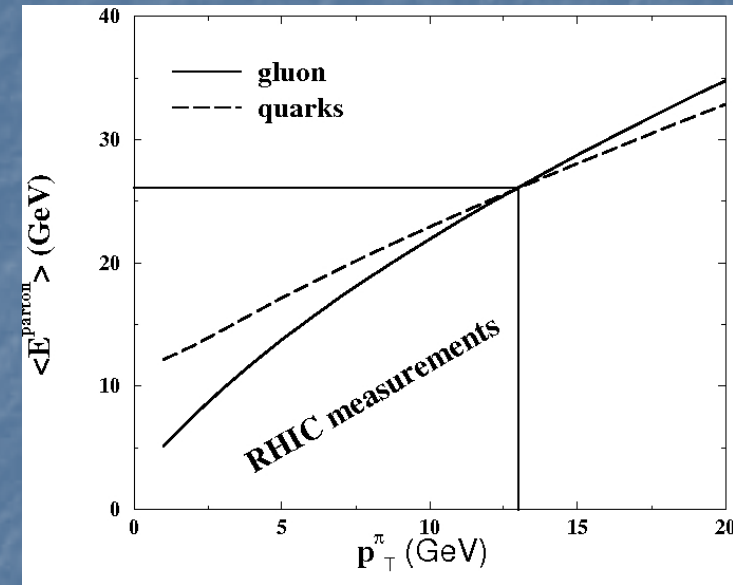
Detail calculation including both collisional and radiative losses should be considered for quenching of high  $p_T$  hadrons



$R_{AA}$  is sensitive to **initial temperature**  
**velocity of sound** initial thermalization  
 time



Reasonable combination :  
 $T_i = 450$  MeV,  $c_s^2 = 0.2$ , and  
 $t_i = 0.15$  fm/c



Importance of coll. and rad. losses :  
 the average energy of parton  $\langle E \rangle = p_T^h / \langle z \rangle$

$\langle E \rangle_{max} \sim 26$  GeV : maximum average energy to produce pions  
with  $p_T$  in the range of 1-13 GeV



# Summary and Discussion

- Reviewed the current theoretical status of collisional energy loss
- Divergence is removed by plasma effects
- Included all possible diagrams for a given process
- Domains in energy: collisional loss dominates over radiative energy loss
- Nuclear modification factors for  $\pi^0$  and  $\eta$  : fragmentation of dynamically evolved parton distribution
- With parameters relevant for RHIC the data is well reproduced
- Not surprising because for RHIC data quenching factor  $Q(p_T) \sim 0.2$  where collisional energy loss plays important role

**THANK YOU**